

QChips 2026 (Gdansk, Poland, August 26, 2026)

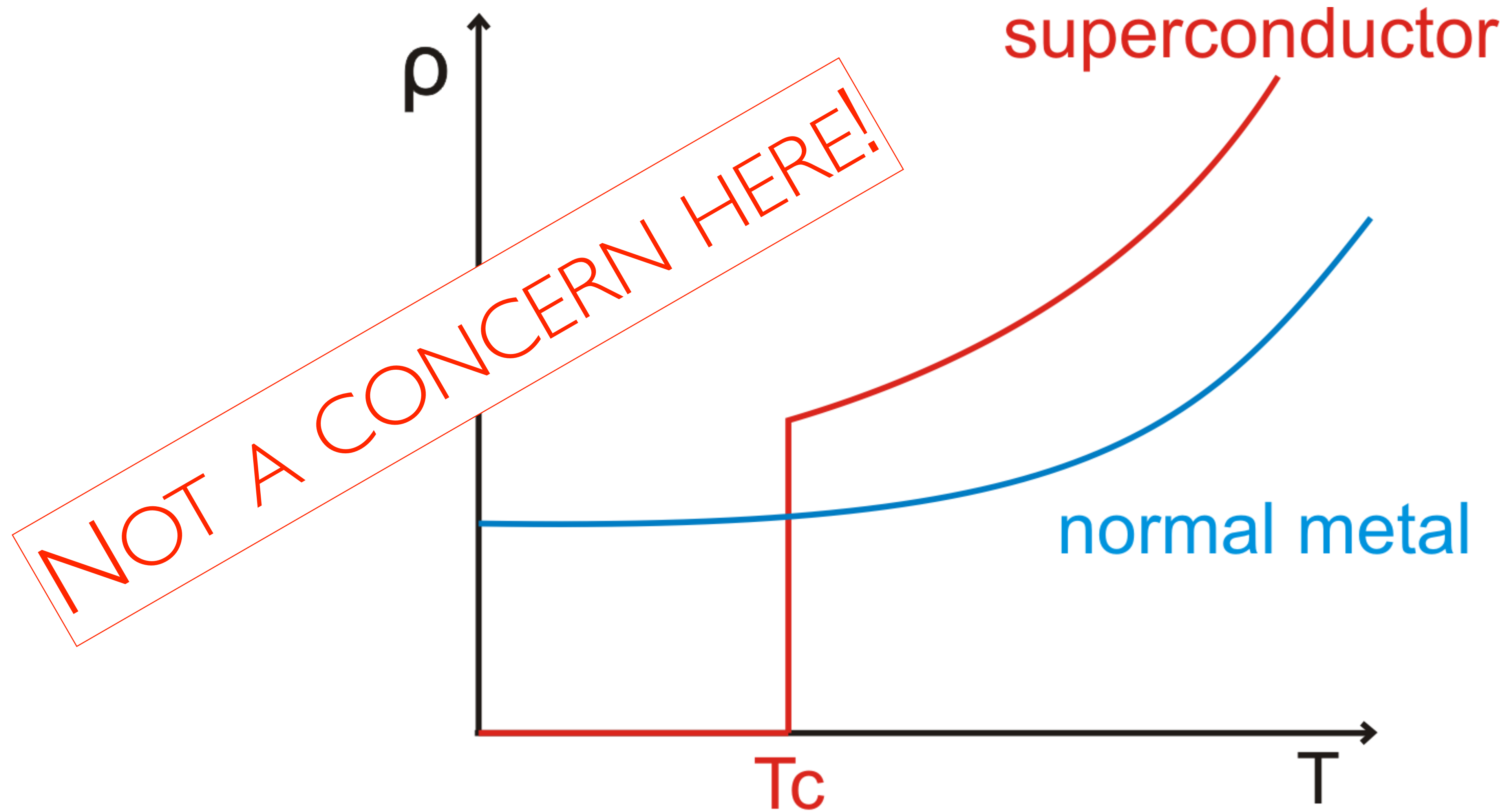
ULTRA-STRONG COUPLING PHYSICS IN CIRCUIT-QED SYSTEMS

MAHN-SOO CHOI
(KOREA UNIVERSITY)

SUPERCONDUCTIVITY

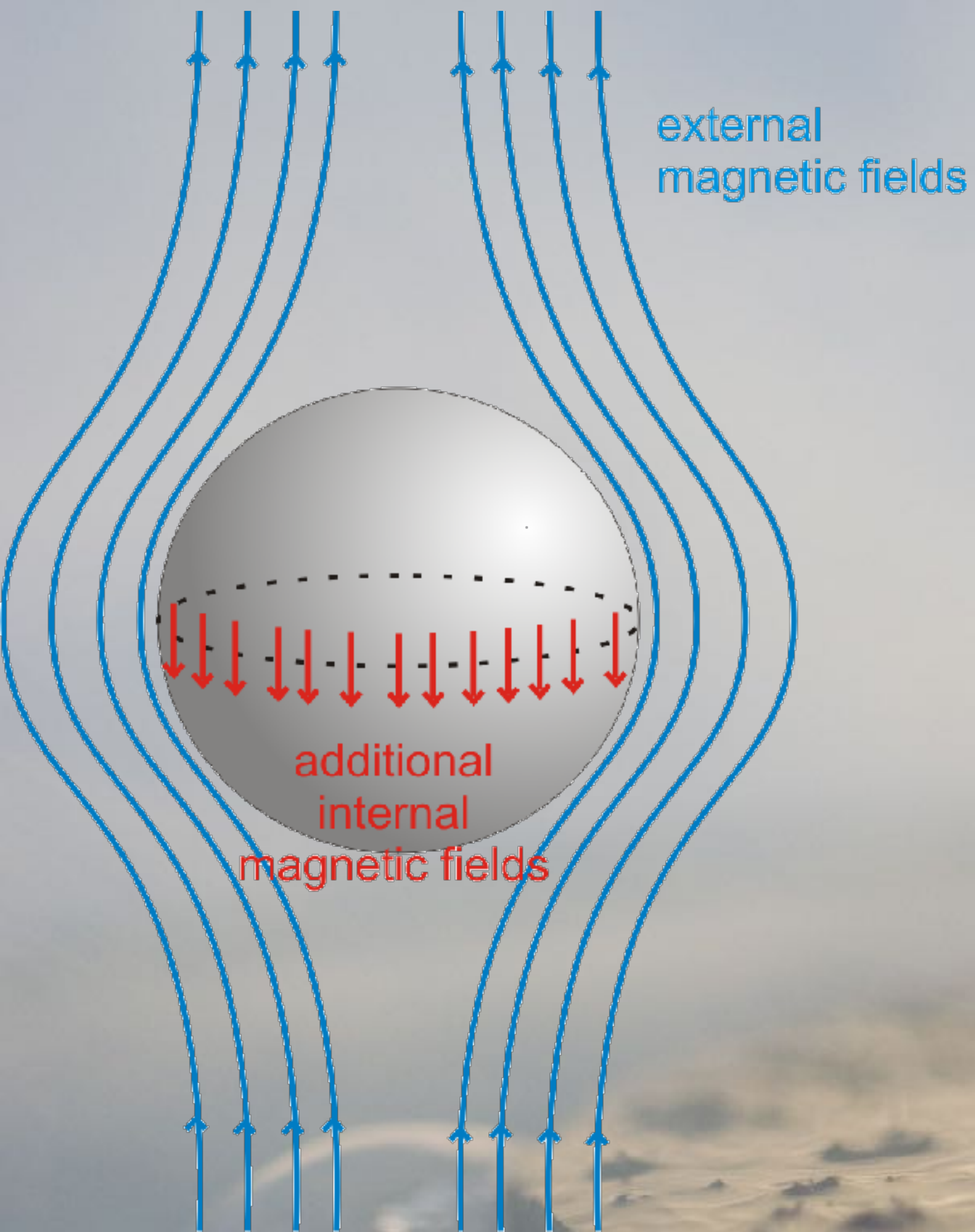
PHENOMENA

PERFECT CONDUCTIVITY



PERFECT DIAMAGNETISM

MEISSNER EFFECT



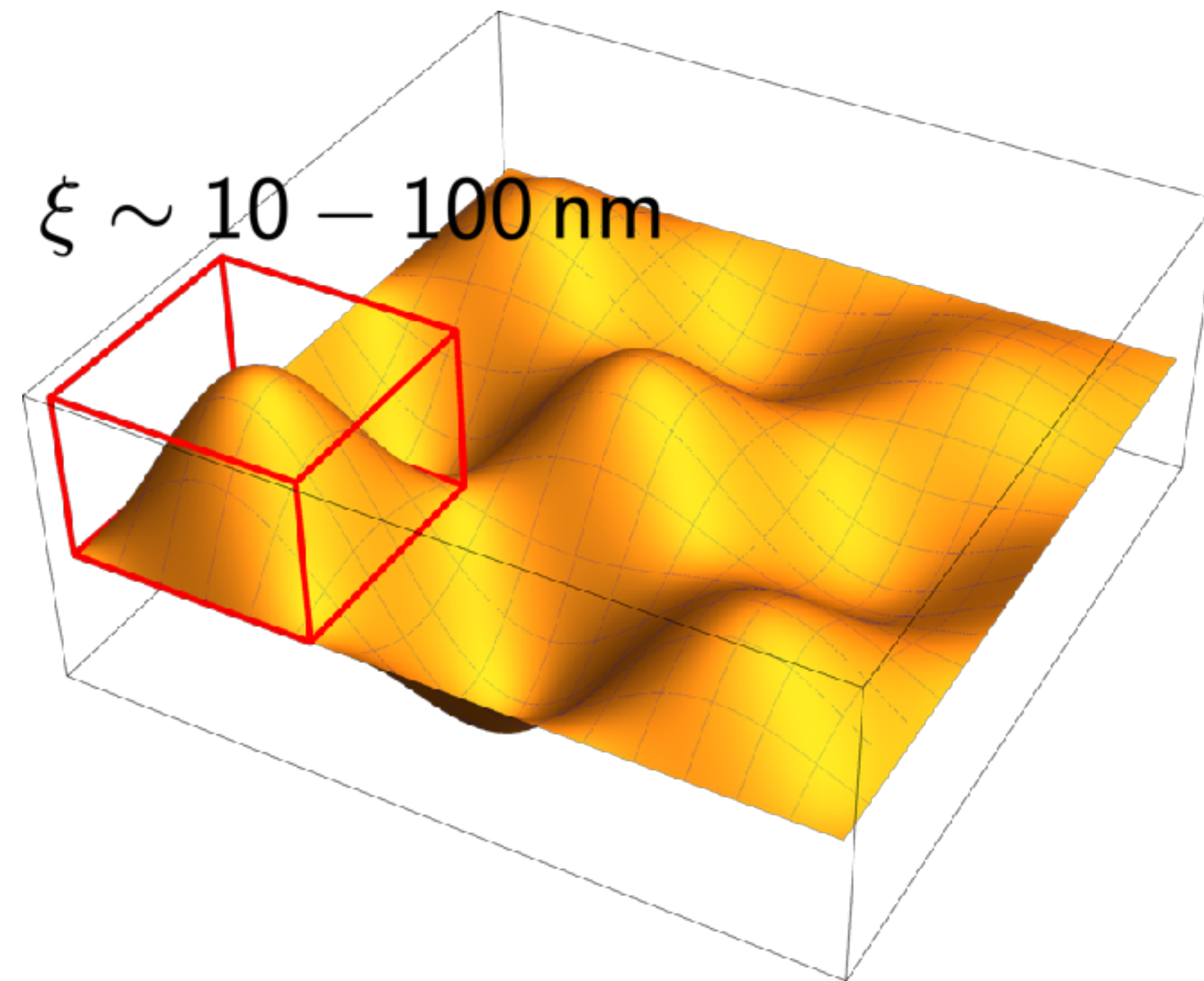
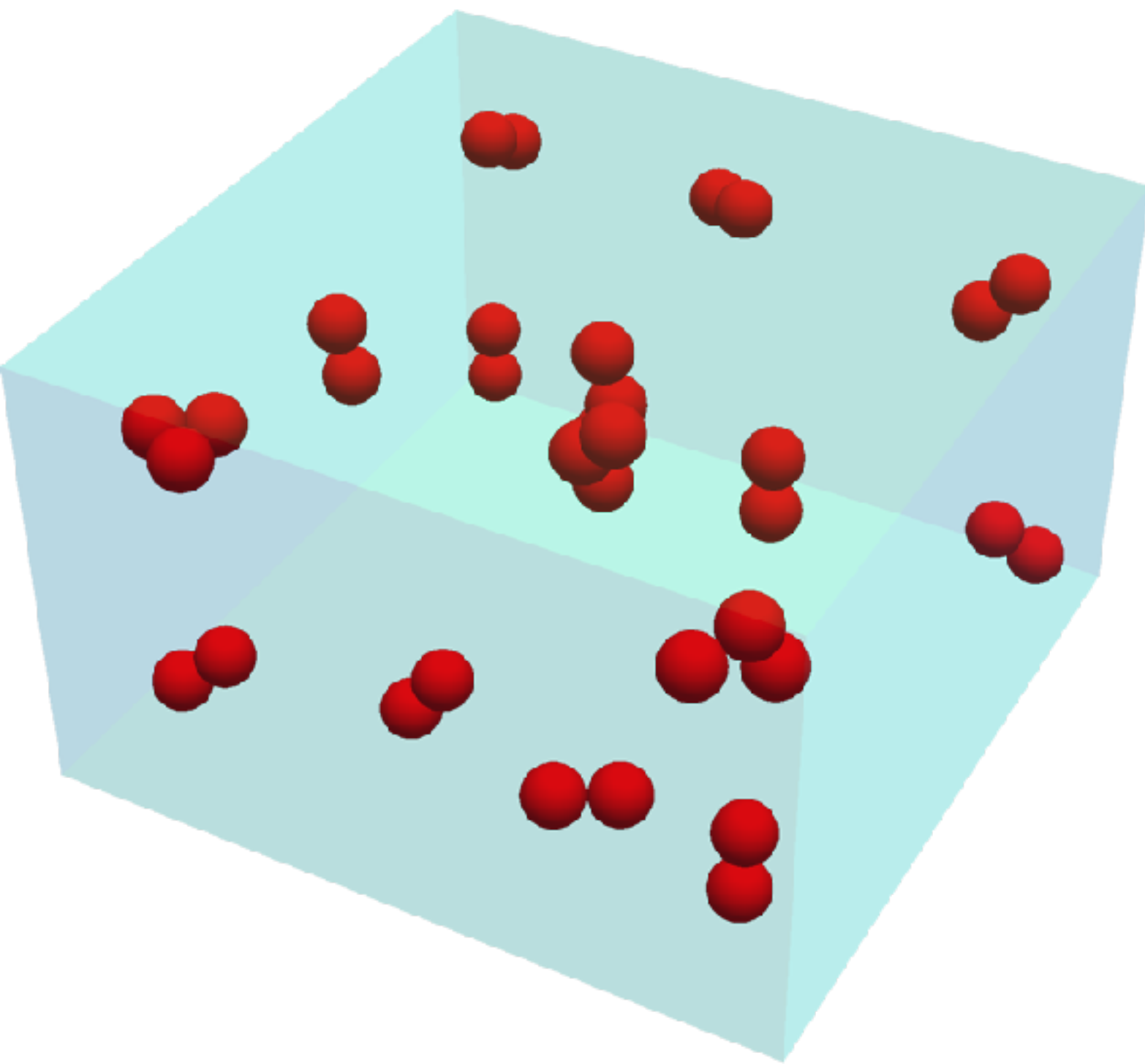
MORE RELEVANT

SUPERCONDUCTIVITY

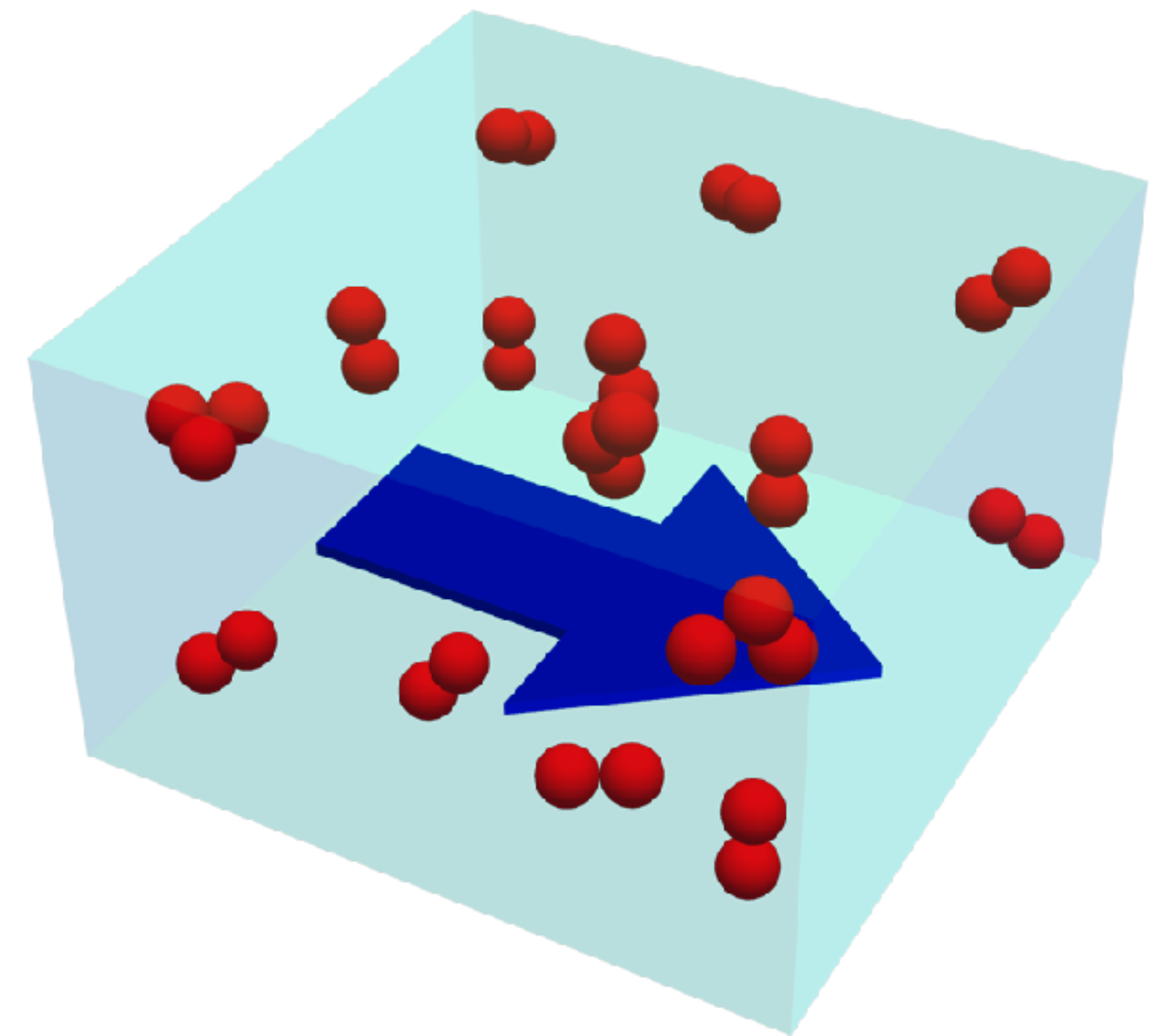
THEORY

THE GINZBURG-LANDAU THEORY

COMPLEX ORDER PARAMETER



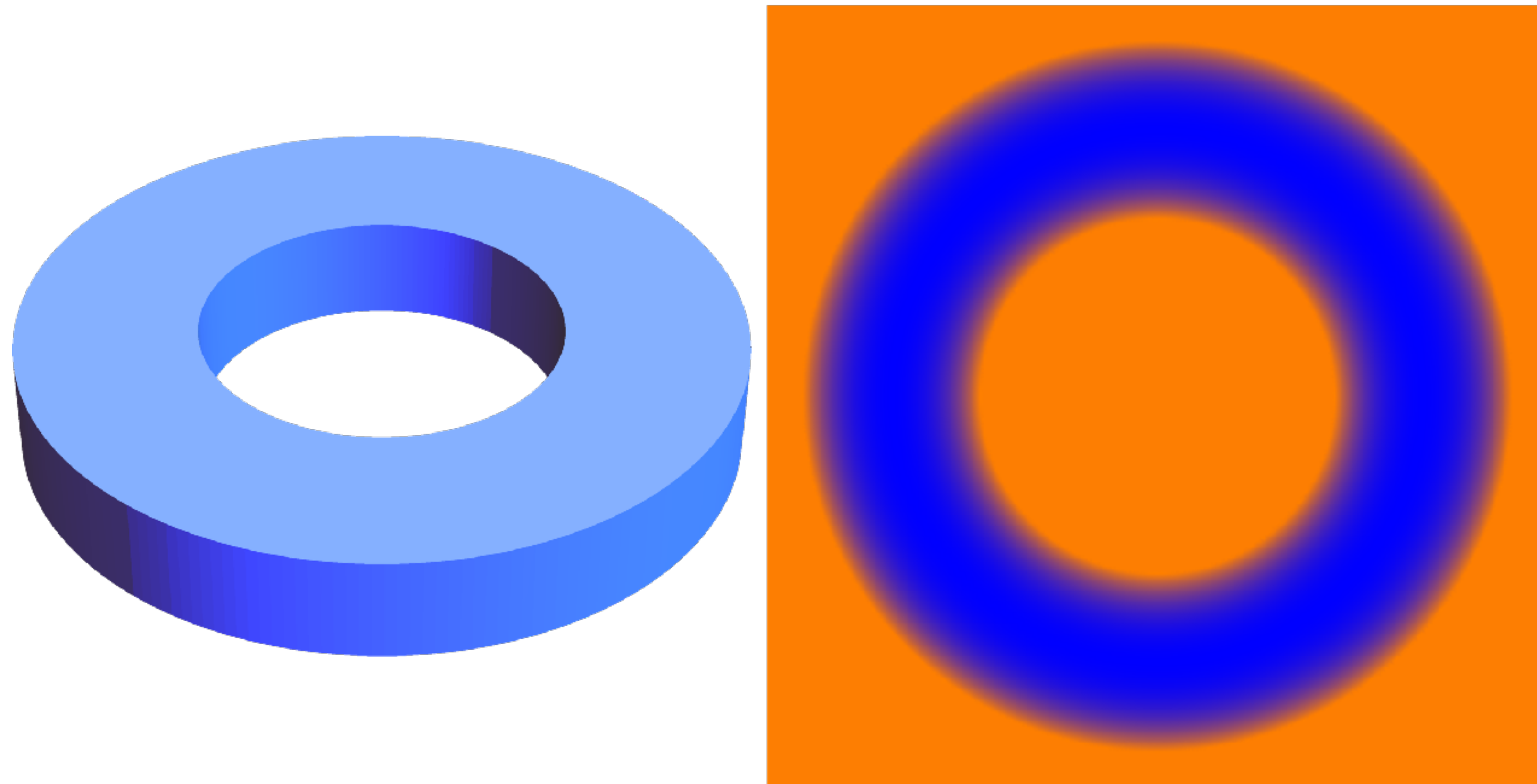
$$\psi(\mathbf{r}) = \sqrt{\rho(\mathbf{r})} \exp[i\phi(\mathbf{r})]$$



$$\mathbf{J}_S(\mathbf{r}) \propto \nabla \phi + 2\pi \mathbf{A} / \Phi_0$$

FLUX QUANTIZATION

(QUANTIZED SUPER-CURRENT)



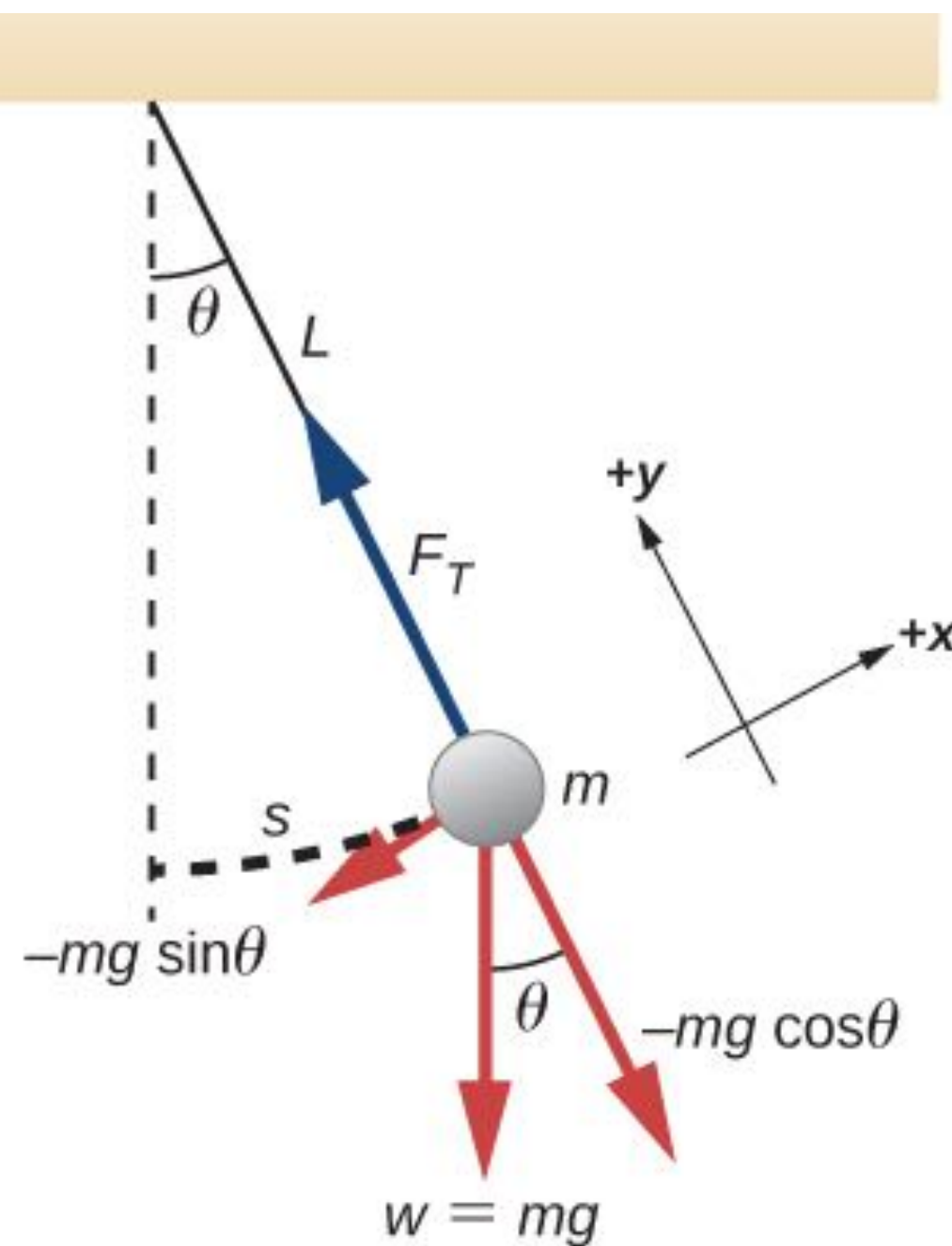
Single-valued-ness

$$\frac{\Phi}{\Phi_0} \in \mathbb{Z}$$

$$\Phi_0 := \frac{h}{2e}$$

SOMETHING VERY DIFFERENT

SIMPLE PENDULUM



$$\mathbf{F} = \frac{d\mathbf{P}}{dt} = m\mathbf{a} = m\ddot{\mathbf{r}}$$

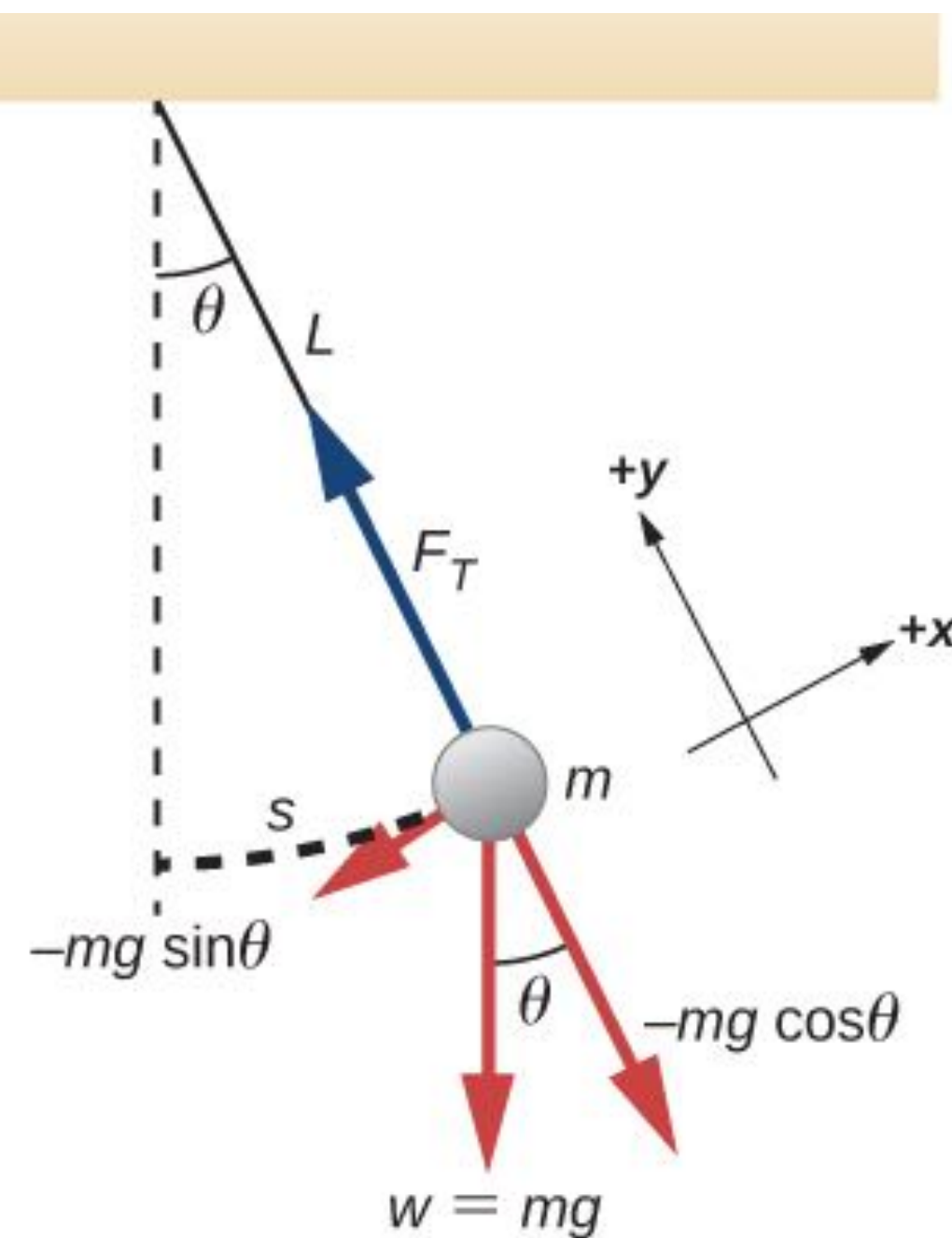
$$\boldsymbol{\tau} = \frac{d\mathbf{L}}{dt}$$

$$\boldsymbol{\tau} := \mathbf{r} \times \mathbf{F}$$

$$\mathbf{L} := \mathbf{r} \times \mathbf{P}$$

$$-mg \sin \theta = m\ell^2 \ddot{\theta}$$

SIMPLE PENDULUM

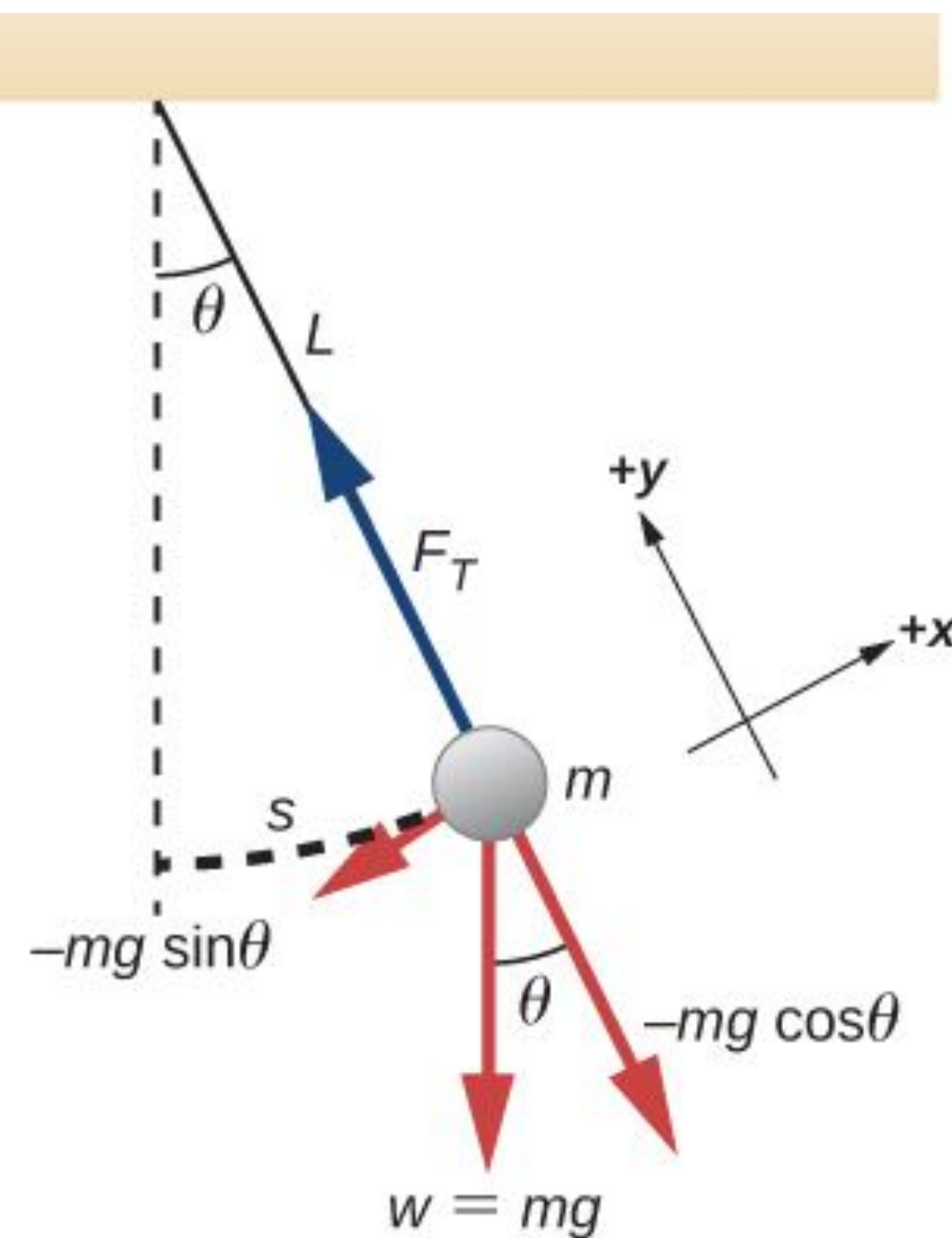


$$\mathcal{L} = T - V = \frac{1}{2}m\ell^2\dot{\theta}^2 - mg\ell(1 - \cos \theta)$$

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\theta}} - \frac{\partial \mathcal{L}}{\partial \theta} = 0$$

$$-mg \sin \theta = m\ell^2\ddot{\theta}$$

SIMPLE PENDULUM

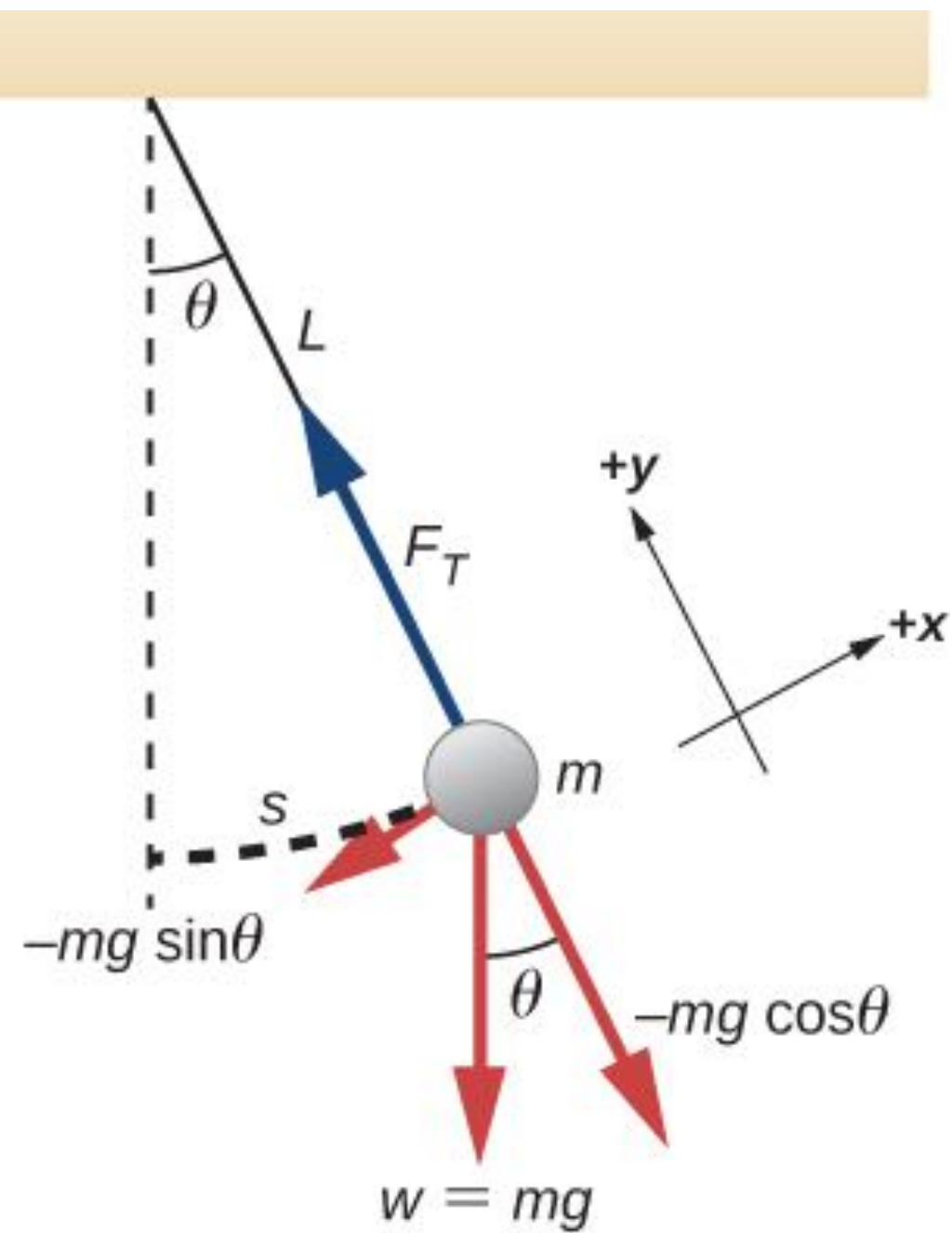


$$H = T + V = \frac{P_\theta^2}{2m\ell^2} + mgl(1 - \cos \theta)$$

$$H = P_\theta \dot{\theta} - \mathcal{L}, \quad P_\theta := \frac{\partial \mathcal{L}}{\partial \dot{\theta}}$$

$$\{\theta, P_\theta\} = 1$$

QUANTUM PENDULUM

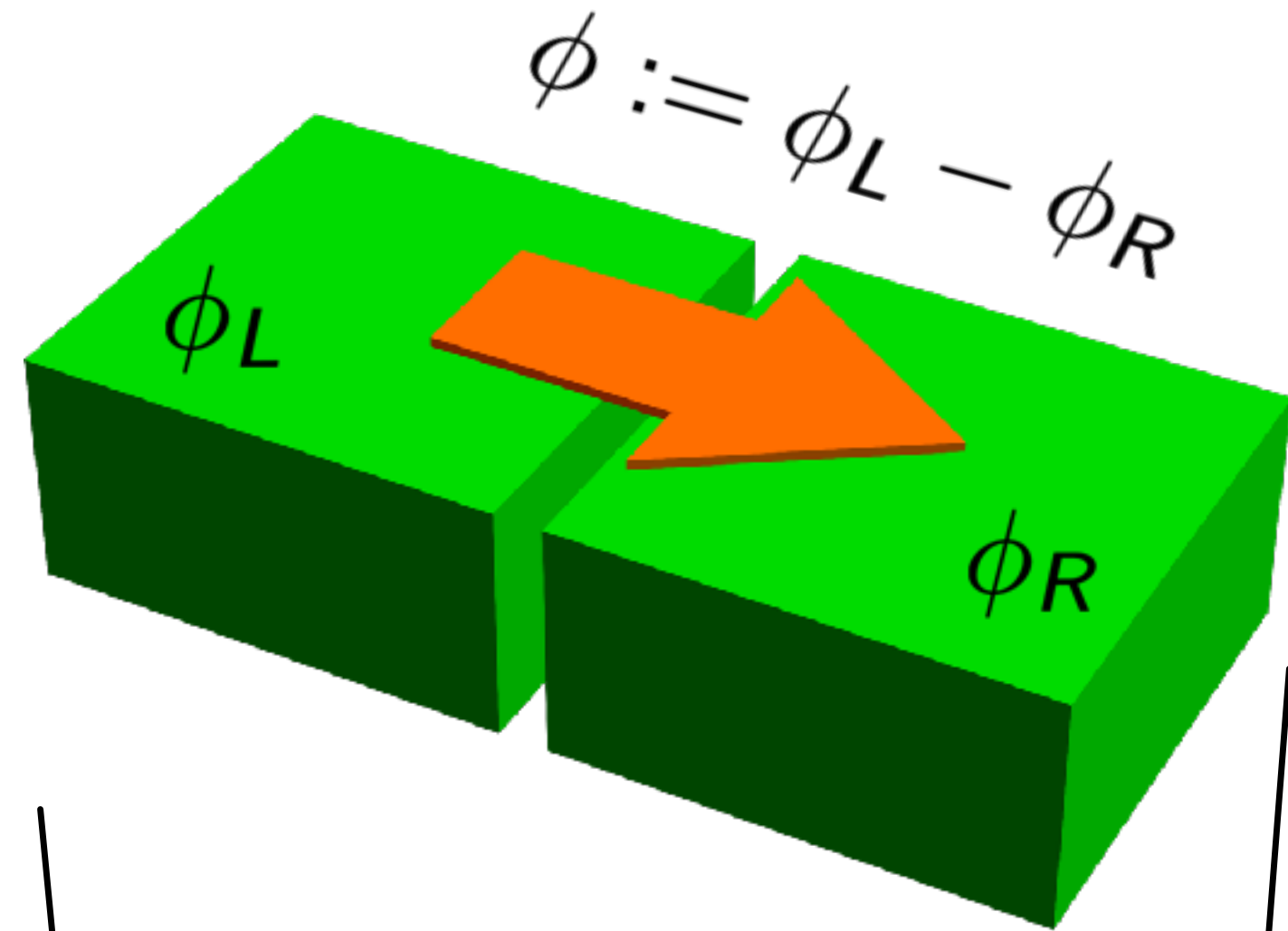


$$\hat{H} = \frac{\hat{P}_\theta^2}{2m\ell^2} - mgl \cos \hat{\theta}$$

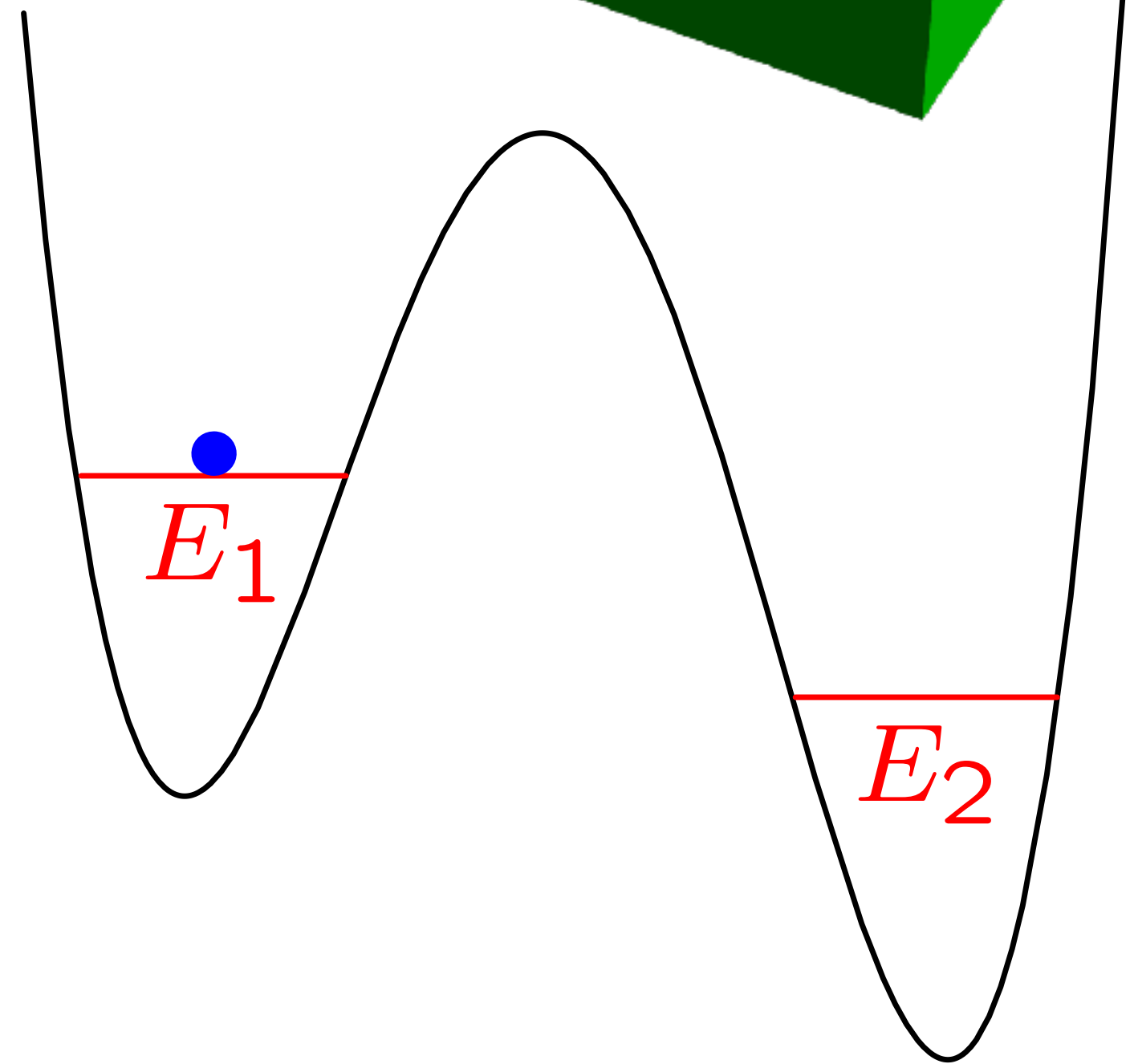
$$[\hat{\theta}, \hat{P}_\theta] = i\hbar$$

JOSEPHSON EFFECTS

JOSEPHSON JUNCTIONS



$$i\hbar\dot{\Psi}_L = E_L\Psi_L - J\Psi_R, \quad \Psi_L = \sqrt{N_L}e^{i\phi_L}$$
$$i\hbar\dot{\Psi}_R = E_R\Psi_R - J\Psi_L, \quad \Psi_R = \sqrt{N_R}e^{i\phi_R}$$



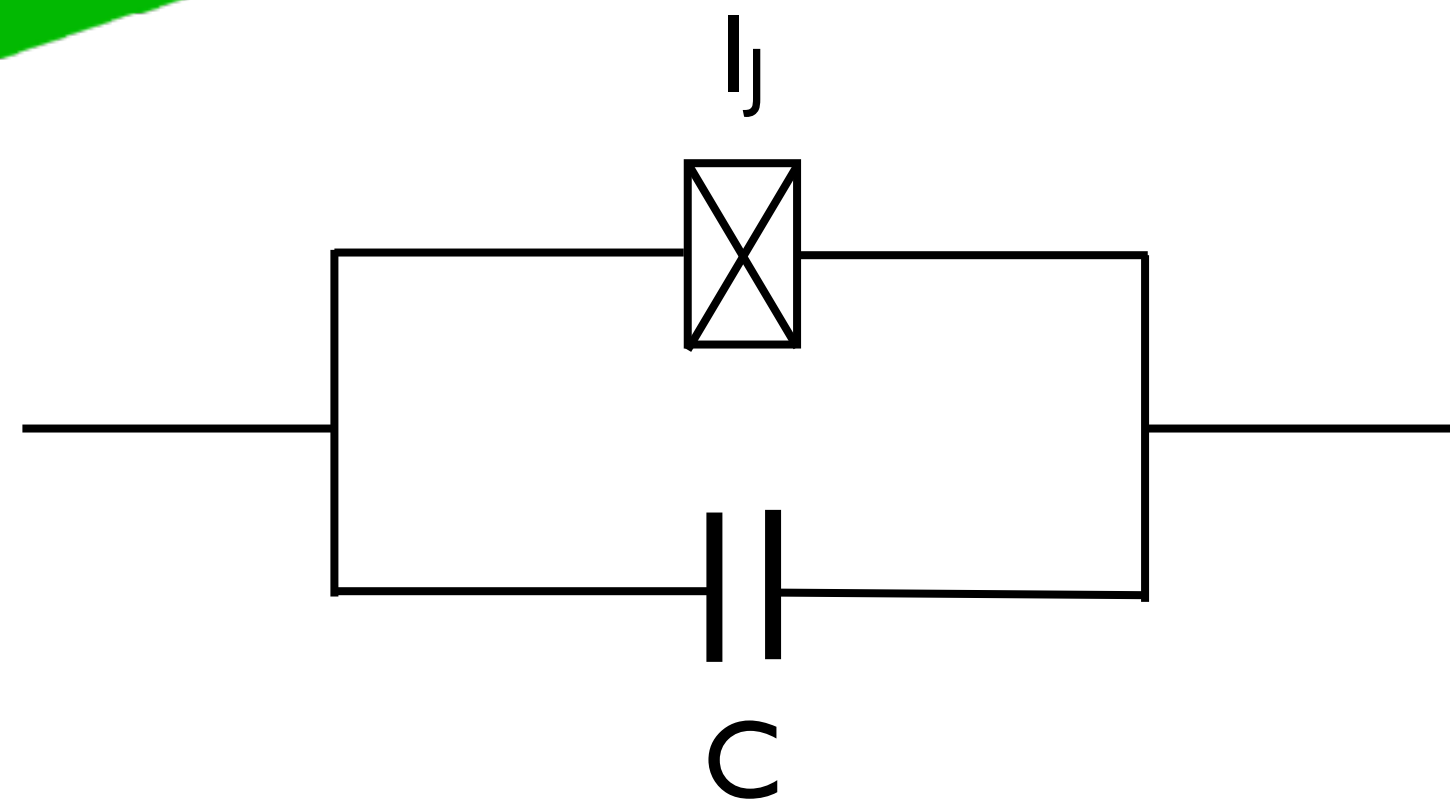
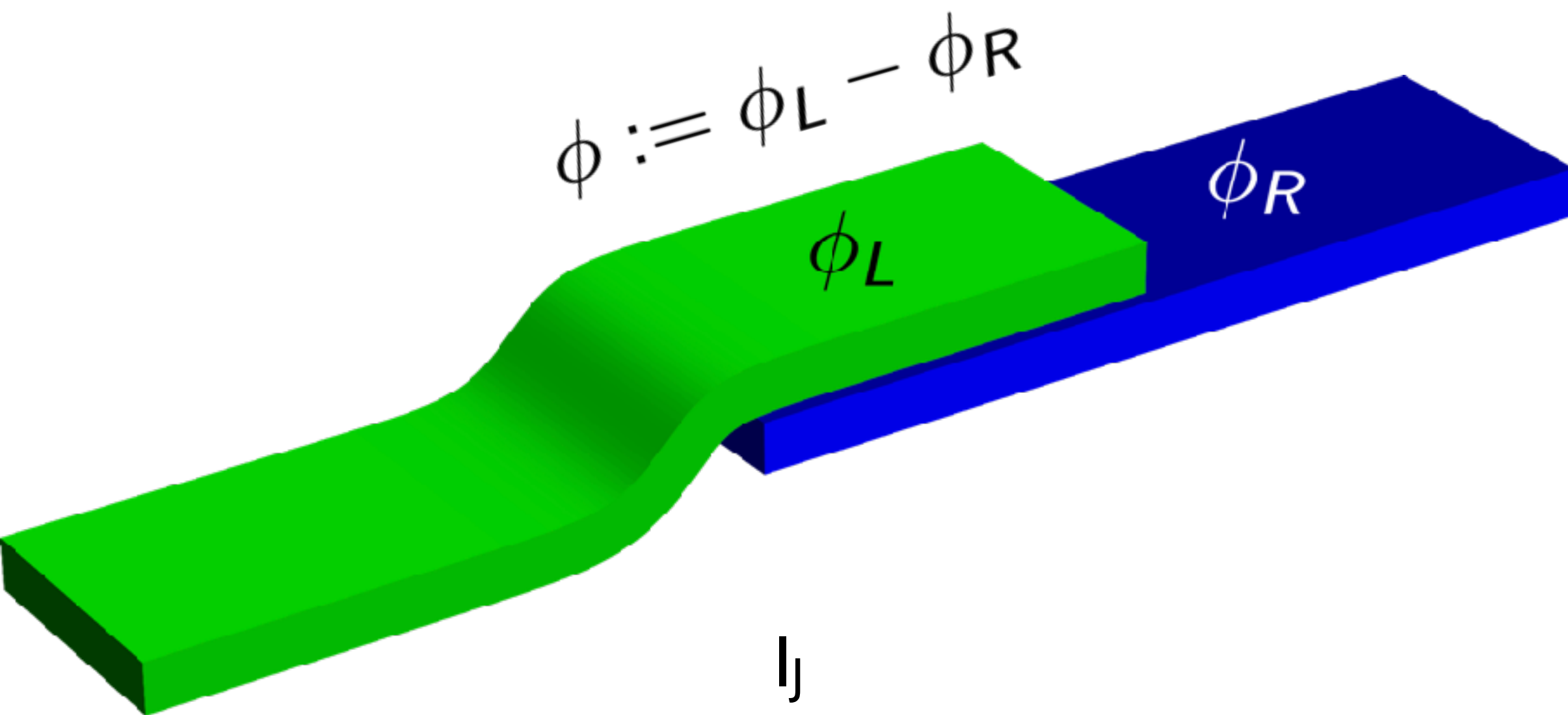
- DC Josephson Effect

$$I_S = \dot{N}_L = -\dot{N}_R = I_J \sin(\phi)$$

- AC Josephson Effect

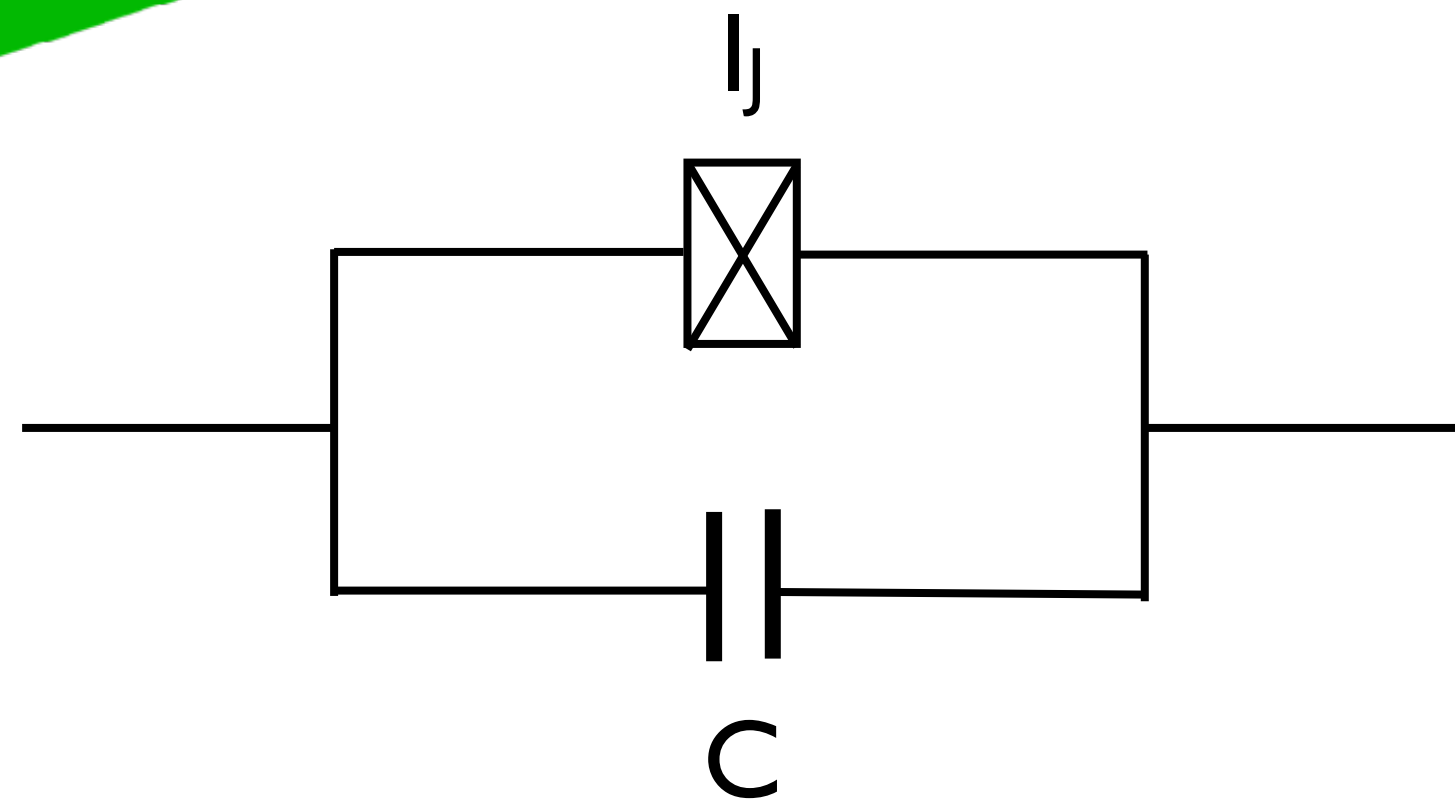
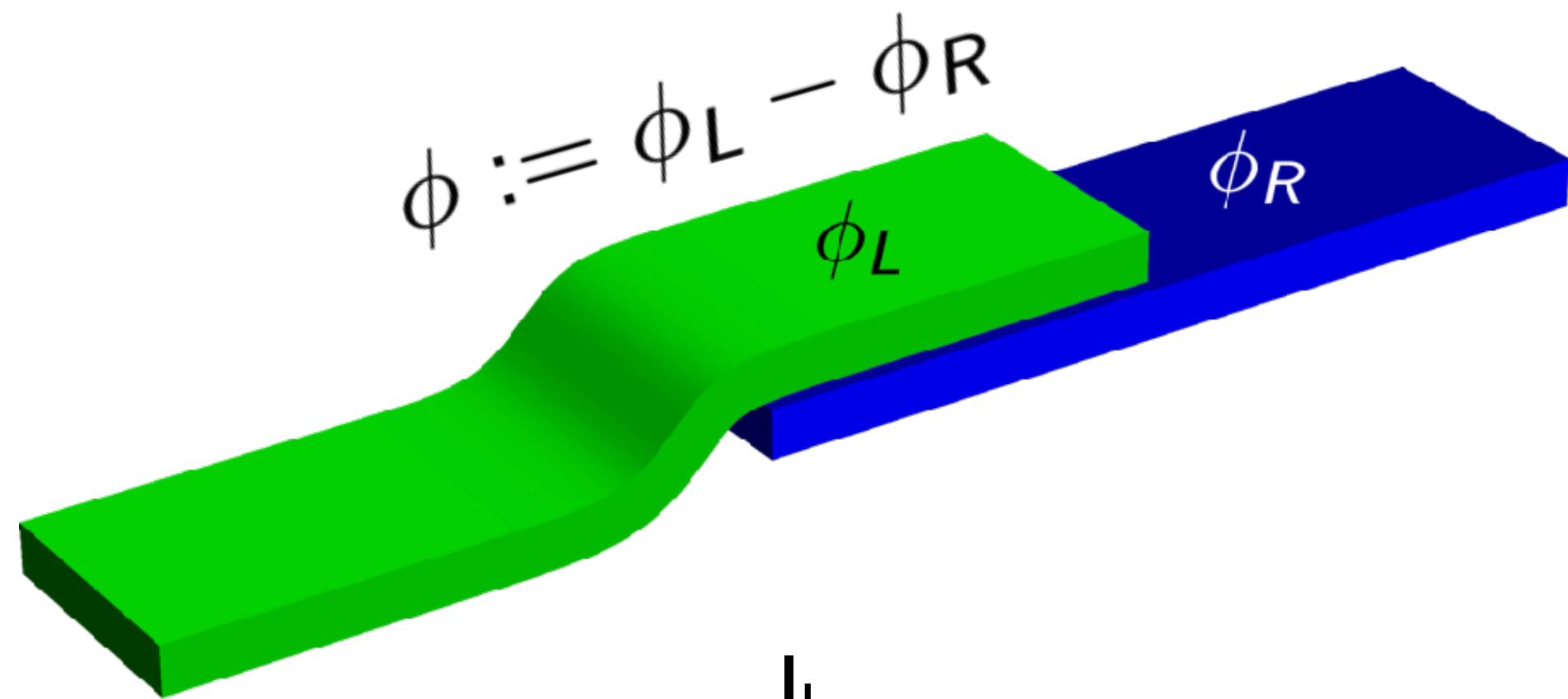
$$\hbar\dot{\phi} = E_L - E_R = 2eV$$

JOSEPHSON JUNCTION



- DC Josephson Effect
 $I_S = I_J \sin(\phi)$
- AC Josephson Effect
 $\dot{\phi} = 2eV/\hbar$

SMALL JOSEPHSON JUNCTION



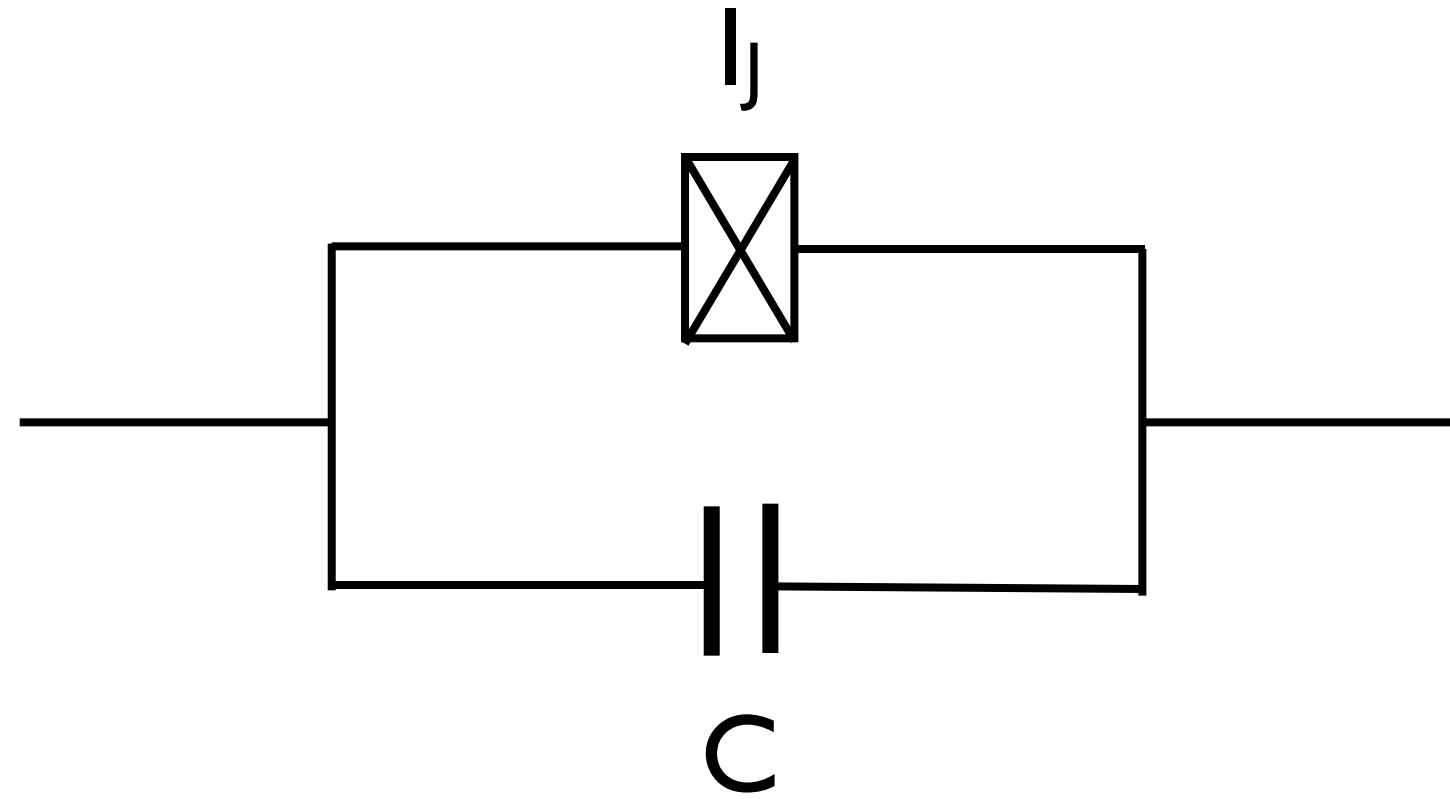
$$\dot{\phi} = \frac{2eV}{\hbar} = \frac{2e}{\hbar} \frac{Q}{C}$$

$$\dot{Q} = -I_J \sin(\phi)$$

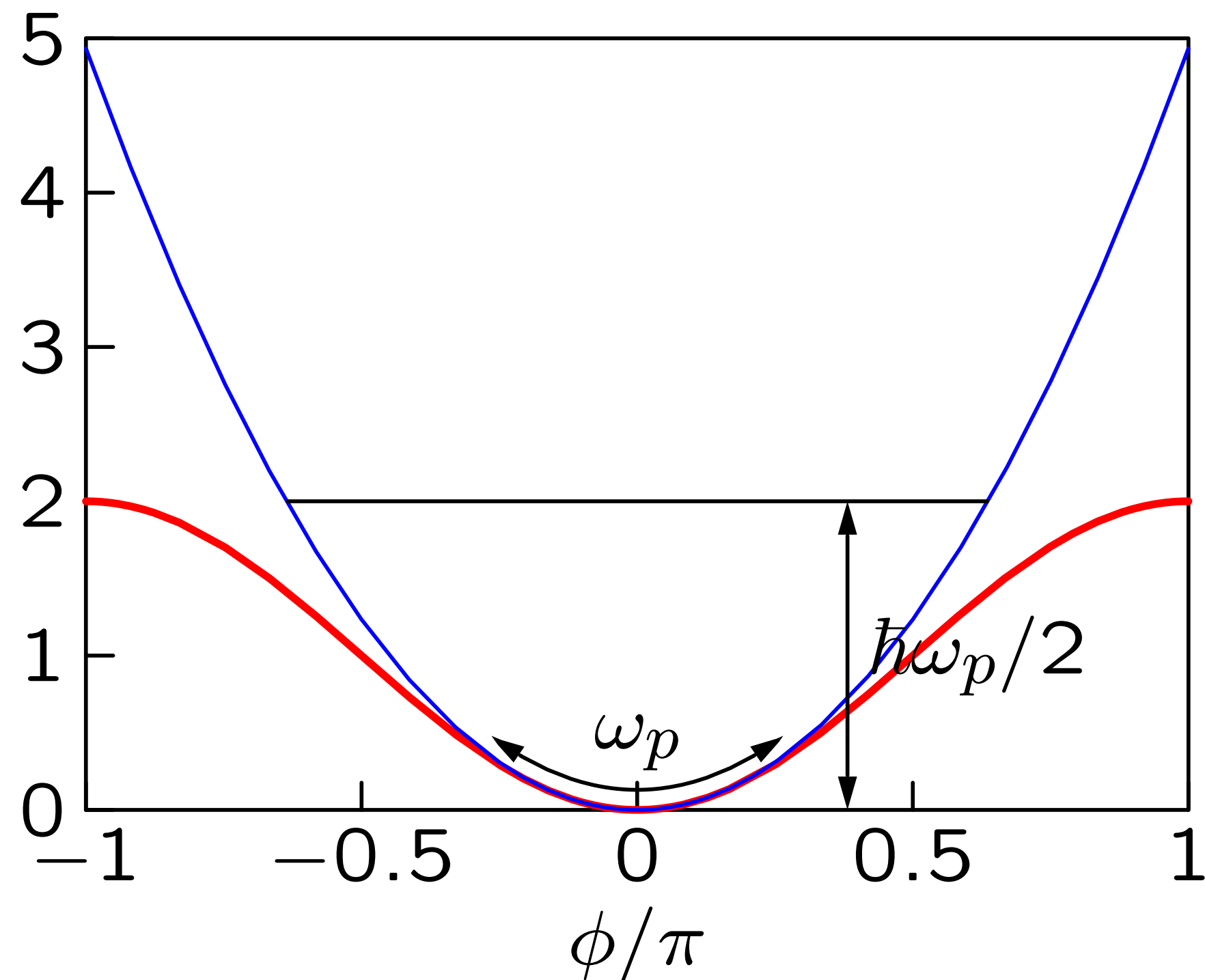
$$\ddot{\phi} = -\omega_p^2 \sin(\phi)$$

$$\omega_p^2 := \frac{2e I_J}{C}$$

MACROSCOPIC QUANTIZATION



$$\hat{H} = E_C \hat{n}^2 - E_J \cos \hat{\phi}, \quad [\hat{\phi}, \hat{n}] = i$$



$$E_C := \frac{(2e)^2}{2C}$$

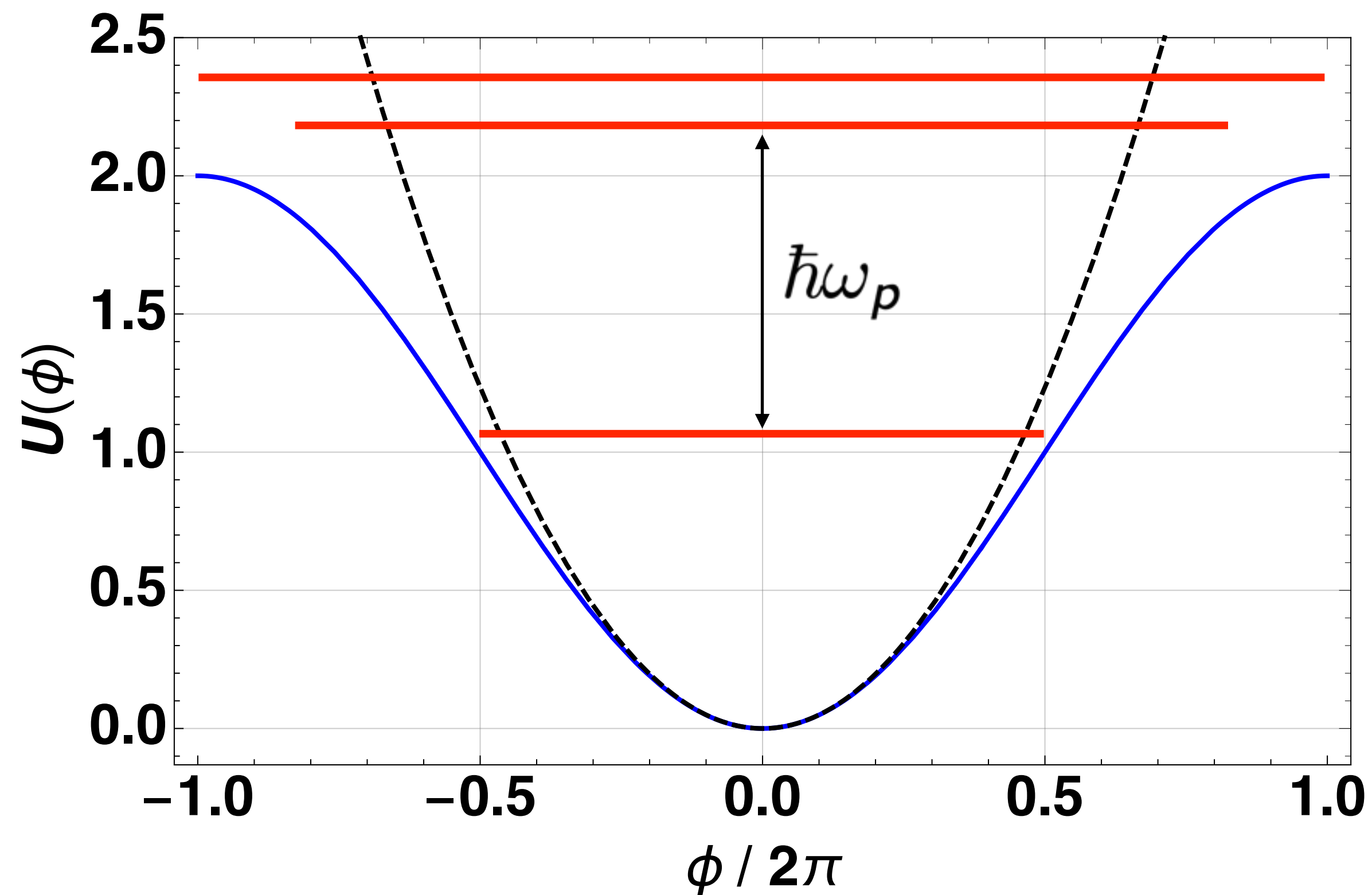
$$E_J := \frac{\hbar I_J}{2e} = \frac{I_J \Phi_0}{2\pi}$$

$$\hbar\omega_p := \sqrt{2E_C E_J}$$

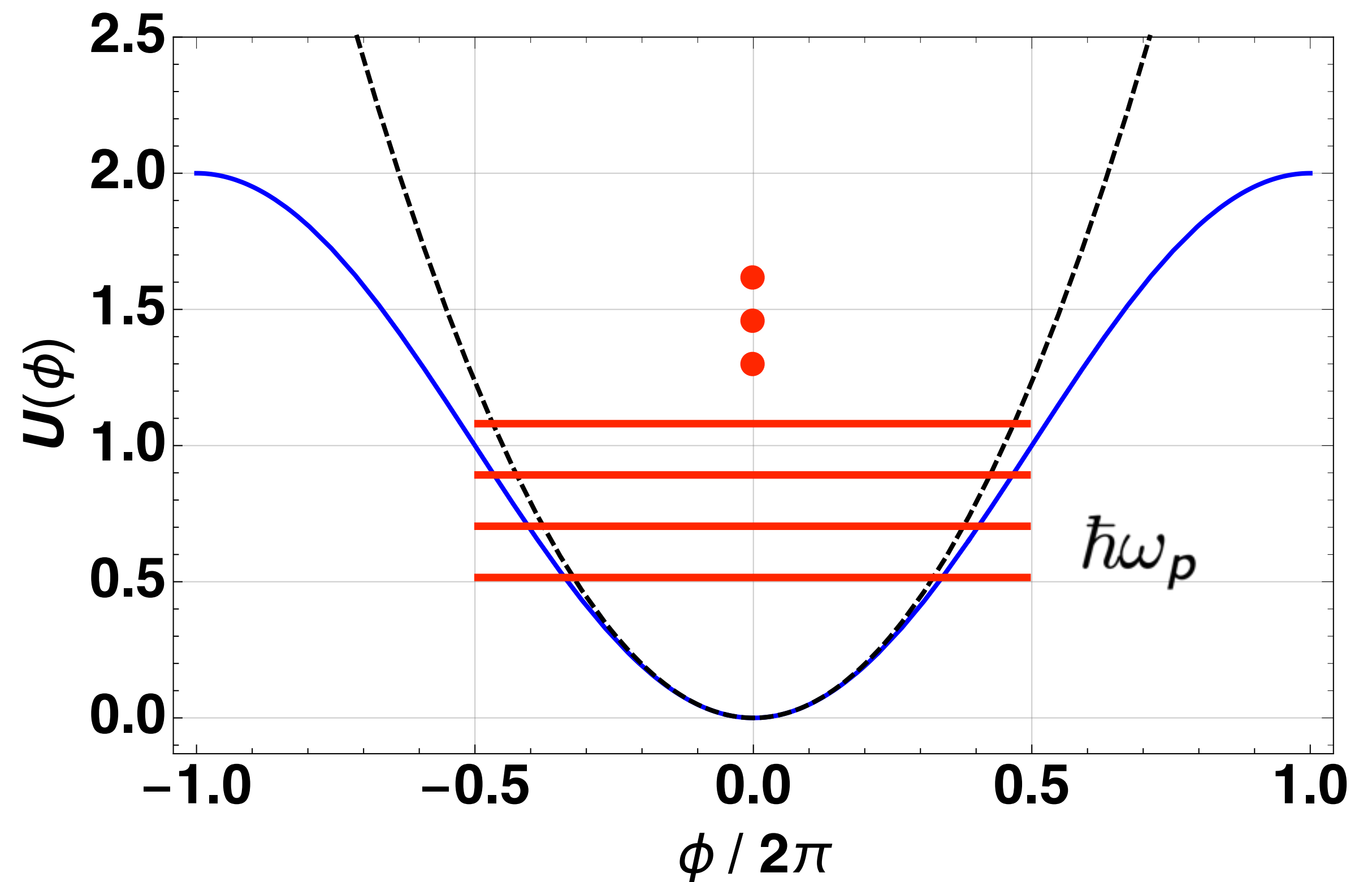
TWO EXTREMES

$$\hat{H} = E_C \hat{n}^2 - E_J \cos \hat{\phi}, \quad [\hat{\phi}, \hat{n}] = i$$

$$E_C \gg E_J$$



$$E_C \ll E_J$$

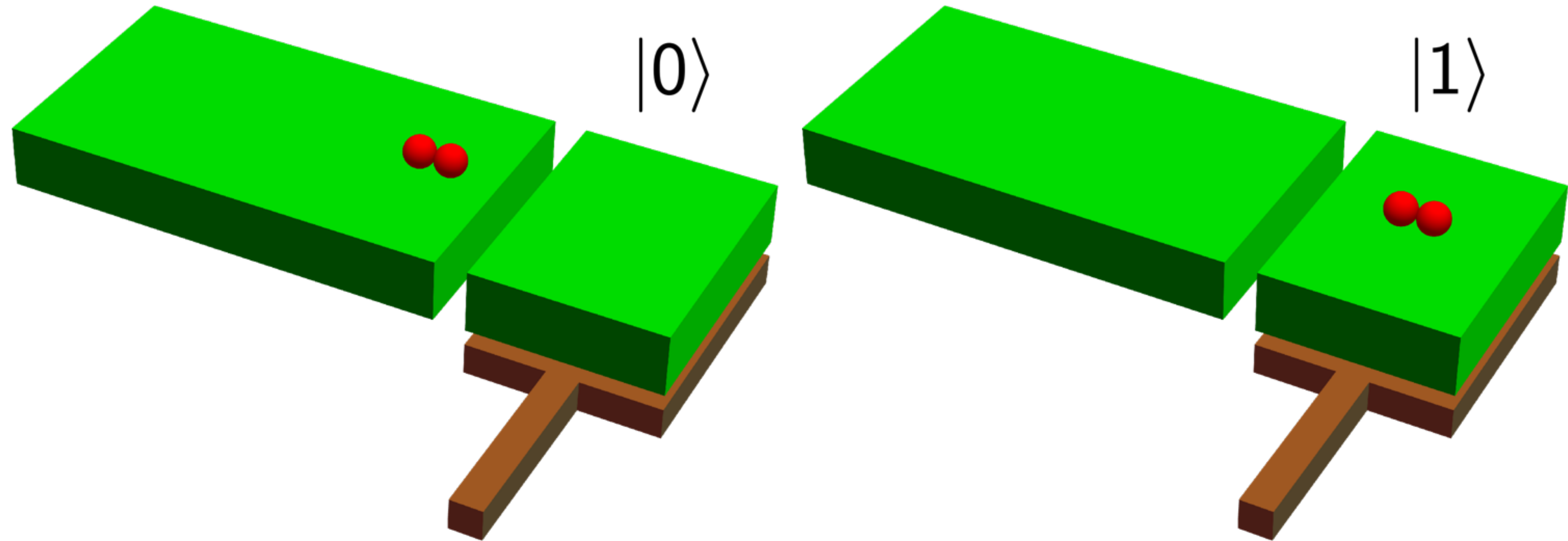


SUPERCONDUCTING QUBITS

CHARGE, PHASE & TRANSMON QUBITS

QUANTIZED CHARGE

(QUANTIZED NUMBER OF COOPER PAIRS)

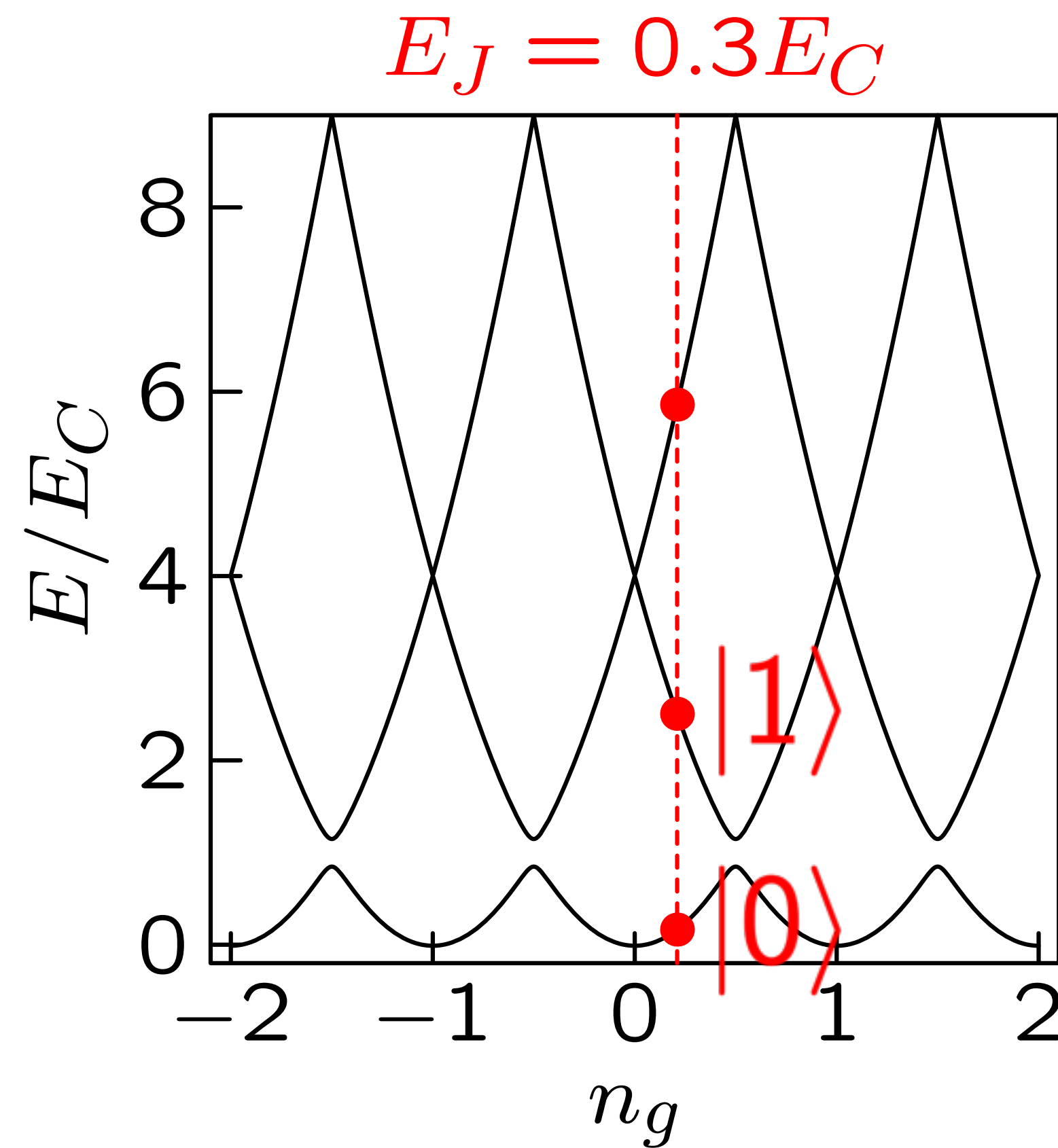


$$H_{\text{qubit}} = \frac{1}{2} \Omega \sigma^z$$

$$\Omega/2\pi \sim 5 \text{ GHz}$$

Bouchiat *et al.* (Phys. Scr., 1998)
Nakamura *et al.* (Nature, 1999)

QUANTIZED CHARGE



$$[\phi, n] = i$$

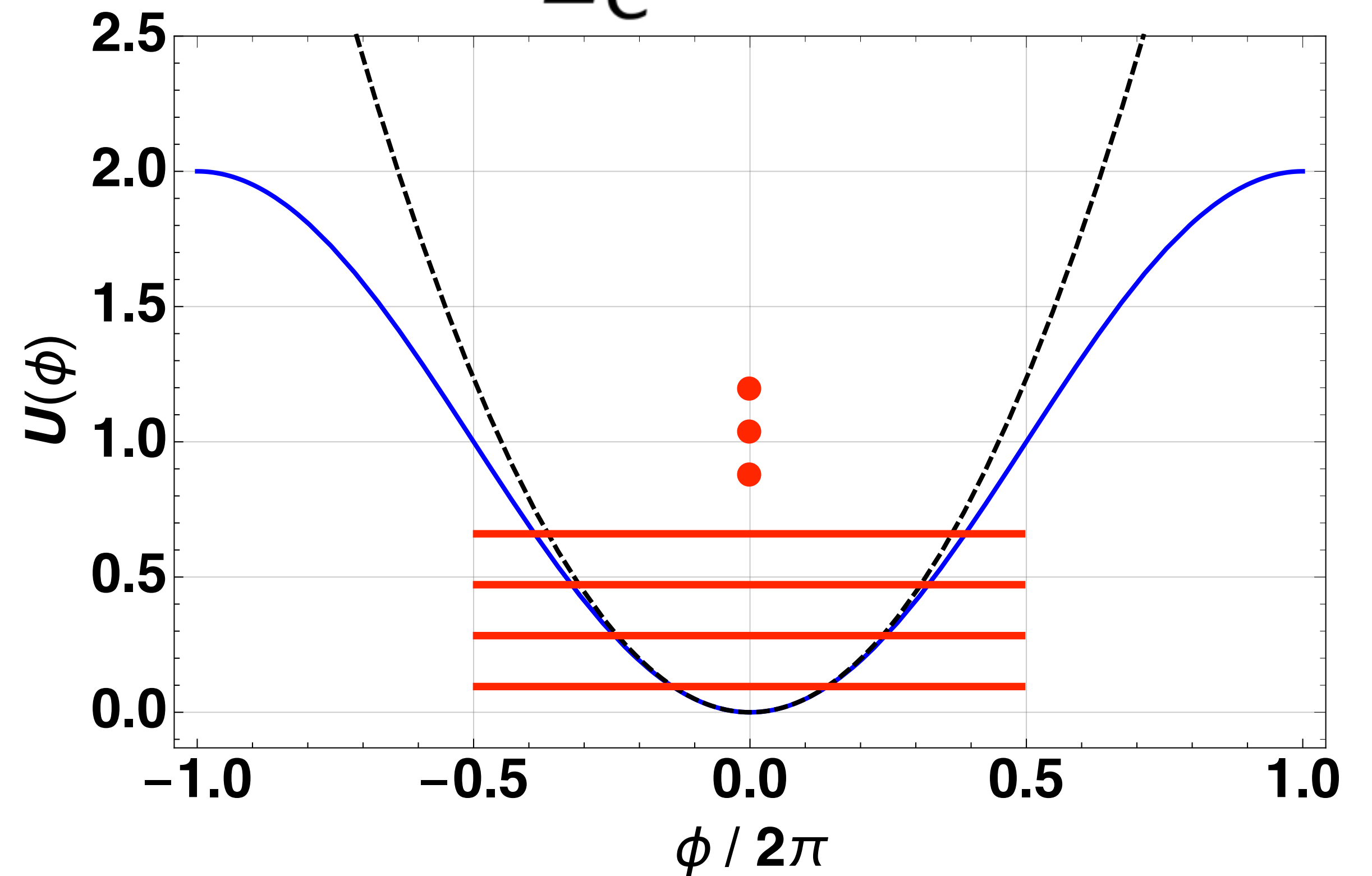
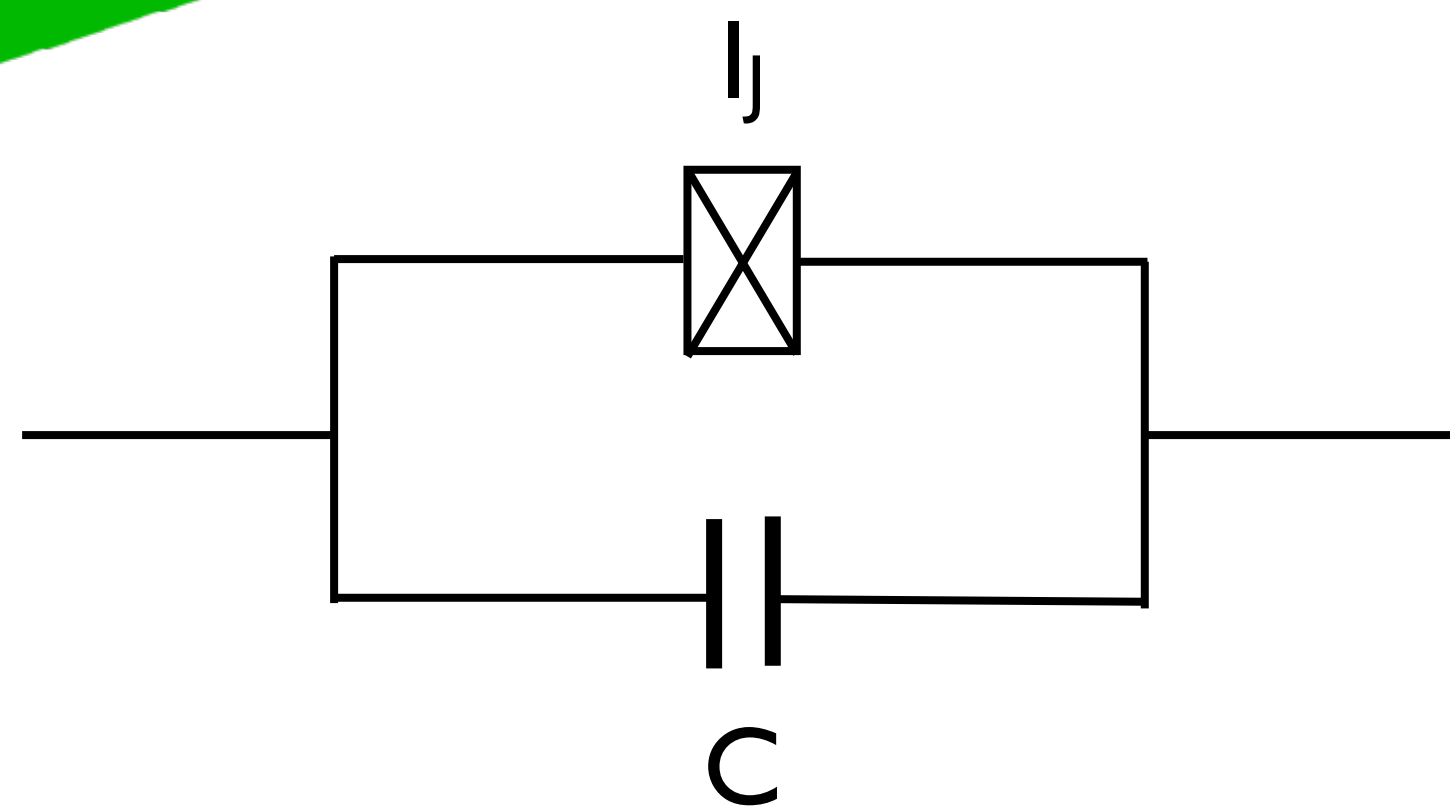
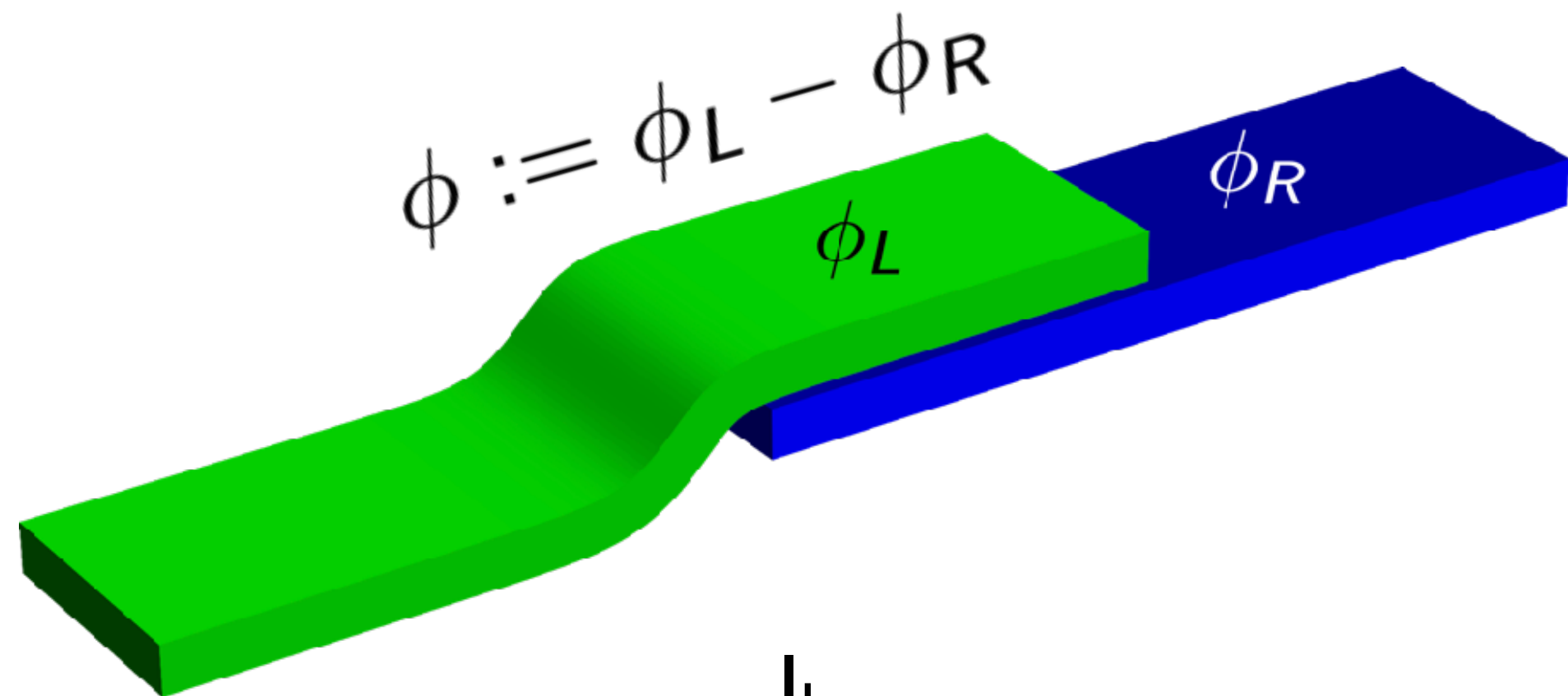
$$H = \overbrace{E_C(n - n_g)^2}^{\text{dominant}} - E_J \cos \phi$$

$$H \approx E_C(n_g - 1)\sigma^z - \frac{1}{2}E_J\sigma^x$$

PHASE QUBIT

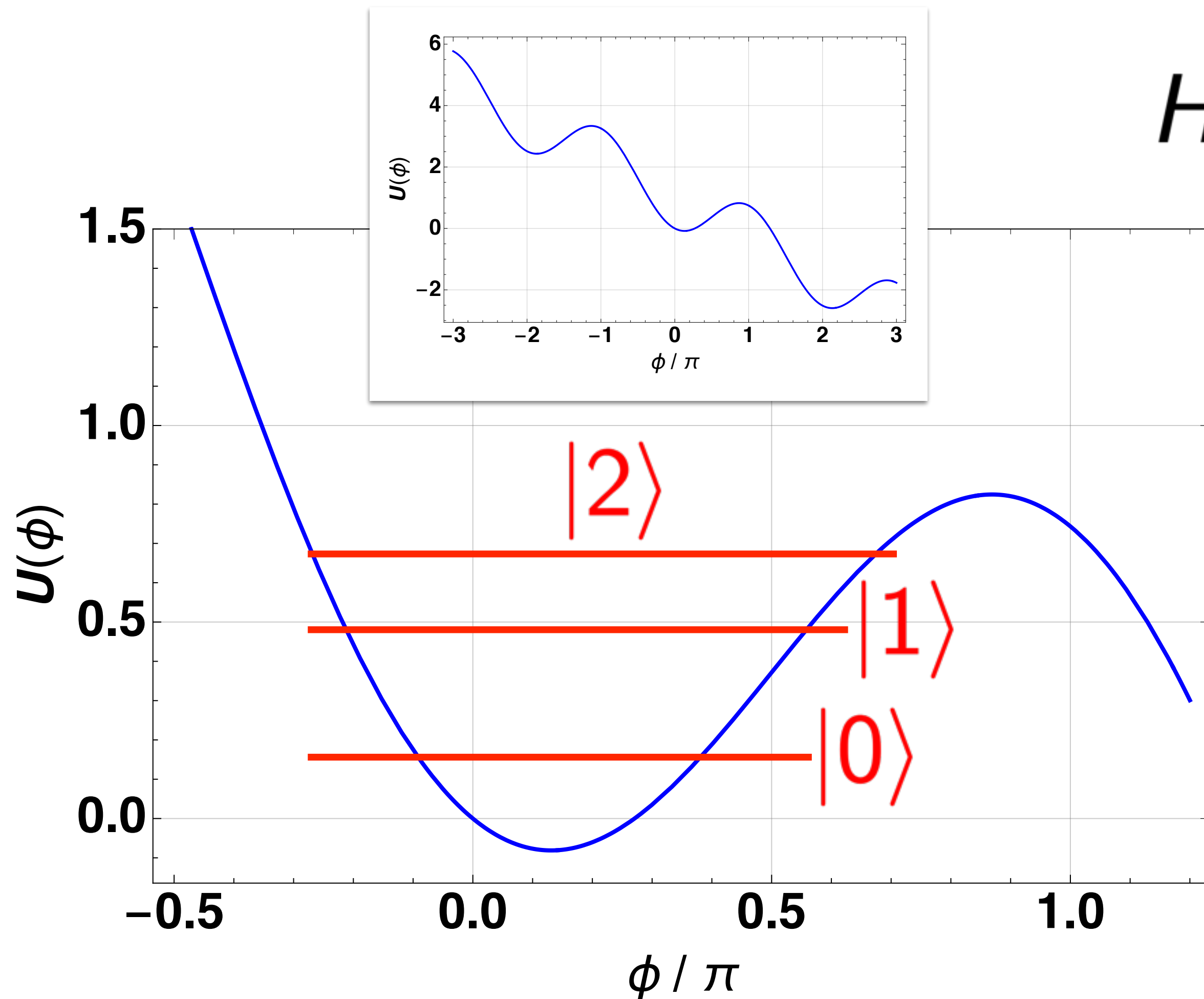
BIASED ANHARMONIC OSCILLATOR

$$\frac{E_J}{E_C} \sim 10^4$$



PHASE QUBIT

BIASED ANHARMONIC OSCILLATOR



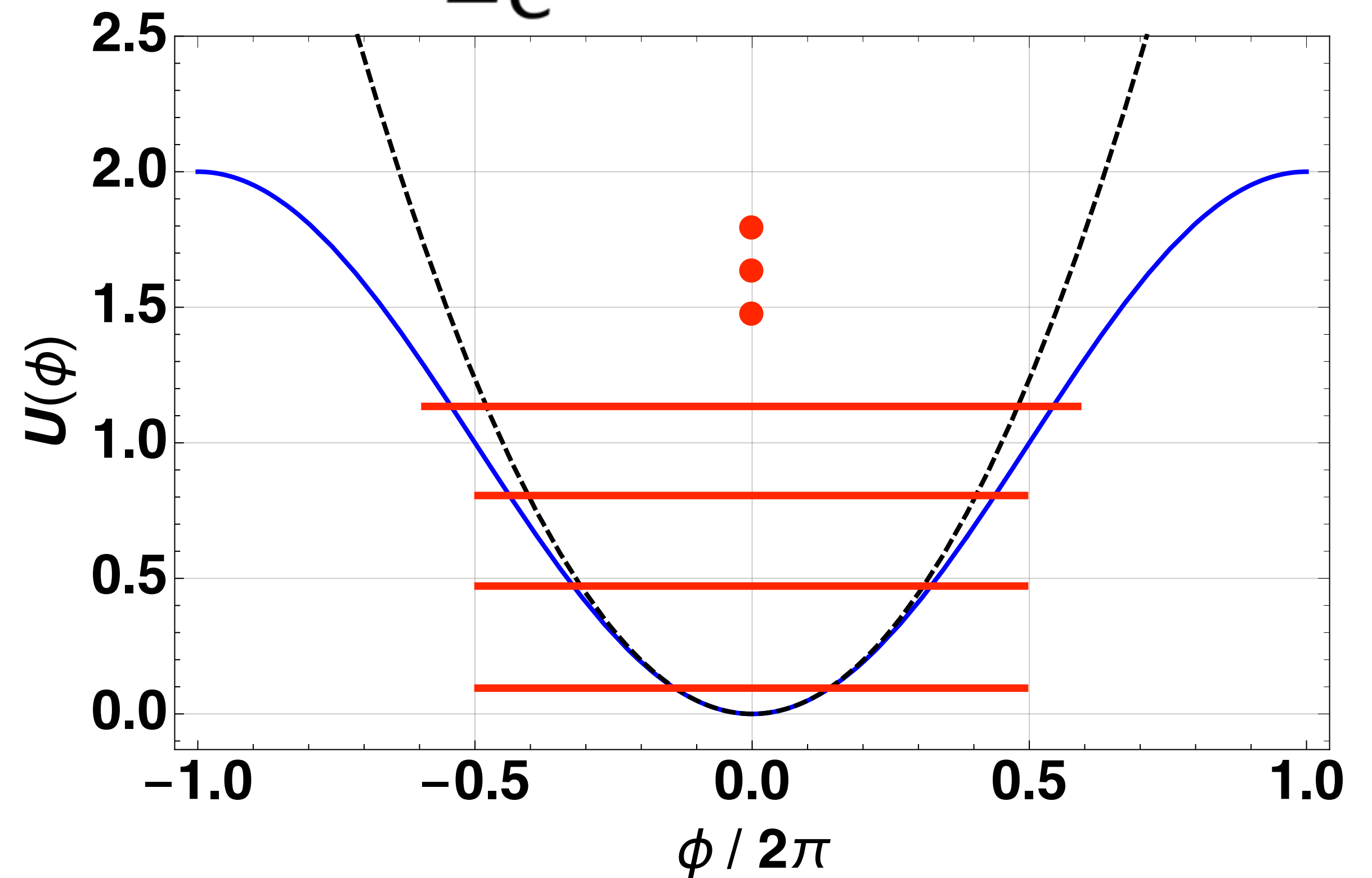
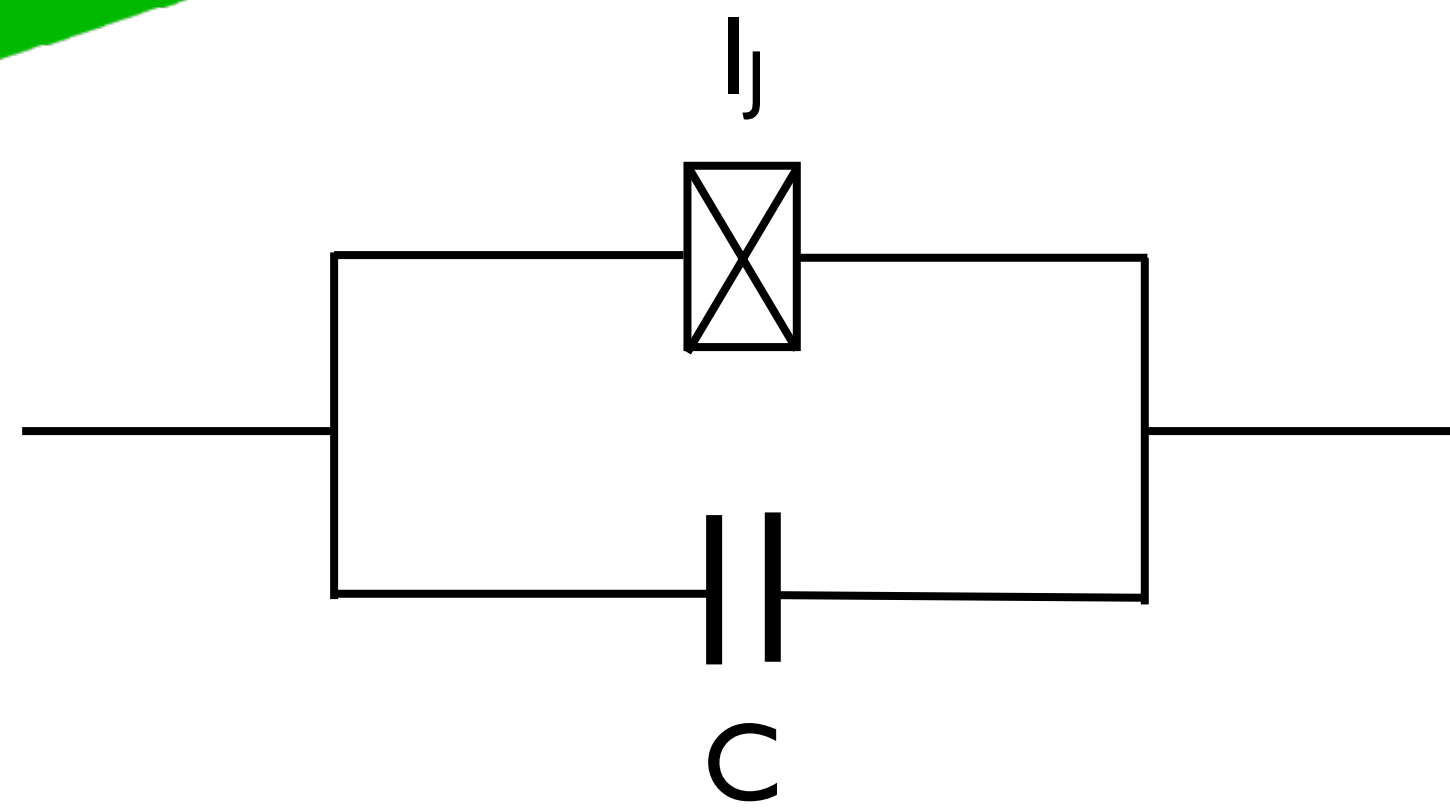
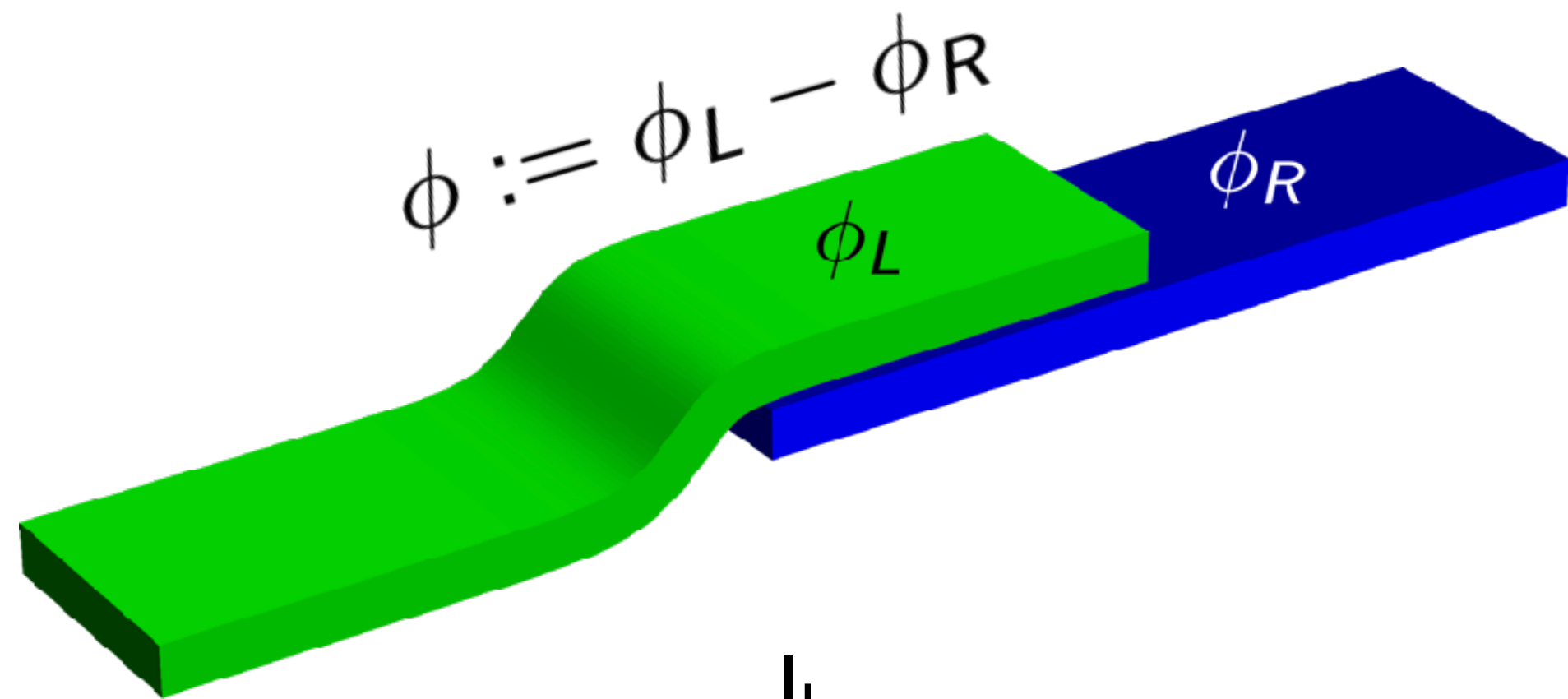
$$H = \underbrace{E_C(n - n_g)^2}_{\text{perturbation}} - E_J \cos \phi - I \phi$$

Martinis *et al.* (PRL, 2002)
Steffen *et al.* (PRL, 2006)

QUANTRONIUM / TRANSMON QUBIT

(UNBIASED) ANHARMONIC OSCILLATOR

$$\frac{E_J}{E_C} \simeq 1 - 100$$



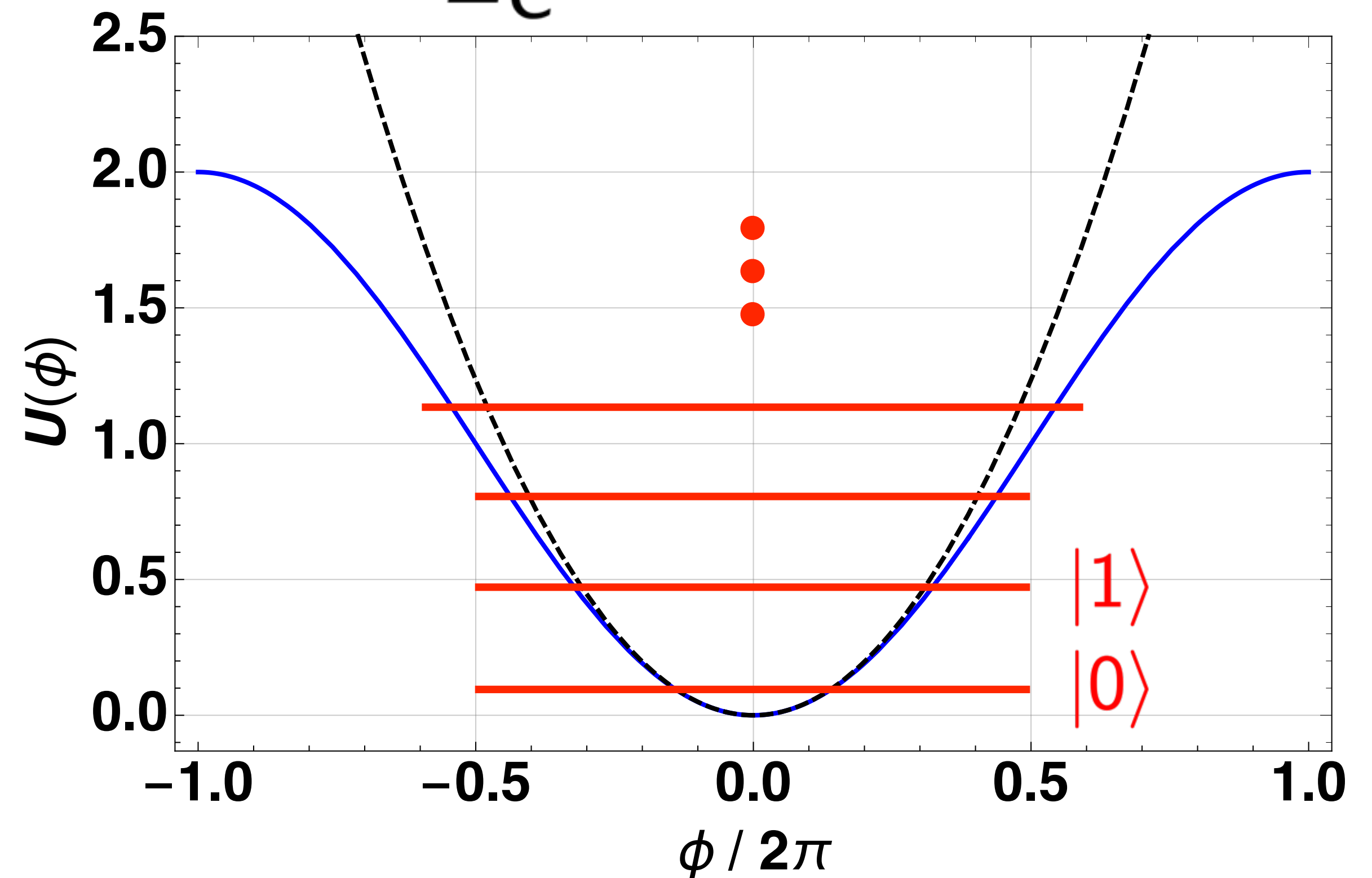
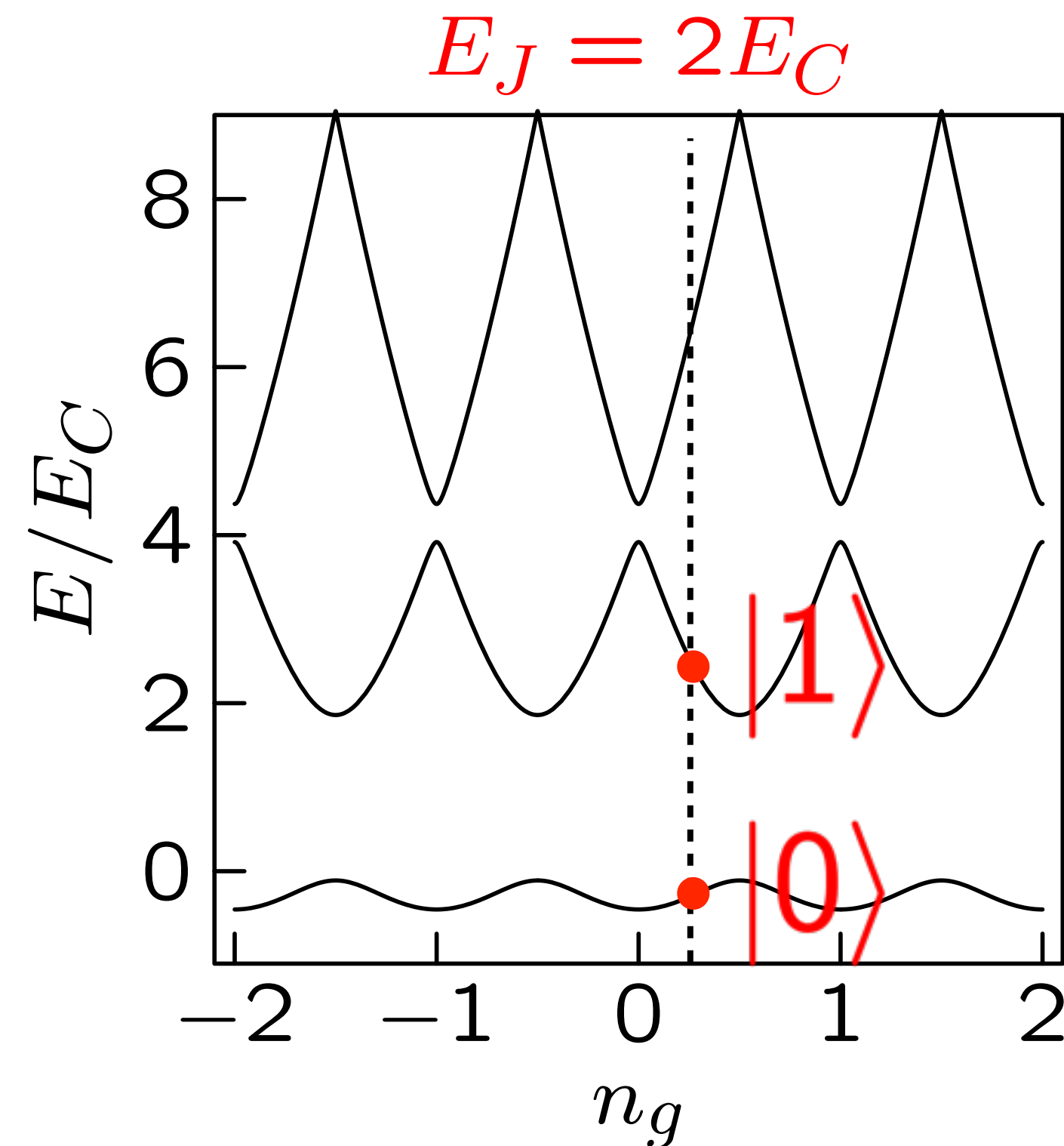
QUANTRONIUM / TRANSMON QUBIT

(UNBIASED) ANHARMONIC OSCILLATOR

Cottet (PhD, 2002)

Vion *et al.* (Science, 2002)

$$\frac{E_J}{E_C} \simeq 1 - 100$$



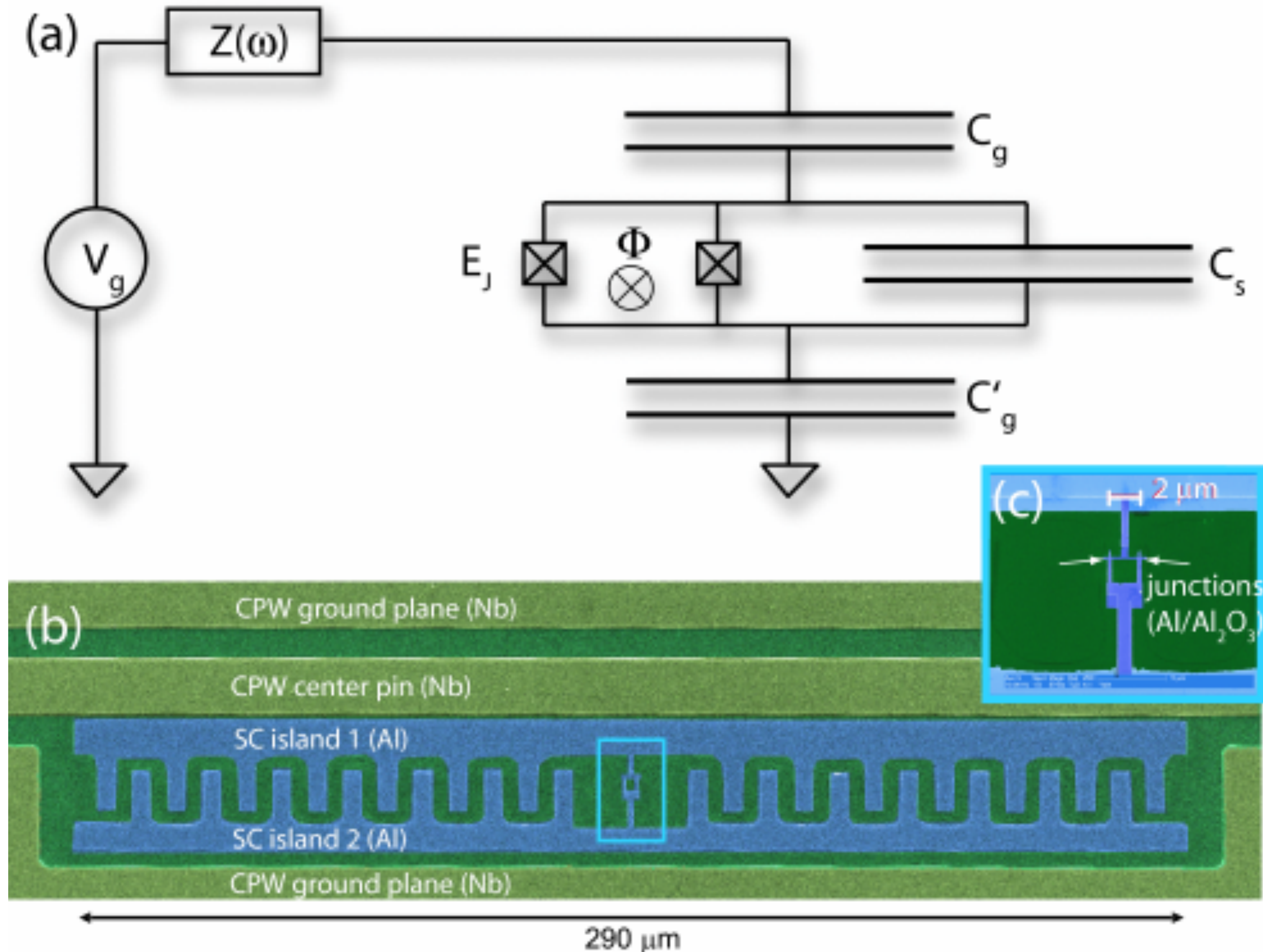
TRANSMON QUBIT (ATOM)

Koch et al. (PRA, 2007)

Houck et al. (Nature, 2007)

$$\Omega/2\pi \sim 5 \text{ GHz}$$

$$H_{\text{qubit}} = \frac{1}{2} \Omega \sigma^z$$

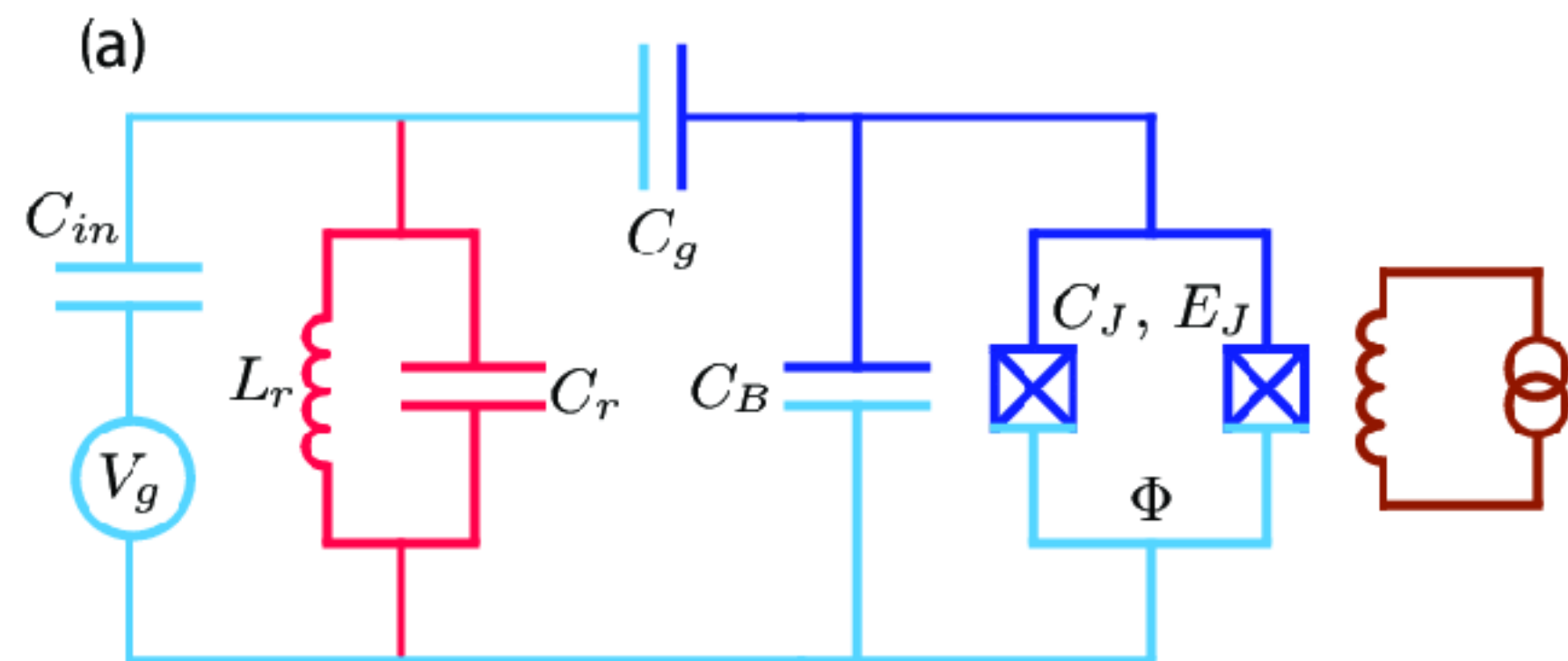


TRANSMON QUBIT

ANHARMONIC OSCILLATOR

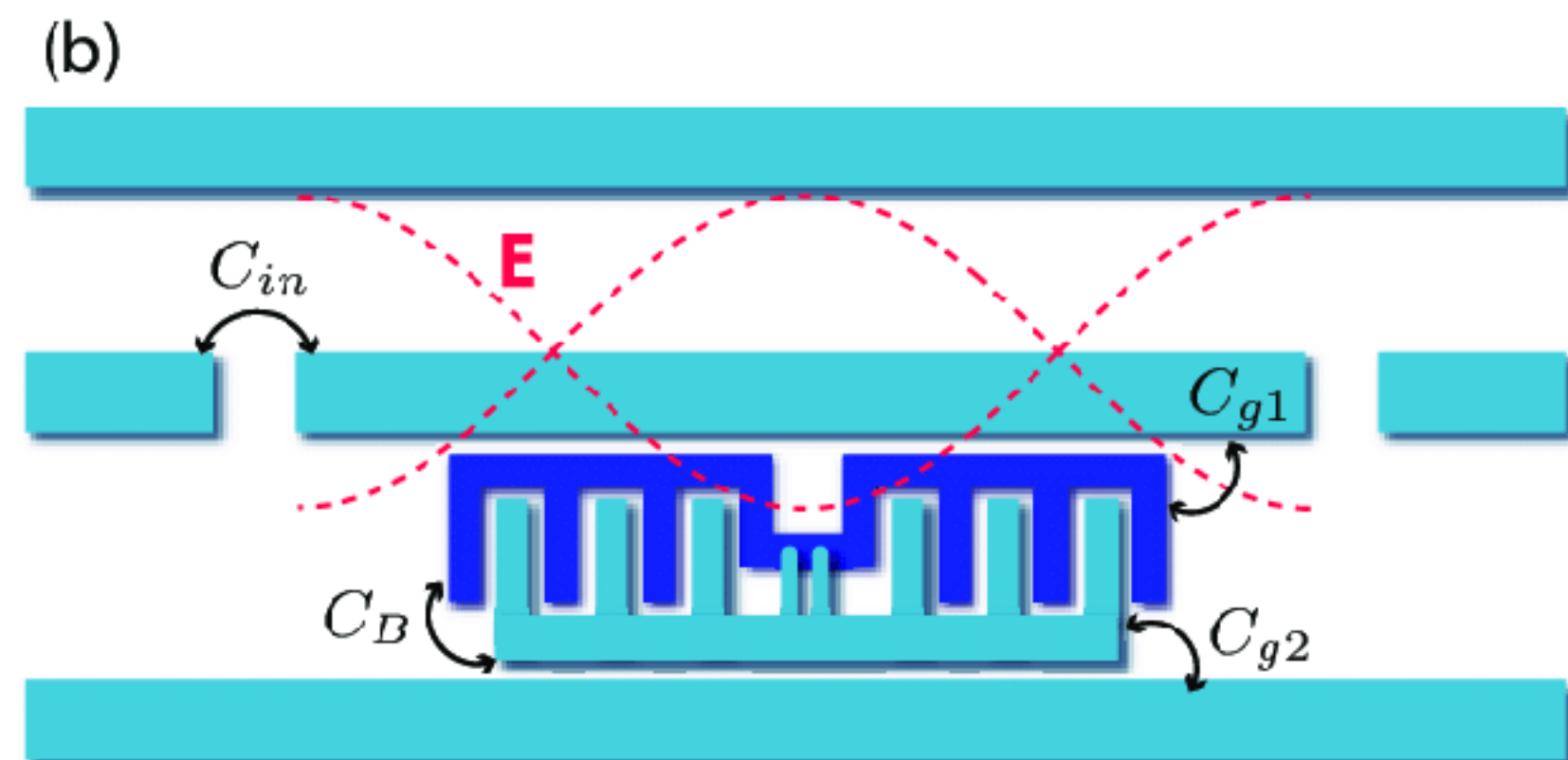
$$\frac{E_J}{E_C} \simeq 100$$

Houck et al. (Nature, 2007)
Koch et al. (PRA, 2007)



$$H_{\text{qubit}} = \frac{1}{2} \Omega \sigma^z$$

$$\Omega/2\pi \sim 5 \text{ GHz}$$

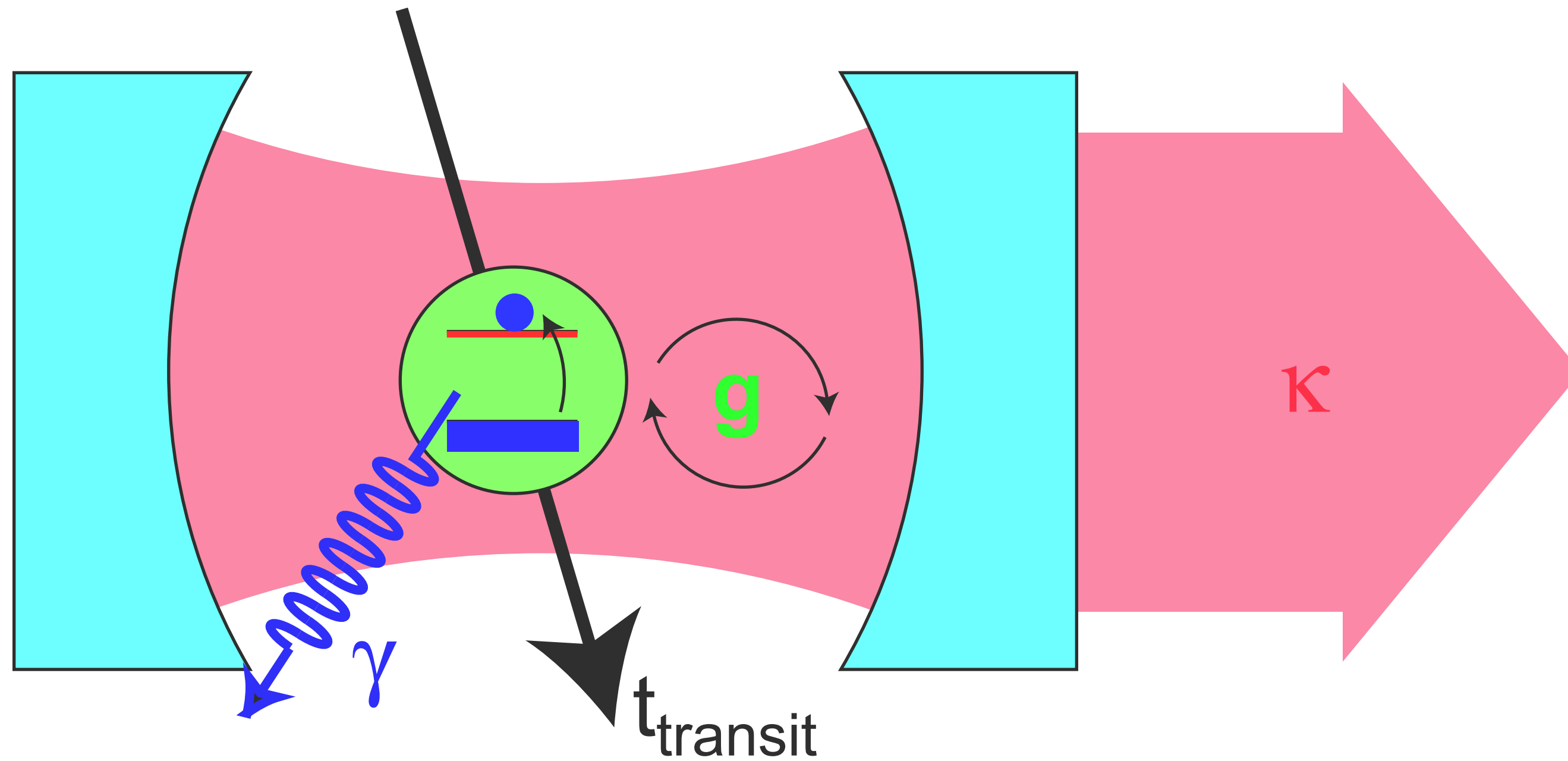


CIRCUIT QED

(Cavity QED system based on superconducting circuit)

“TRADITIONAL” CAVITY QED SYSTEM

FUNDAMENTAL LIGHT-MATTER INTERACTION



ω : photon frequency

Ω : atomic transition frequency

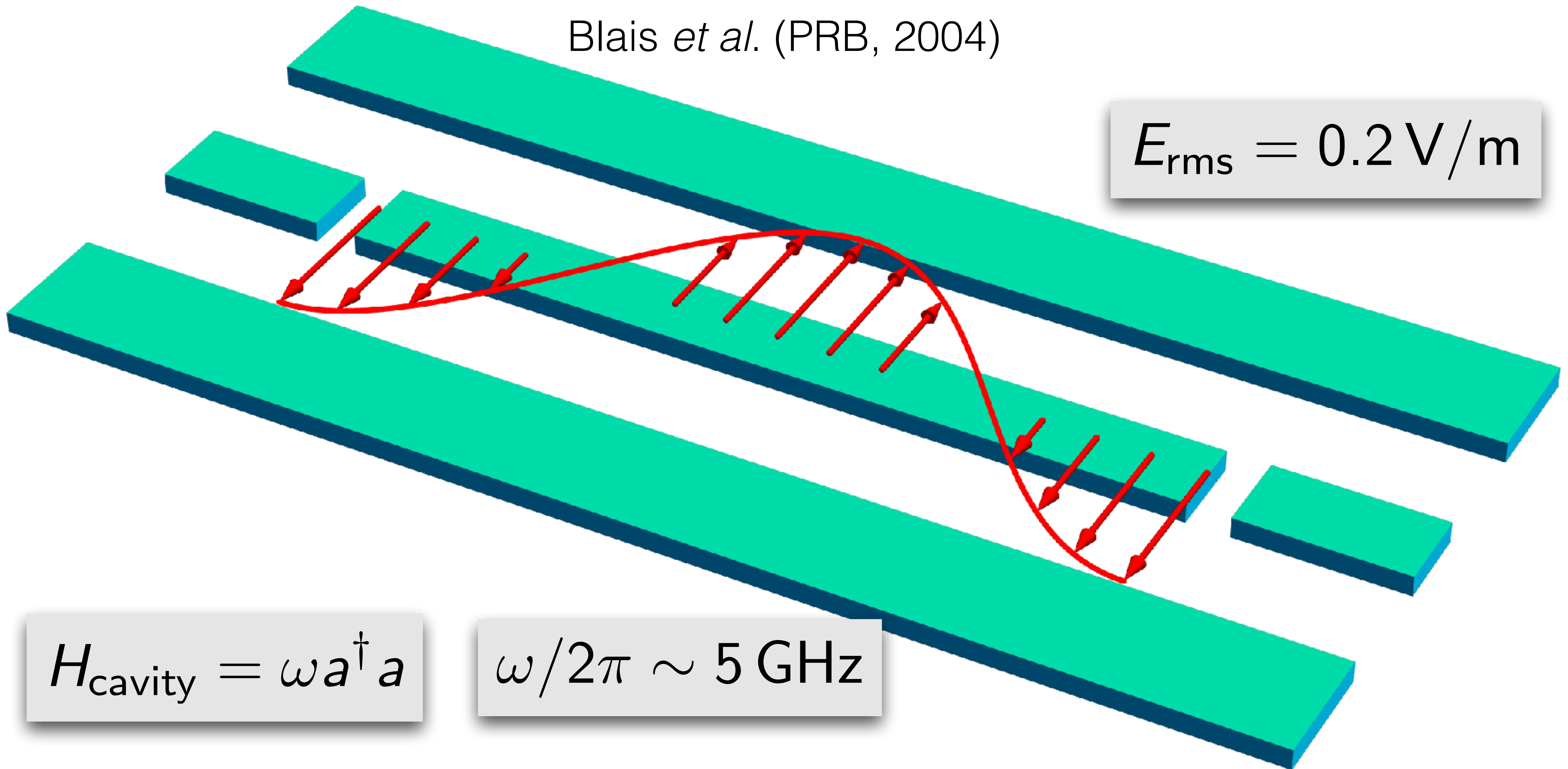
κ : photon loss rate

γ : atomic decoherence rate

RESONATOR (CAVITY)

Blais *et al.* (PRB, 2004)

$$E_{\text{rms}} = 0.2 \text{ V/m}$$



$$H_{\text{cavity}} = \omega a^\dagger a$$

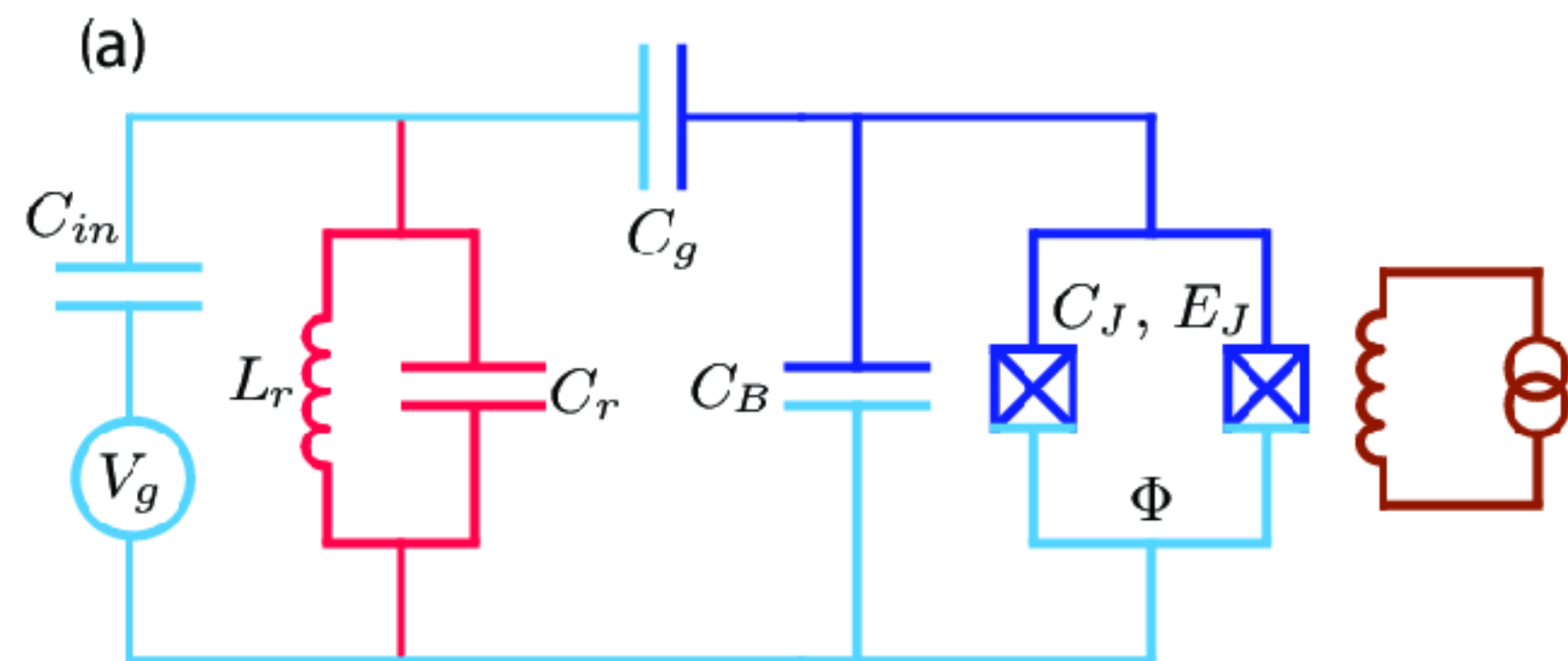
$$\omega/2\pi \sim 5 \text{ GHz}$$

TRANSMON QUBIT

ANHARMONIC OSCILLATOR

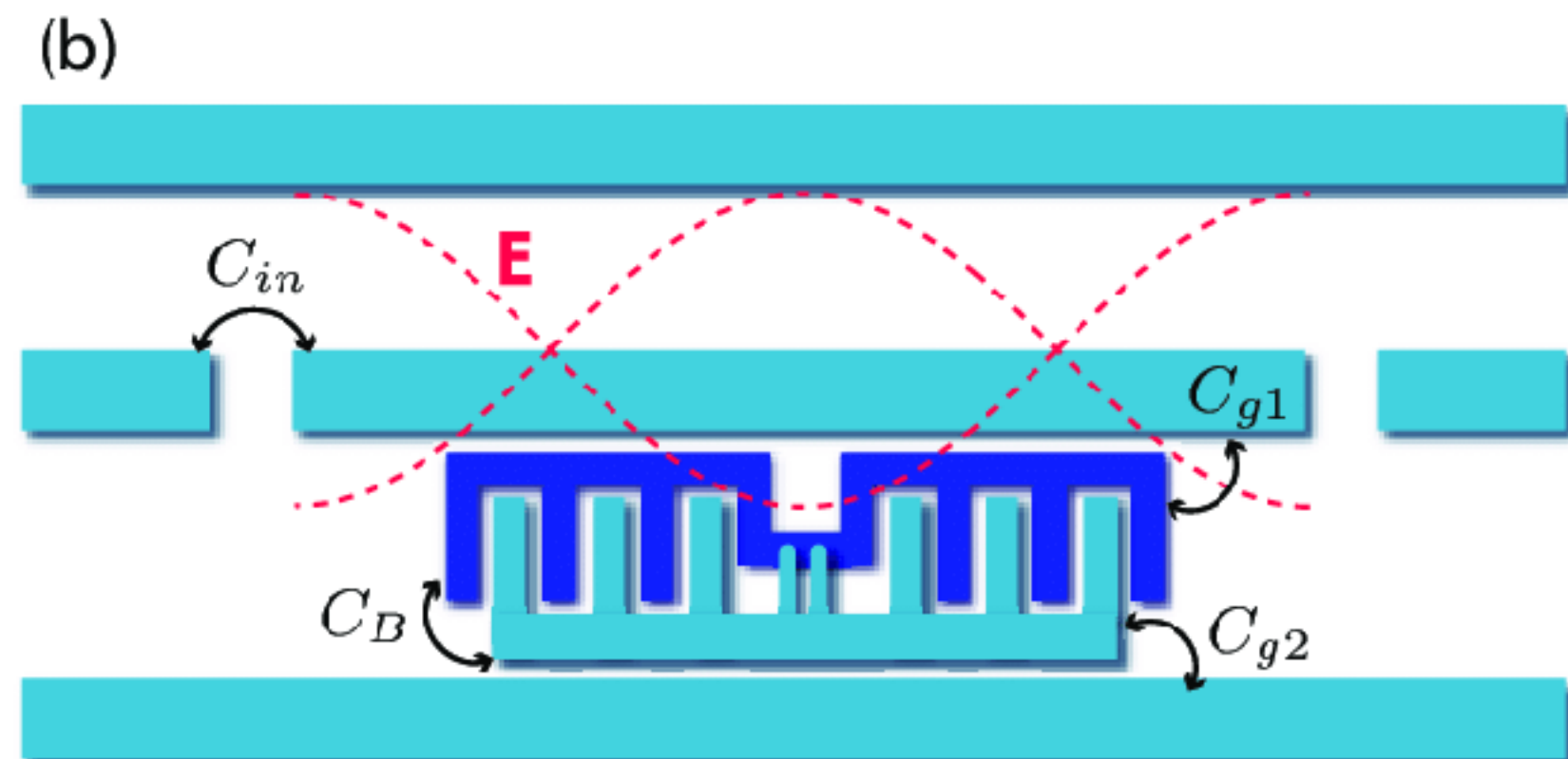
$$\frac{E_J}{E_C} \simeq 100$$

Houck et al. (Nature, 2007)
Koch et al. (PRA, 2007)



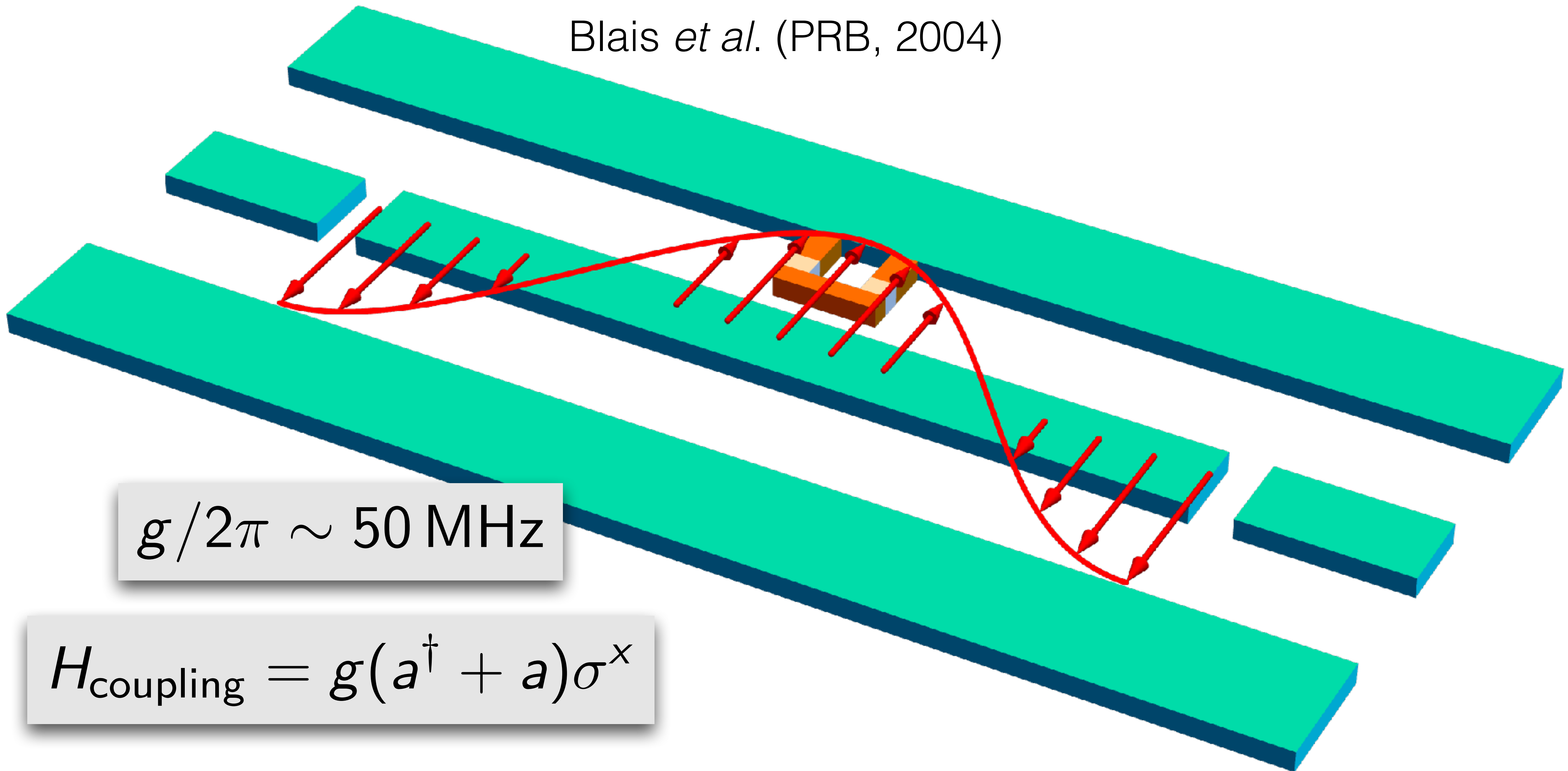
$$H_{\text{qubit}} = \frac{1}{2} \Omega \sigma^z$$

$$\Omega/2\pi \sim 5 \text{ GHz}$$



CIRCUIT QED SYSTEM

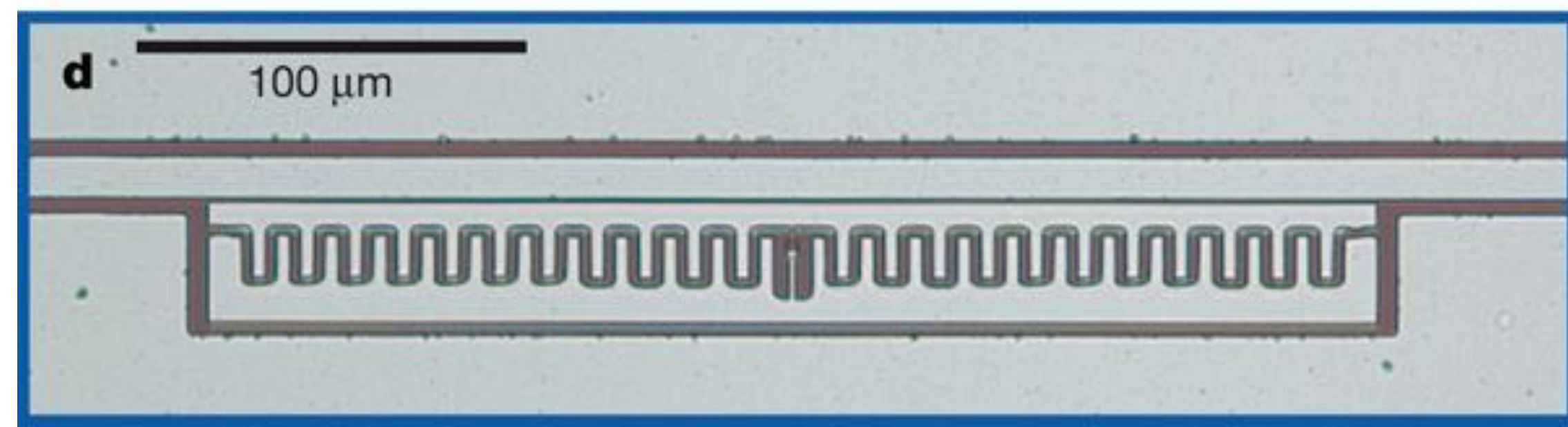
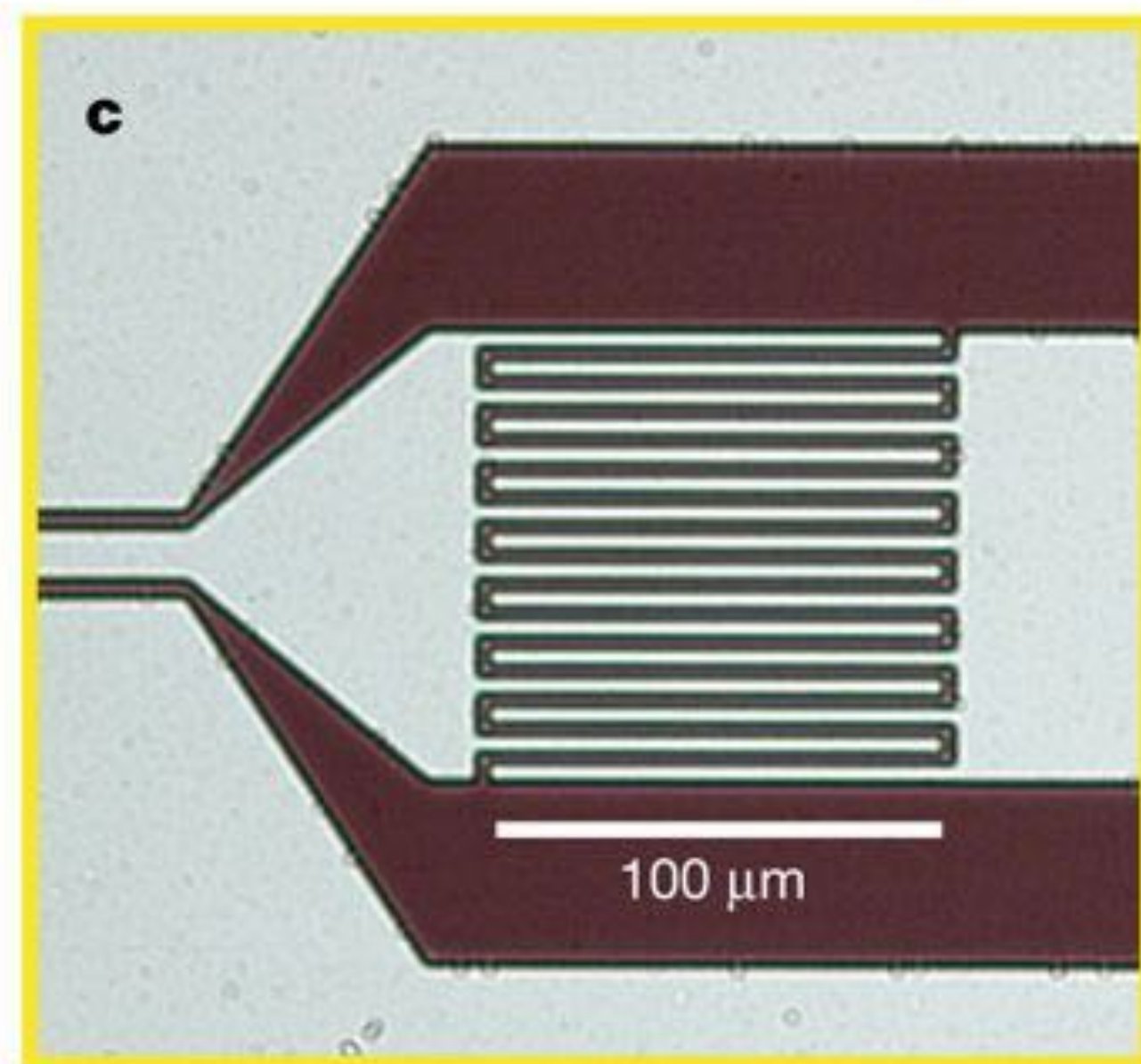
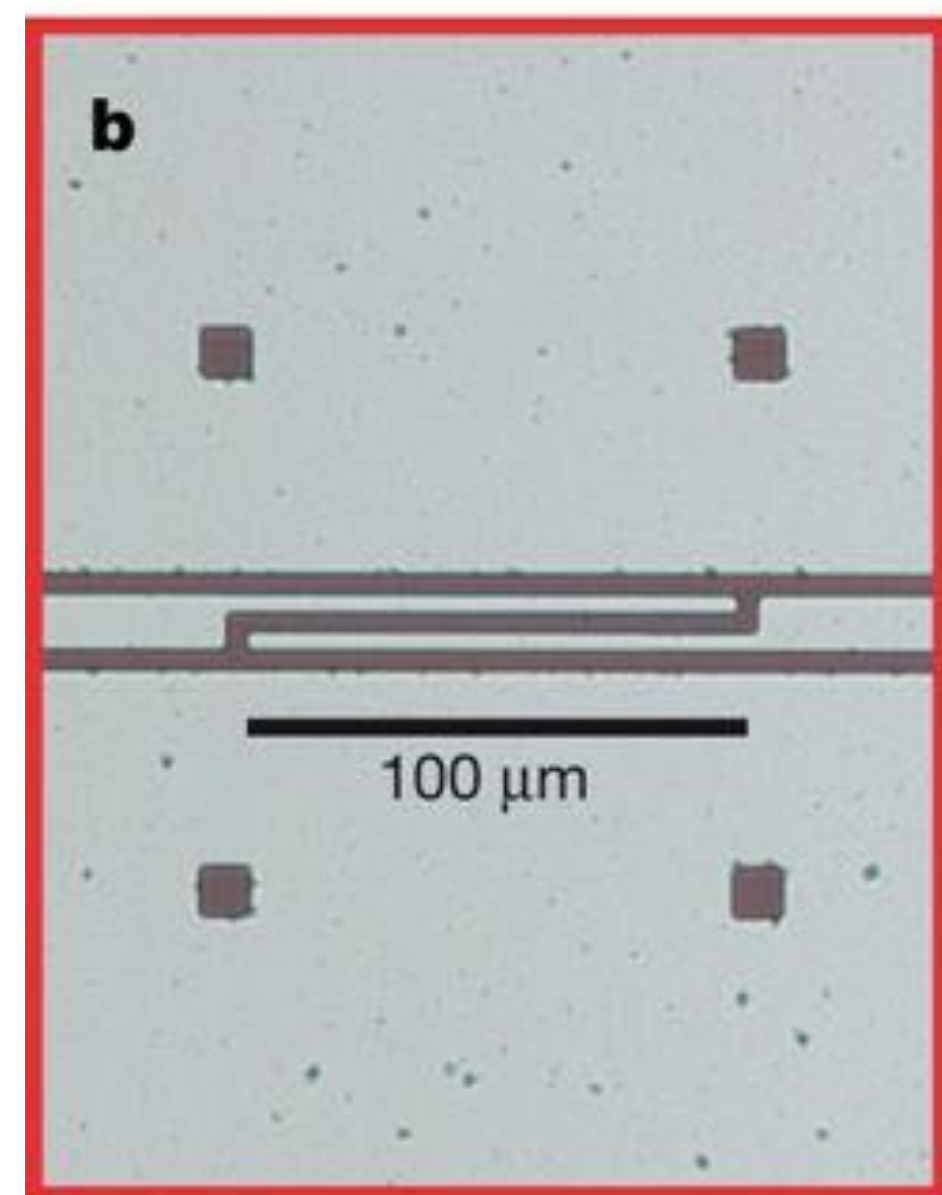
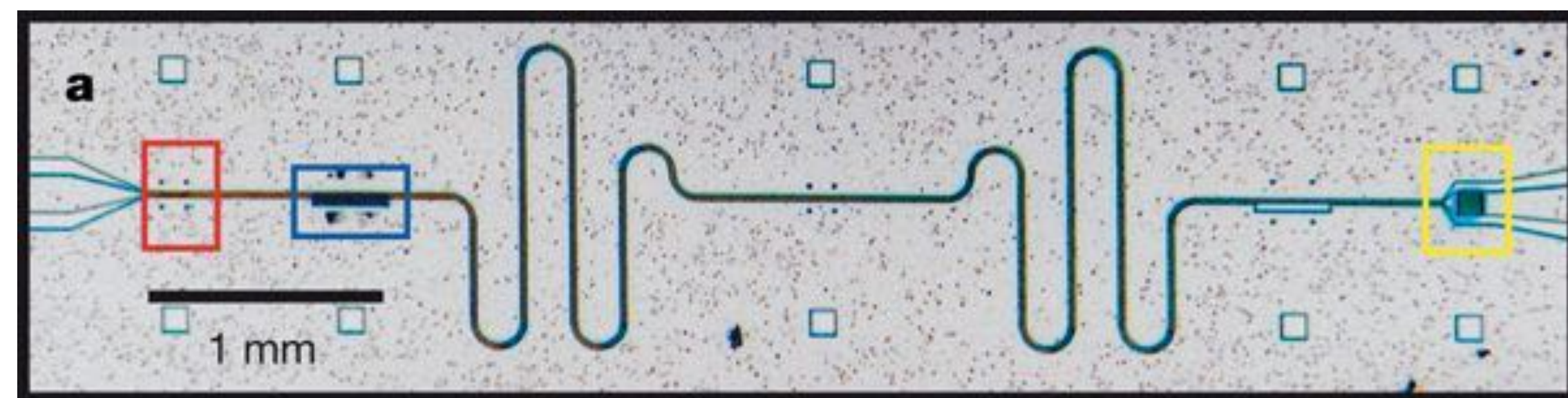
Blais *et al.* (PRB, 2004)



$$g/2\pi \sim 50 \text{ MHz}$$

$$H_{\text{coupling}} = g(a^\dagger + a)\sigma^x$$

CIRCUIT QED SYSTEM



$$g/2\pi \approx 100 \text{ MHz}$$

$$H_{\text{coupling}} = g(a^\dagger + a)\sigma^x$$

Houck et al. (Nature, 2007)

THE RABI MODEL

$$H_{\text{Rabi}} = \omega a^\dagger a + \frac{1}{2} \Omega \sigma^z + g(a^\dagger + a) \sigma^x$$

- κ : cavity photon loss rate
- γ : qubit decoherence rate

$$\omega/2\pi \sim \Omega/2\pi \sim 5 \text{ GHz}$$

$$g/2\pi \sim 50 \text{ MHz}$$

$$g/\omega \sim 10^{-3} - 1$$

- Bishop *et al.*, Nature Physics (2009)
- Niemczyk *et al.*, Nature Physics (2010)
- ...

SO, LET'S SOLVE THE RABI

THE RABI MODEL

$$H_{\text{Rabi}} = \omega a^\dagger a + \frac{1}{2} \Omega \sigma^z + g(a^\dagger + a) \sigma^x$$

- κ : cavity photon loss rate
- γ : qubit decoherence rate

STRONG COUPLING, ULTRA STRONG COUPLING, RESONANCE, ...

$$g \gg \kappa, \gamma$$

“strong” coupling

$$g \sim \sqrt{\omega\Omega}$$

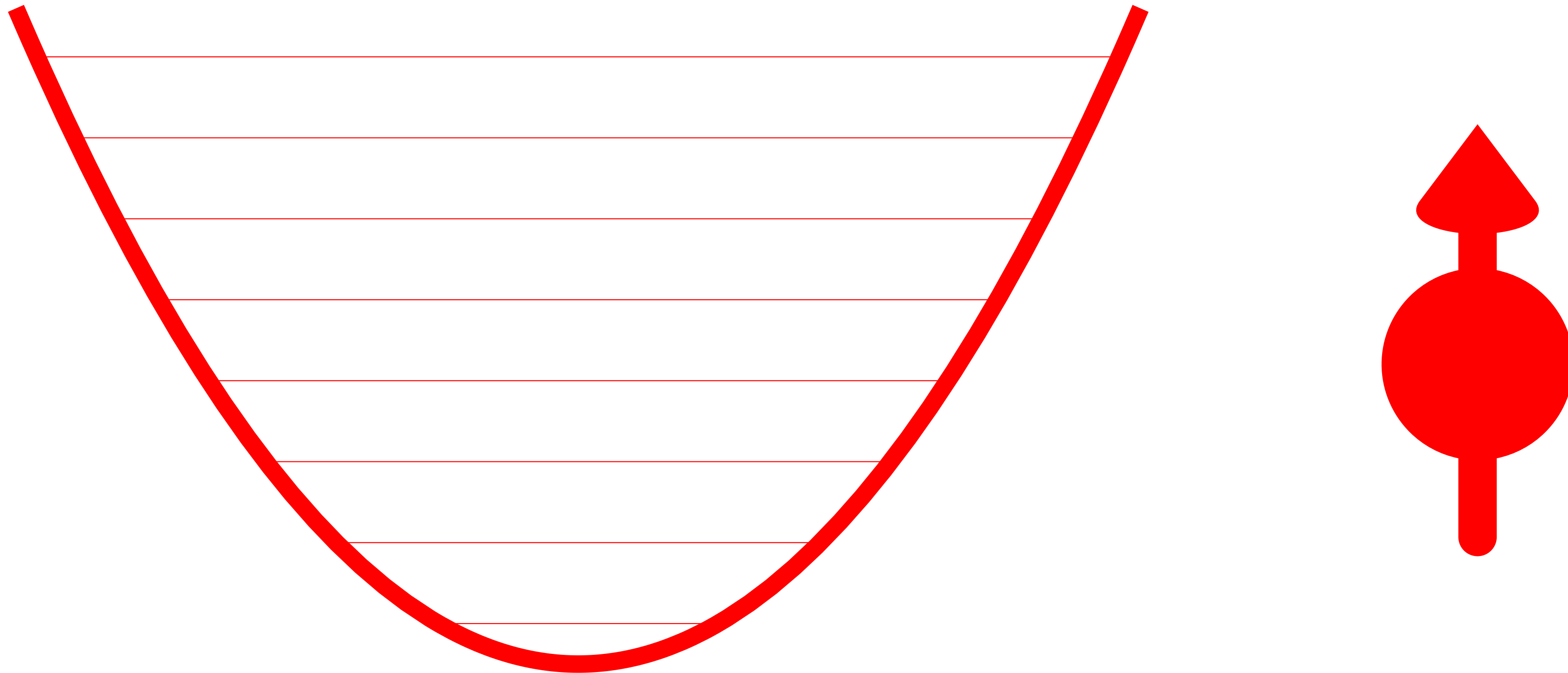
“ultra strong” coupling

$$\omega = \Omega$$

resonance

WEAK-COUPPLING LIMIT

$$g = 0$$



$$H_{\text{Rabi}} = \omega a^\dagger a + \frac{1}{2} \Omega \sigma^z + \cancel{g(a^\dagger + a) \sigma^x}$$

WEAK-COUPPLING LIMIT

$$g = 0$$

$|3, \uparrow\rangle$

$|4, \downarrow\rangle$

$|2, \uparrow\rangle$

$|3, \downarrow\rangle$

$|1, \uparrow\rangle$

$|2, \downarrow\rangle$

$|0, \uparrow\rangle$

$|1, \downarrow\rangle$

$$\{ |n, \sigma\rangle \mid n = 0, 1, 2, \dots ; \sigma = \uparrow, \downarrow \}$$

$|0, \downarrow\rangle$

“STRONG” COUPLING

Jaynes-Cummings Model (Rotating Wave Approximation)

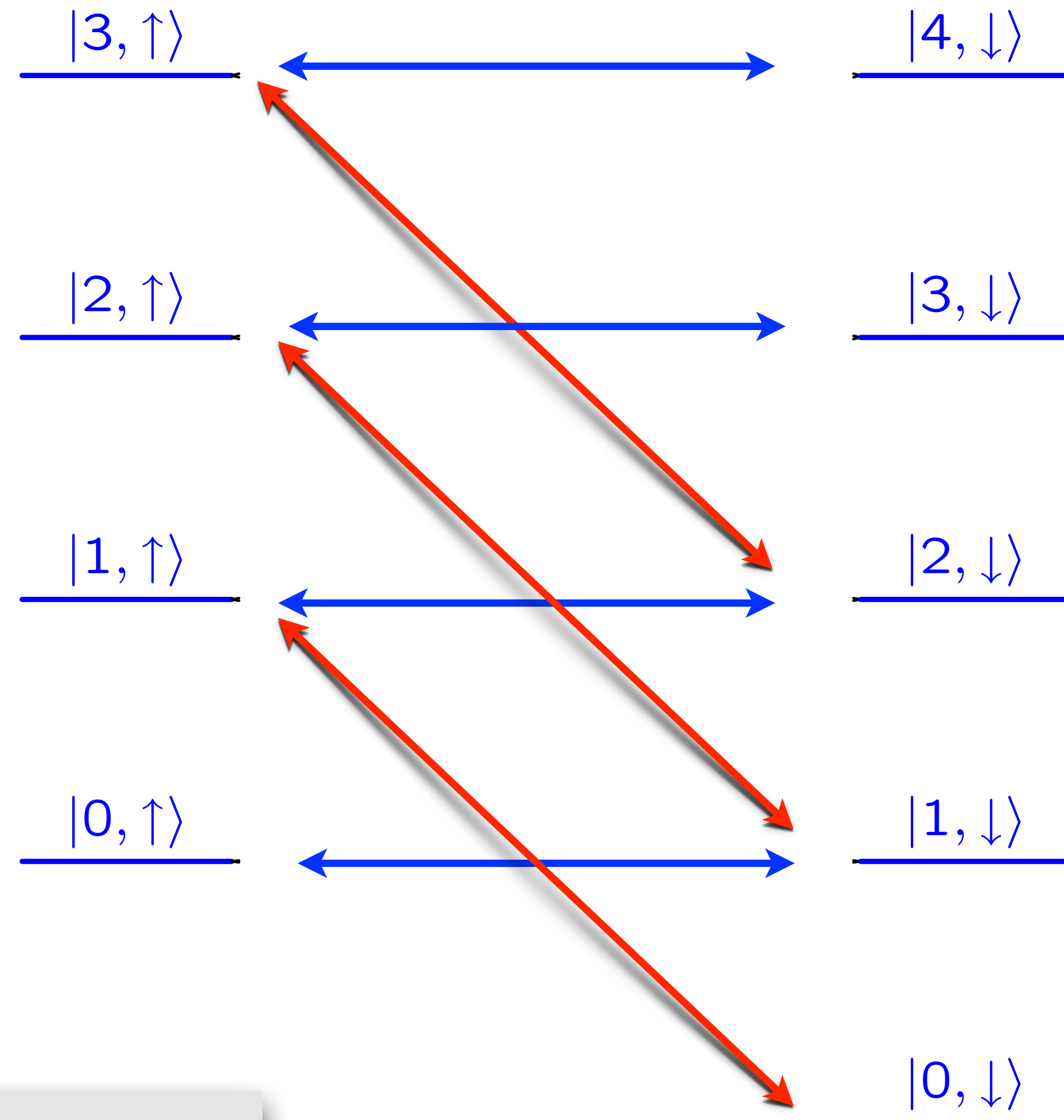
$$\kappa, \gamma \ll g \ll \omega$$

$$(a^\dagger + a)(\sigma^+ + \sigma^-) = \underbrace{(a^\dagger \sigma^- + a \sigma^+)}_{\text{rotating}} + \underbrace{(a^\dagger \sigma^+ + a \sigma^-)}_{\text{counter-rotating}}$$

$$H_{\text{Rabi}} = \omega a^\dagger a + \frac{1}{2} \Omega \sigma^z + g(a^\dagger + a) \sigma^x$$

ROTATING VS. COUNTER-ROTATING

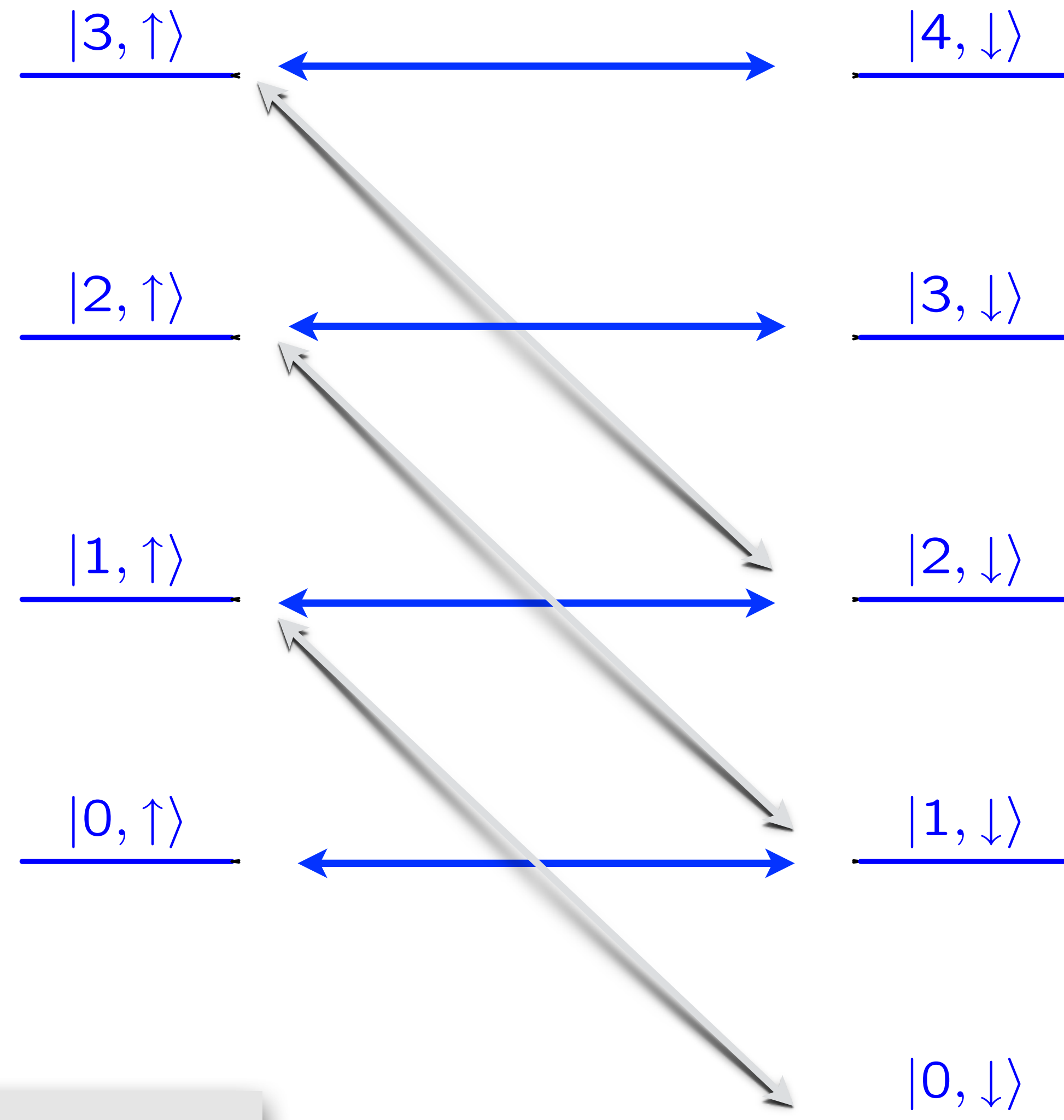
$$\underbrace{(a^\dagger \sigma^- + a \sigma^+)}_{\text{rotating}} + \underbrace{(a^\dagger \sigma^+ + a \sigma^-)}_{\text{counter-rotating}}$$



$$\{ |n, \sigma\rangle \mid n = 0, 1, 2, \dots ; \sigma = \uparrow, \downarrow \}$$

ROTATING VS. ~~COUNTER-ROTATING~~

$$\underbrace{(a^\dagger \sigma^- + a \sigma^+)}_{\text{rotating}} + \underbrace{(a^\dagger \sigma^+ + a \sigma^-)}_{\text{counter-rotating}}$$



$$\{ |n, \sigma\rangle \mid n = 0, 1, 2, \dots ; \sigma = \uparrow, \downarrow \}$$

THE JAYNES-CUMMINGS MODEL

(ROTATING WAVE APPROXIMATION)

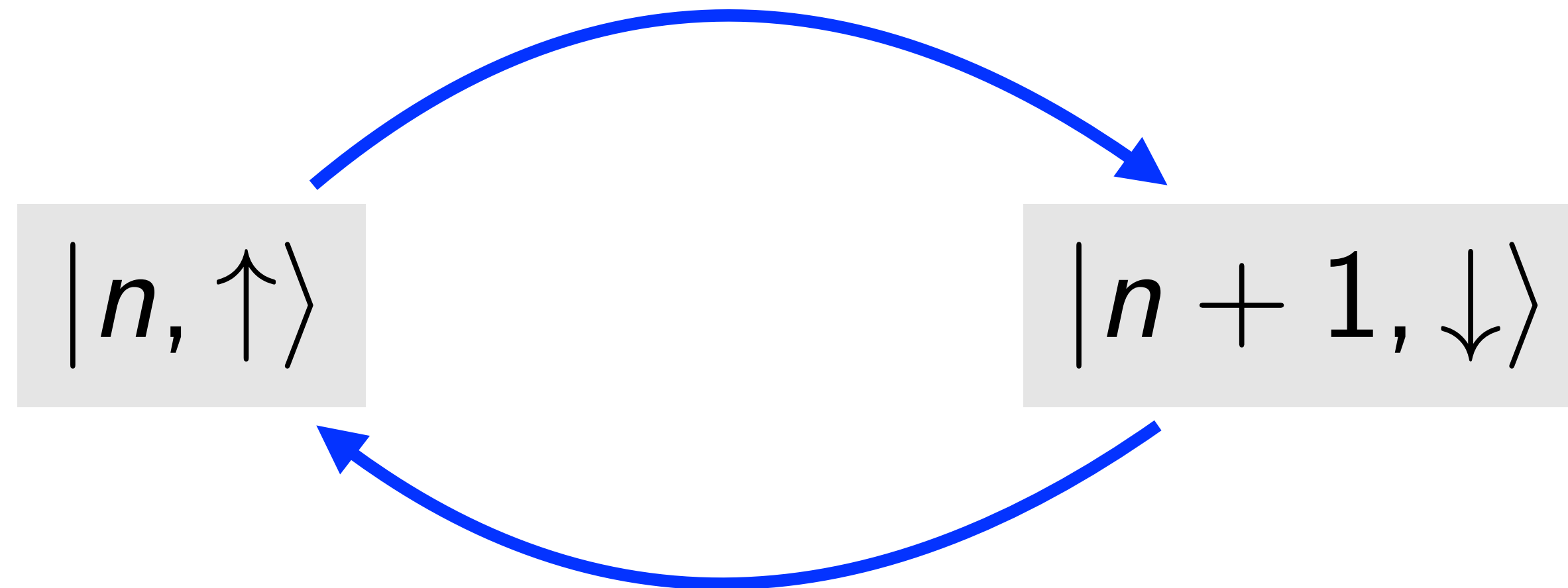
$$H_{JC} = \omega a^\dagger a + \frac{1}{2}\Omega\sigma^z + g(a\sigma^+ + a^\dagger\sigma^-) + \underbrace{g(a\sigma^- + a^\dagger\sigma^+)}_{\text{counter-rotating}}$$

- κ : cavity photon loss rate
- γ : qubit decoherence rate

THE JAYNES-CUMMINGS MODEL

(ROTATING WAVE APPROXIMATION)

$$H_{JC} = \omega a^\dagger a + \frac{1}{2}\Omega\sigma^z + g(a\sigma^+ + a^\dagger\sigma^-) + \underbrace{g(a\sigma^- + a^\dagger\sigma^+)}_{\text{counter-rotating}}$$

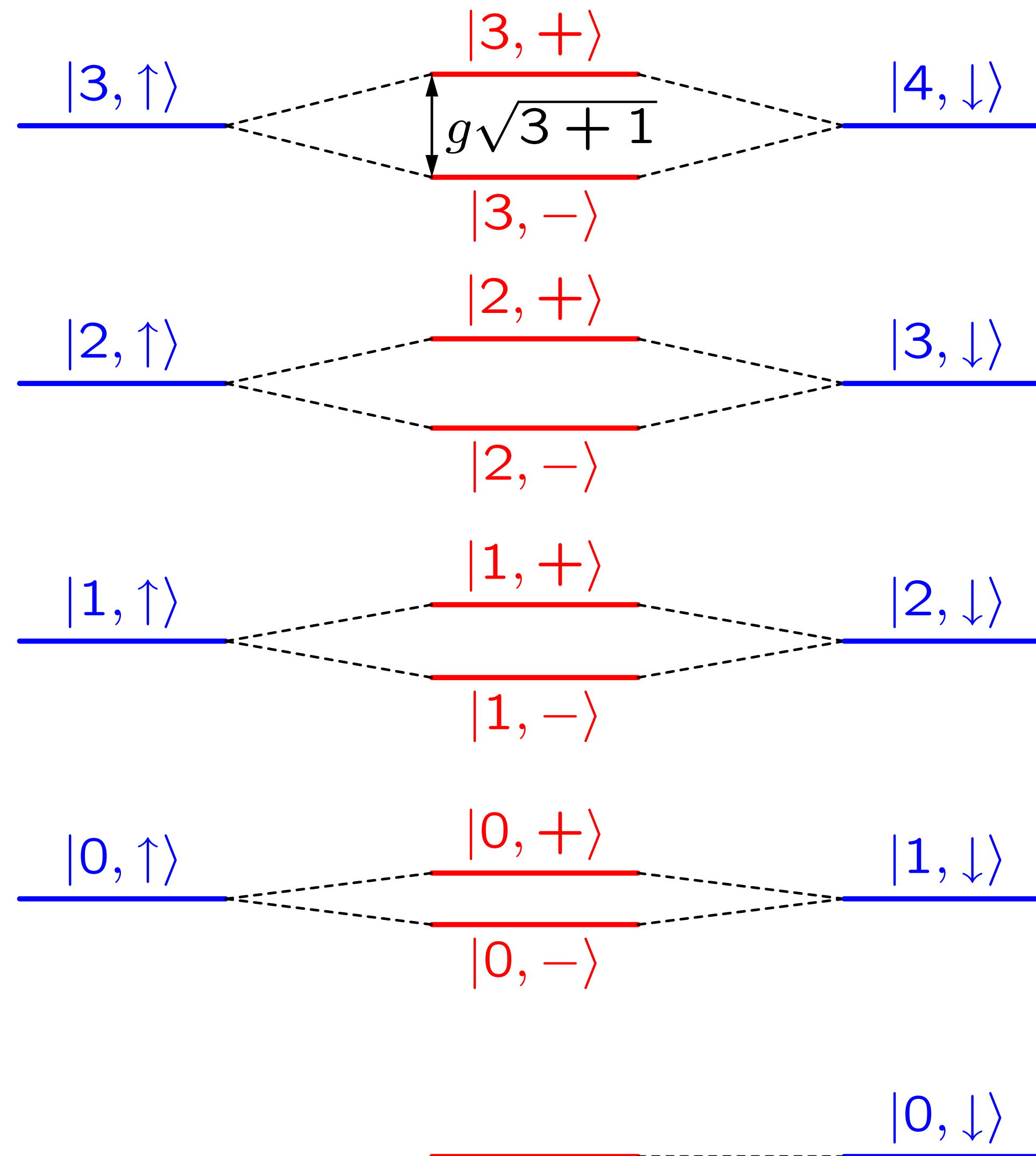


$$\begin{bmatrix} |n, \uparrow\rangle \\ |n+1, \downarrow\rangle \end{bmatrix} = \left\{ (n+1/2)\omega + \begin{bmatrix} (\Omega - \omega)/2 & g\sqrt{n+1} \\ g\sqrt{n+1} & -(\Omega - \omega)/2 \end{bmatrix} \right\} \begin{bmatrix} |n, \uparrow\rangle \\ |n+1, \downarrow\rangle \end{bmatrix}$$

$$E_{n,\pm} = (n+1/2)\omega \pm \sqrt{g^2(n+1) + (\Omega - \omega)^2/4}$$

“STRONG” COUPLING

Jaynes-Cummings Model (Rotating Wave Approximation)

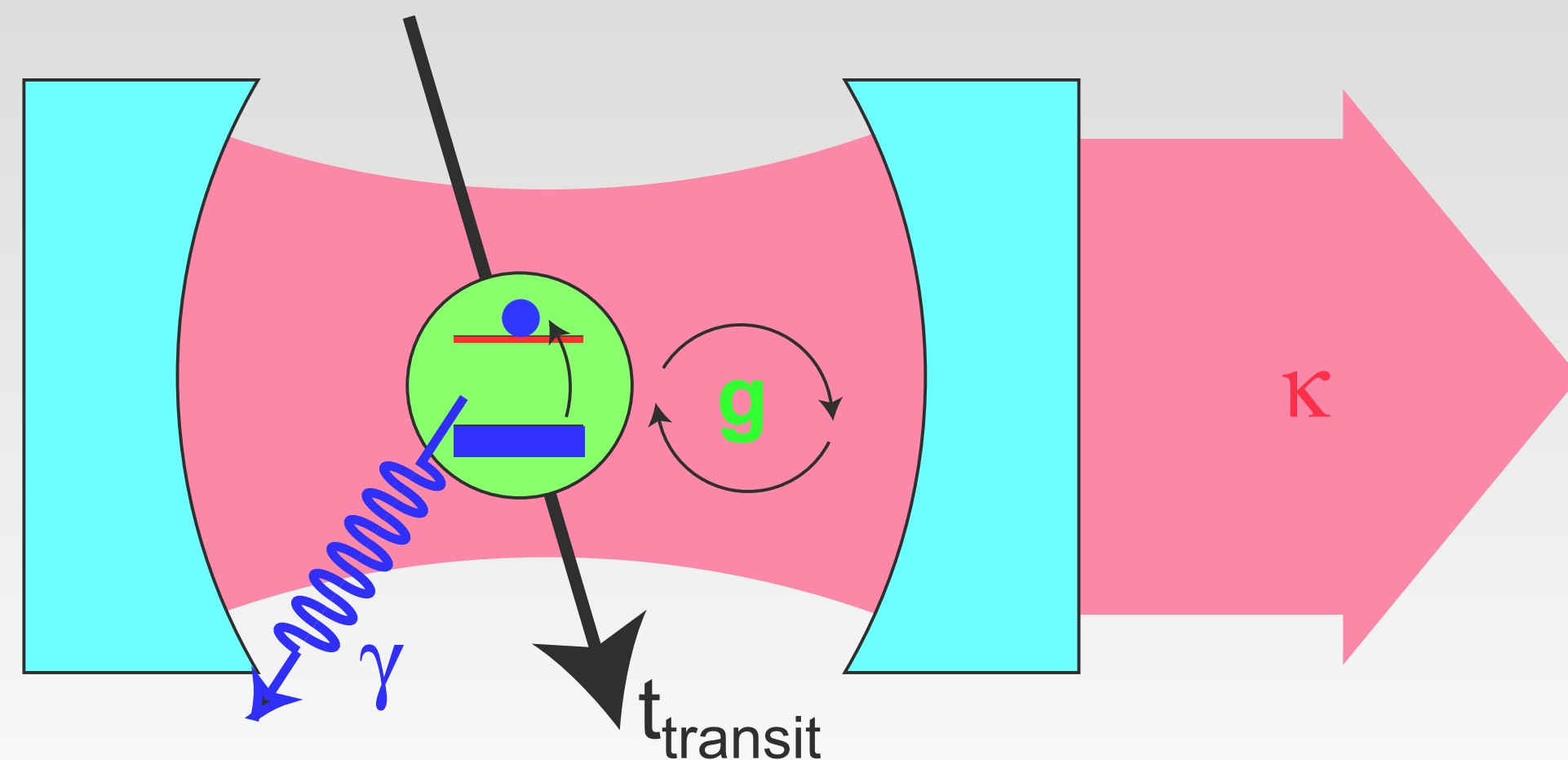


$$g \ll \omega$$

$$g(a^\dagger + a)(\sigma^+ + \sigma^-) \rightarrow \underbrace{g_1(a^\dagger \sigma^- + a \sigma^+)}_{\text{rotating}} + \underbrace{g_2(a^\dagger \sigma^+ + a \sigma^-)}_{\text{counter-rotating}}$$

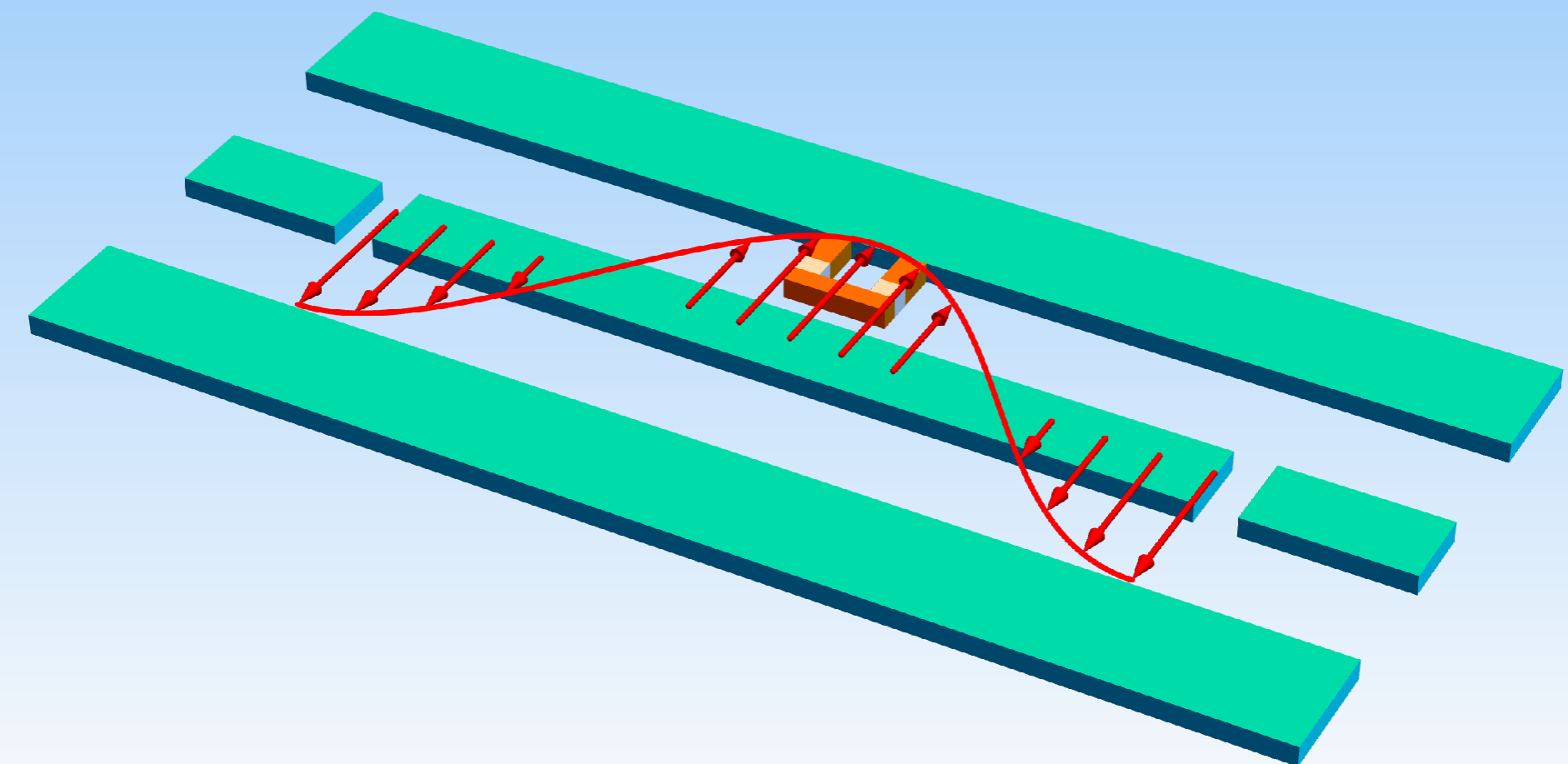
$$\kappa, \gamma \ll g \ll \omega$$

$$g = g_1 = g_2$$



$$g \sim \sqrt{\omega \Omega}$$

$$g_1 \neq g_2$$

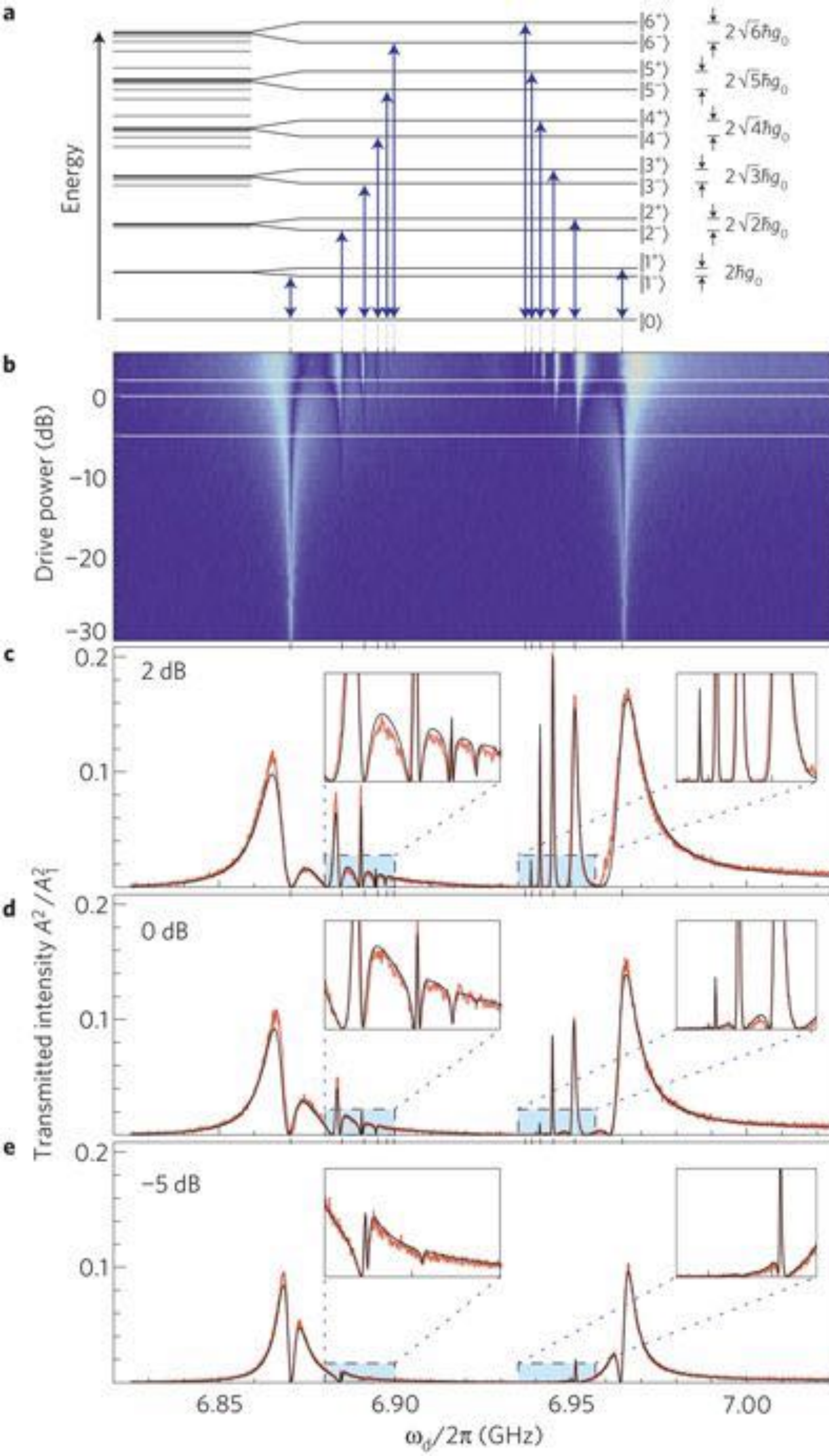


ULTRA-STRONG COUPLING

ADV. QUANT. TECH. 3, 2000085 (2020)

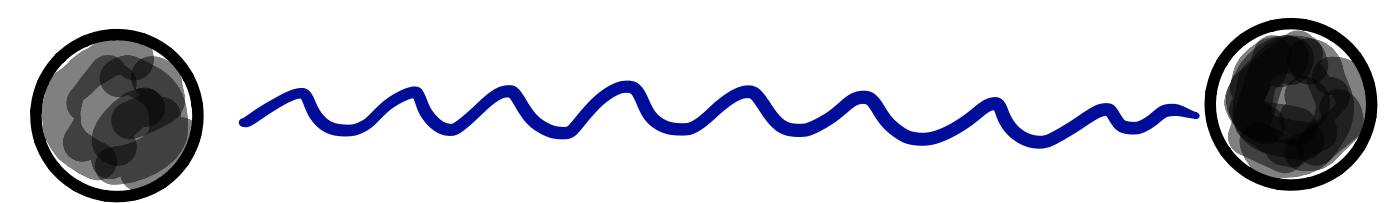
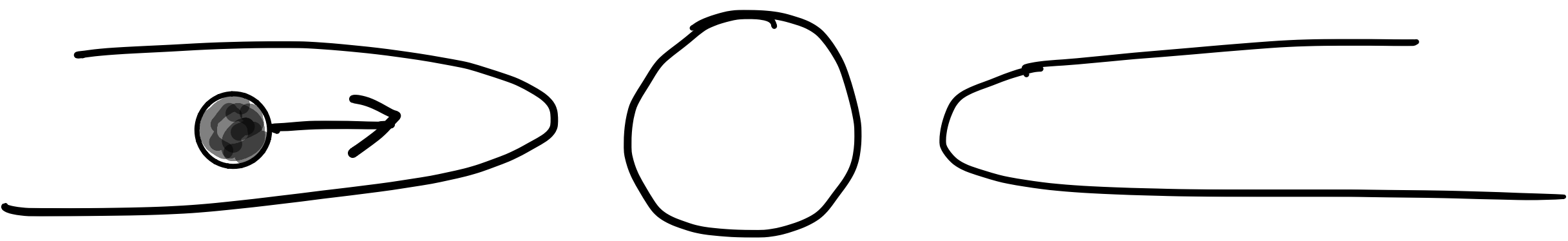


Bishop et al., Nature Physics (2009)
Bishop, PhD Thesis (2010)



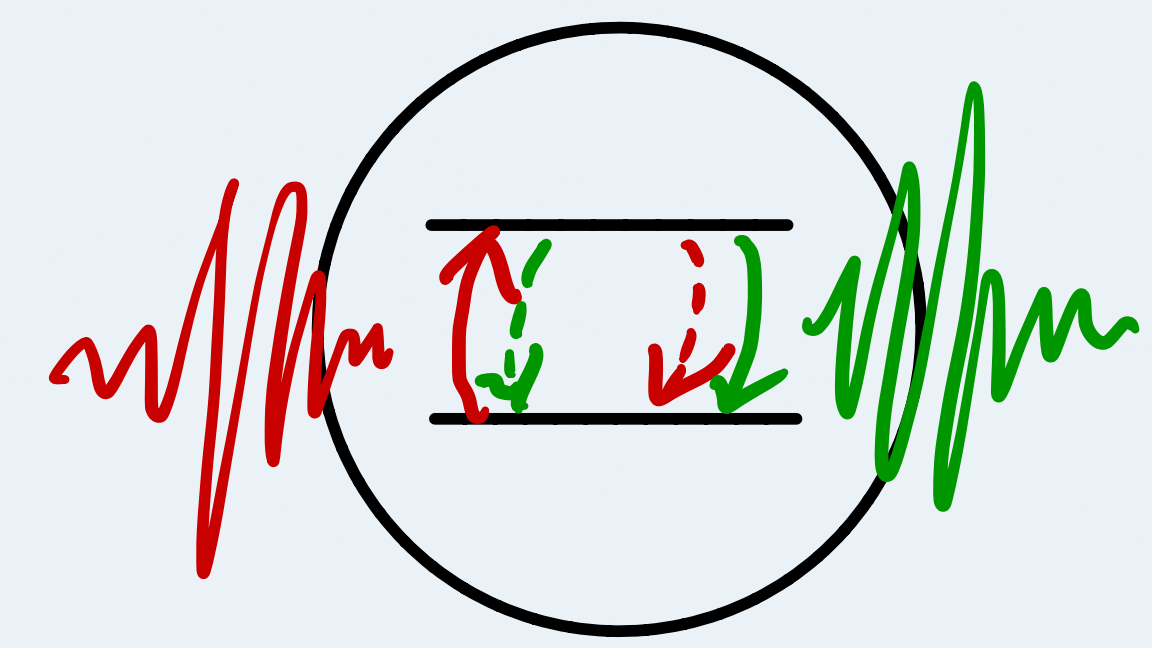
VACUUM RABI OSCILLATION

Bishop et al., Nature Physics (2009)
Bishop, PhD Thesis (2010)



$$\sim \frac{e^2}{2C} \equiv E_C$$

COULOMB BLOCKADE

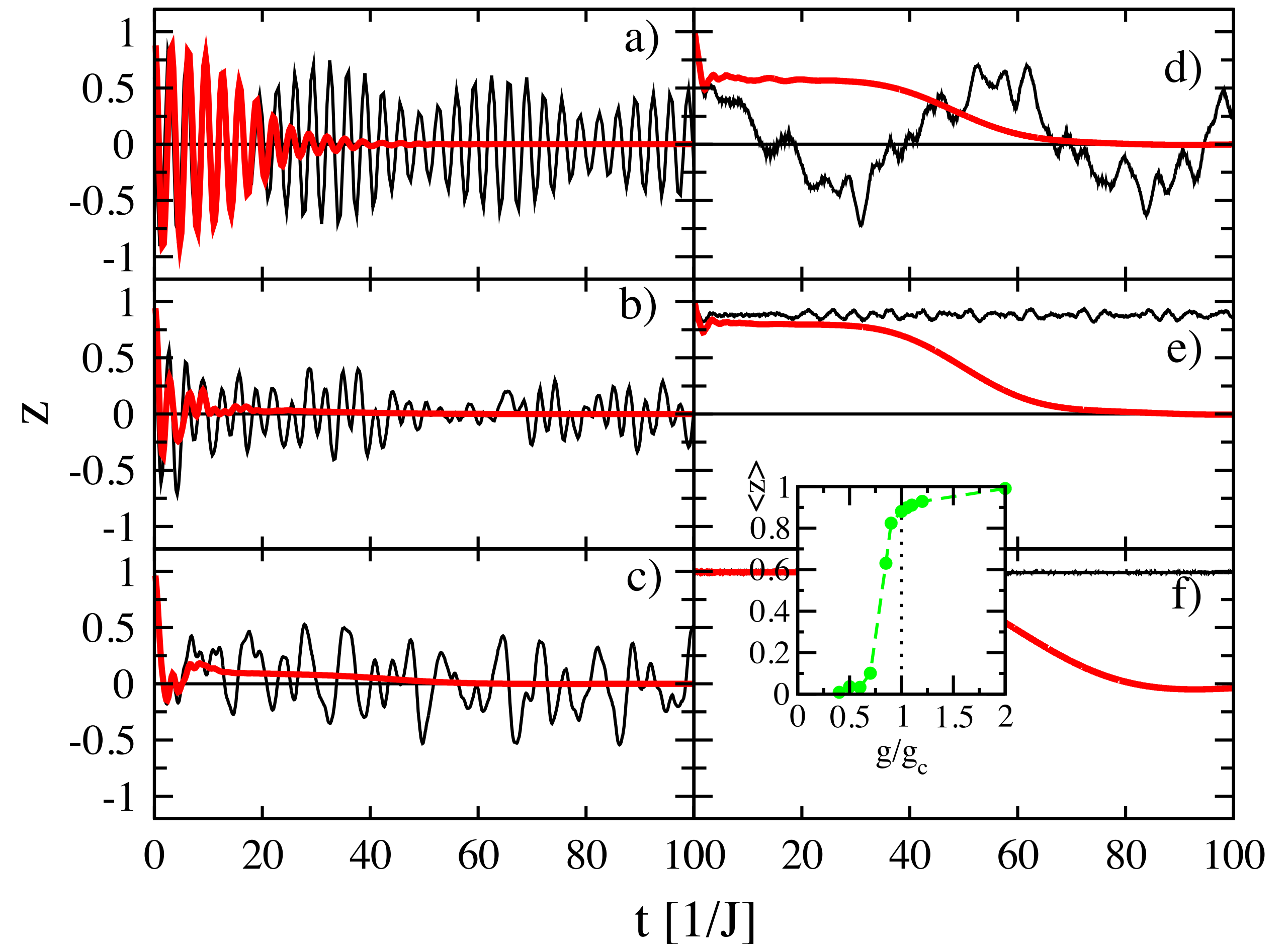
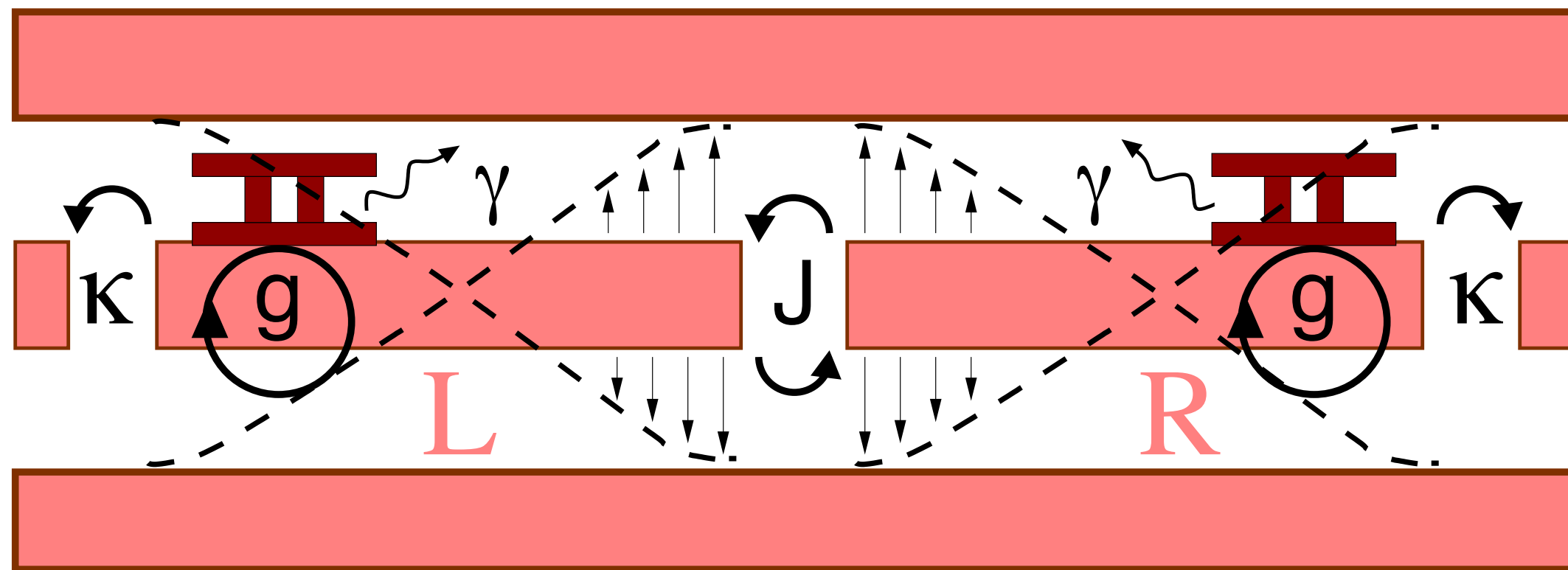


$$\sim g(a^\dagger + a)(\sigma^+ + \sigma^-)$$

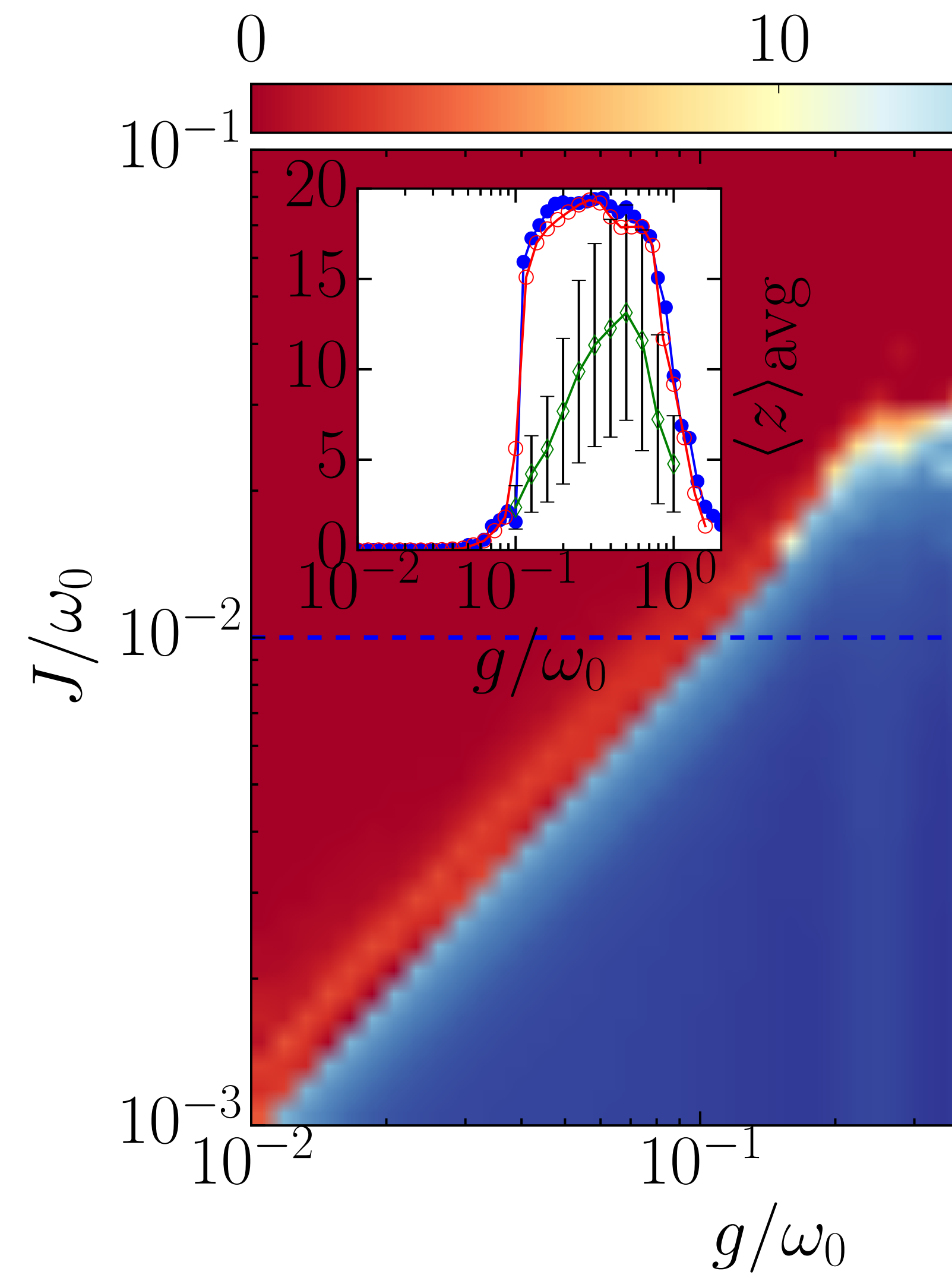
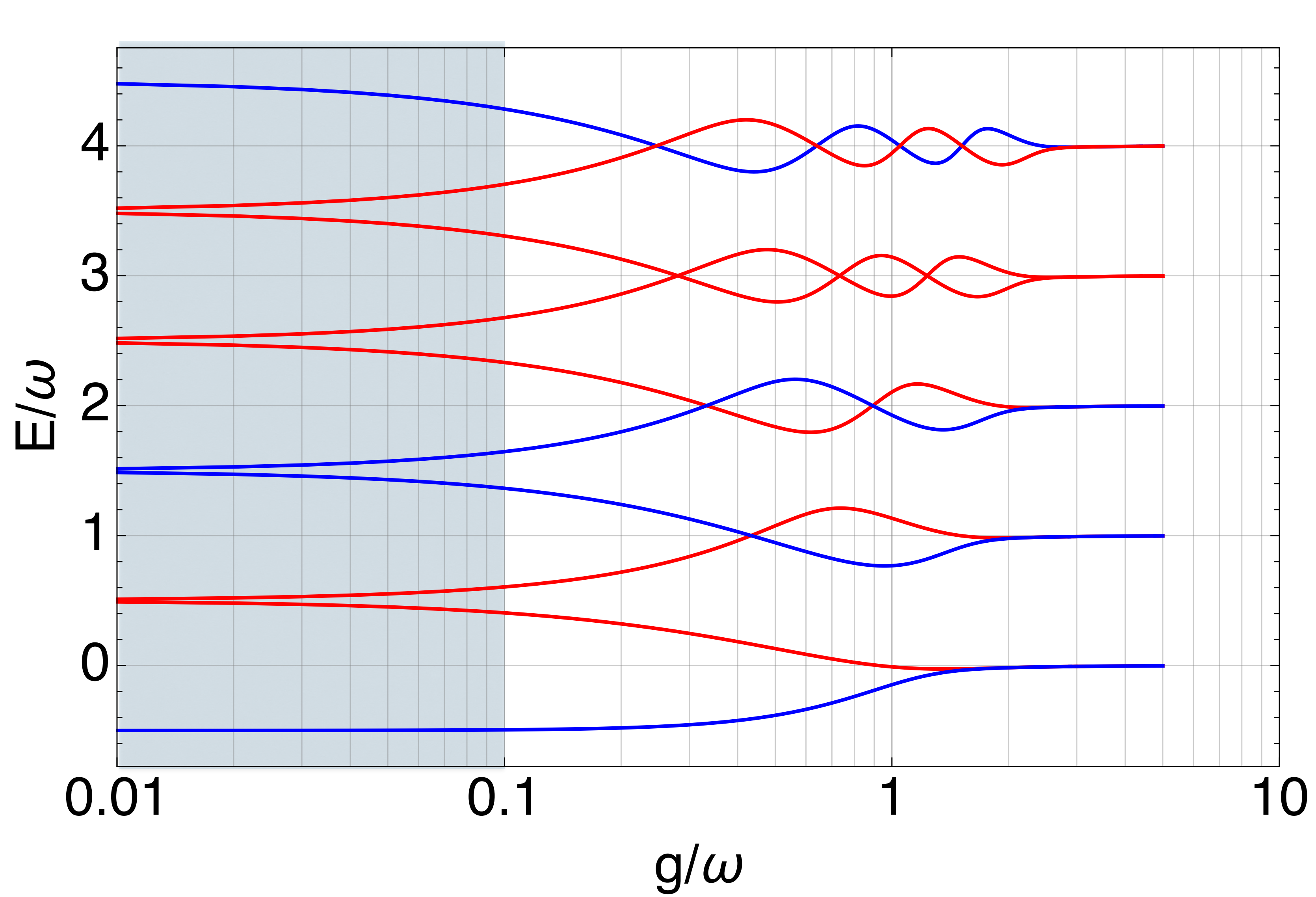
PHOTON BLOCKADE

PHOTON BLOCKADE

Schmidt *et al.*, (PRB, 2010); Raftery *et al.*, (PRX, 2014)

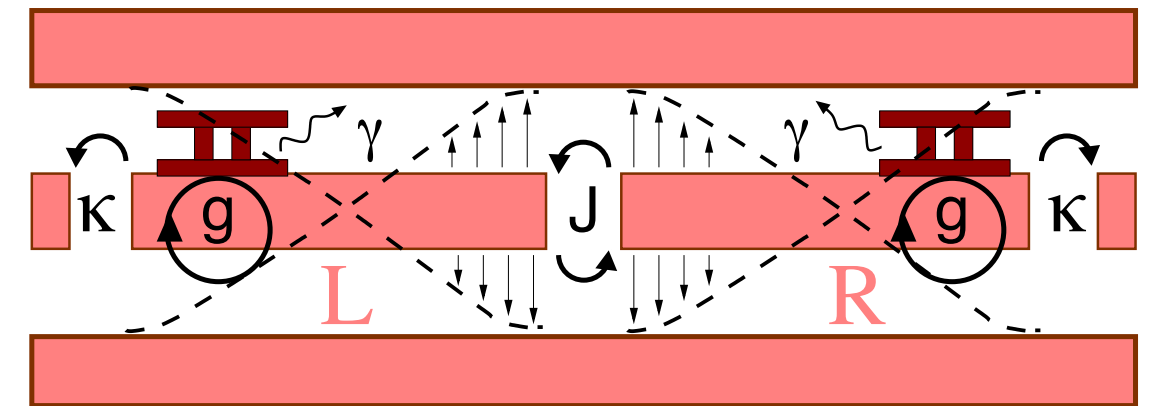
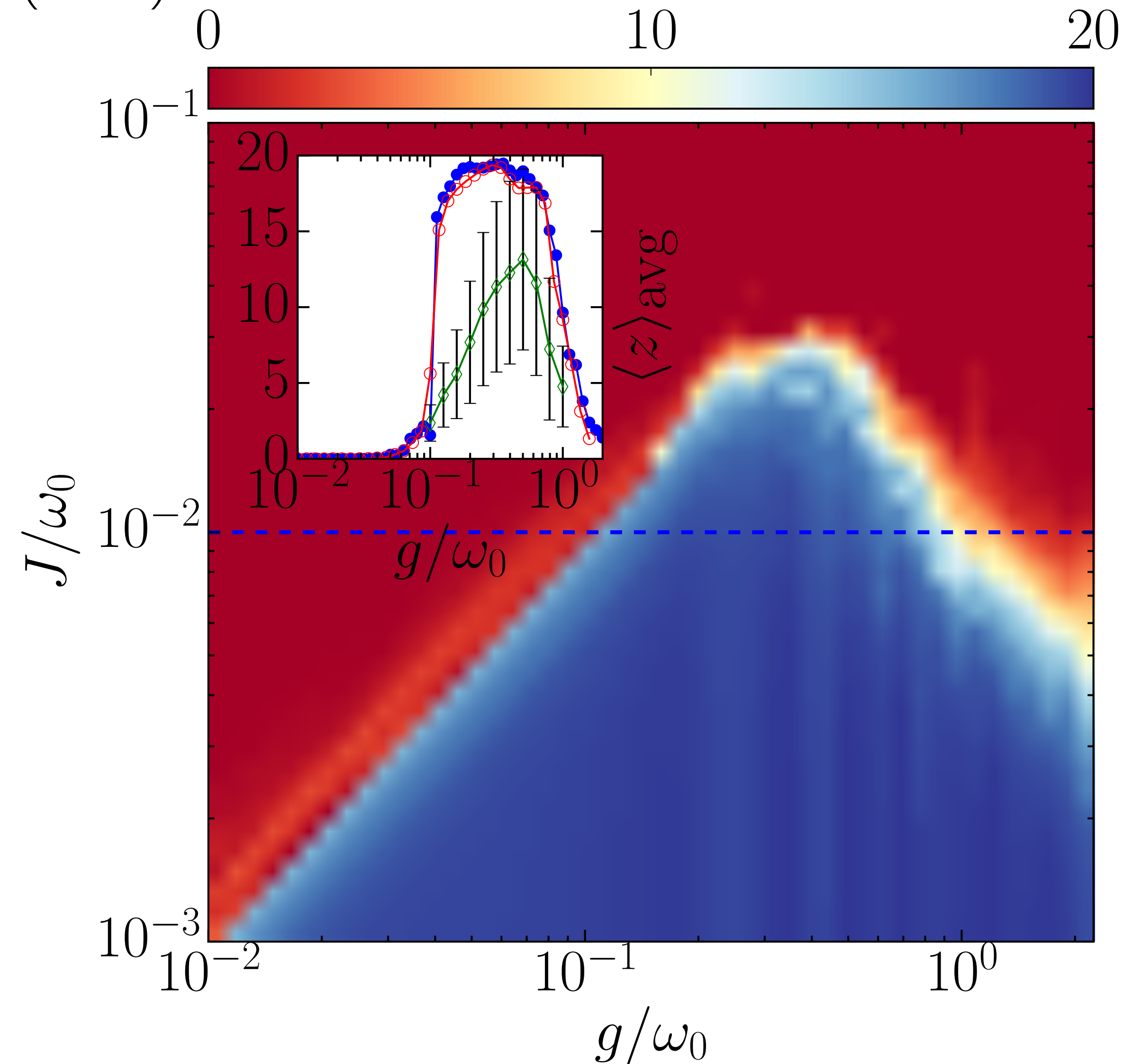


red curve: no decoherence, no photon loss
black curve: small decoherence+photon loss

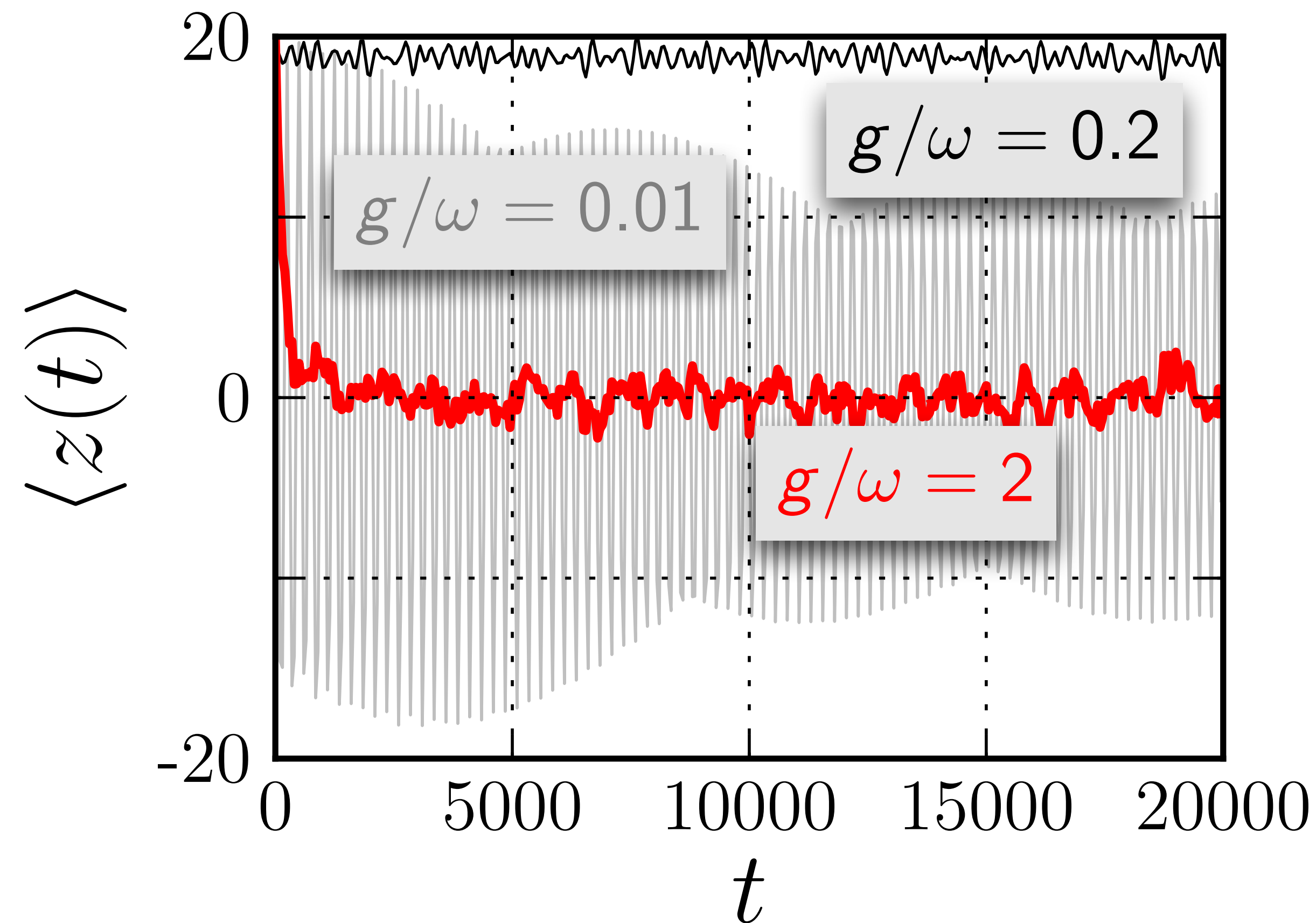
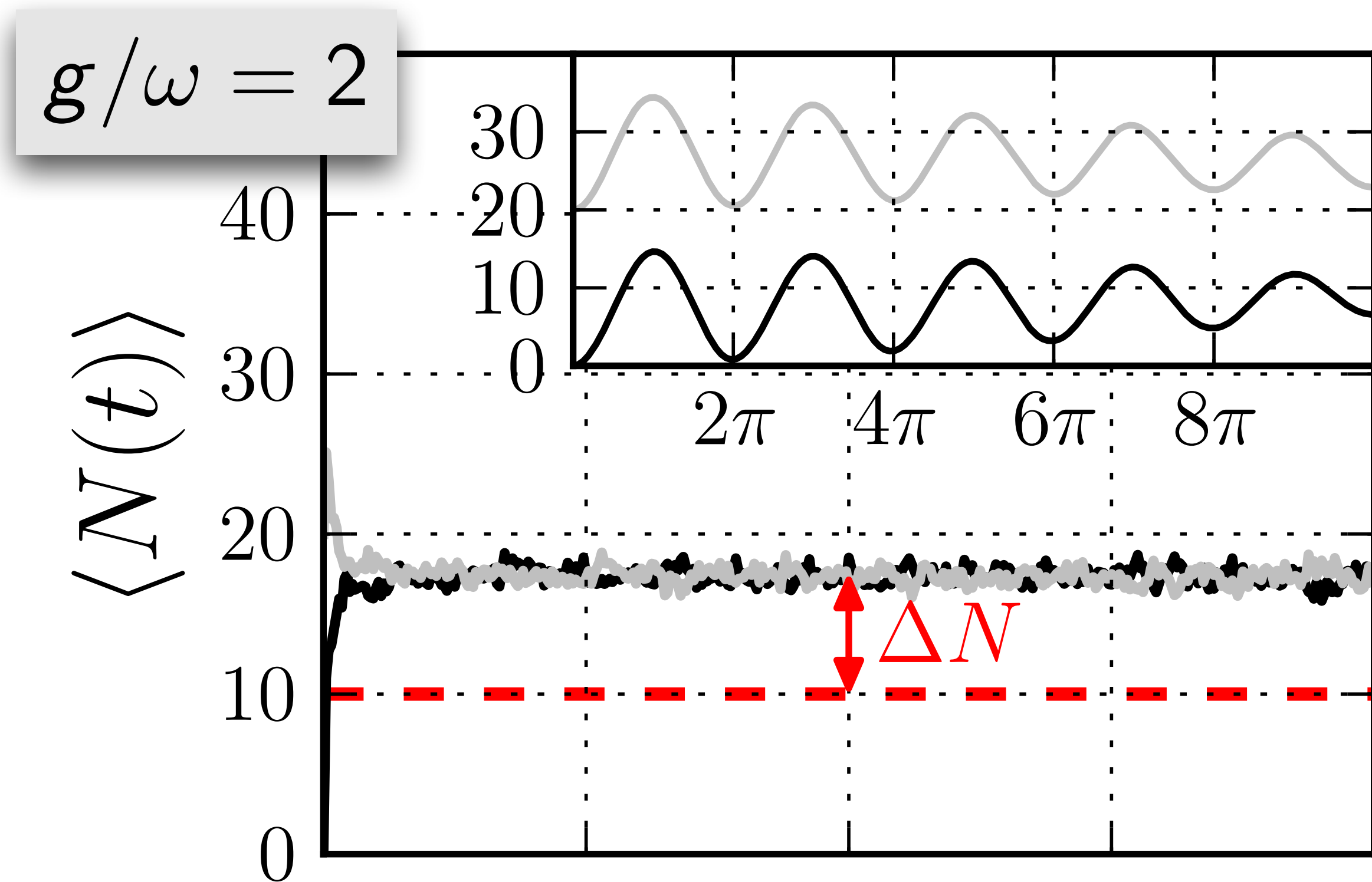


RECURRENT DELOCALIZATION

Phys. Rev. Lett. 116, 153601 (2016)



QUASI-EQUILIBRIUM



Phys. Rev. Lett. 116, 153601 (2016)

BACK TO THE RABI MODEL
but from the (infinitely) strong-coupling limit

INFINITELY STRONG-COUPPLING LIMIT

$$g \rightarrow \infty$$

$$H_{\text{Rabi}} = \omega a^\dagger a + \cancel{\frac{1}{2}\Omega\sigma^z} + g(a^\dagger + a)\sigma^x$$

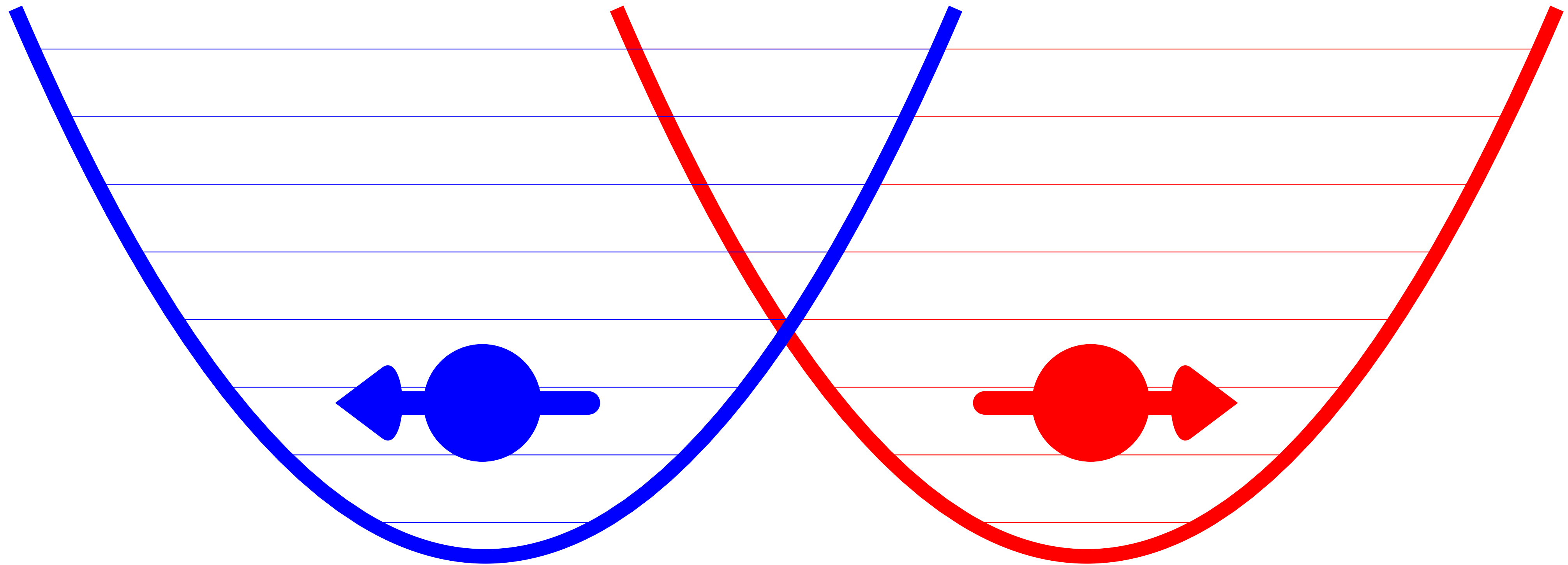
$$|n, \pm\rangle = D(\pm g/\omega) |n\rangle \times |\pm\rangle$$

$$D(z) \equiv \exp(za^\dagger - z^* a)$$

$$E_{n,\pm} \approx \omega n + \mathcal{O}(e^{-g^2/2\omega^2})$$

INFINITELY STRONG-COUPPLING LIMIT

$$g \rightarrow \infty$$



(Franck-Condon effect in polaron physics)

PARITY CONSERVING REPRESENTATION

Hwang & Choi (PRA, 2010; PRB, 2013)

$$H_{\text{Rabi}} = \omega a^\dagger a + \frac{1}{2} \Omega \sigma^z + g(a^\dagger + a) \sigma^x$$

$$\tau^z \equiv \cos(\pi a^\dagger a) \sigma^z$$

$$b \equiv a \sigma^x$$

$$H_{\text{Rabi}} = \omega b^\dagger b - g(b^\dagger + b) + \frac{1}{2} \Omega \cos(\pi b^\dagger b) \tau^z$$

DRESSED STATES

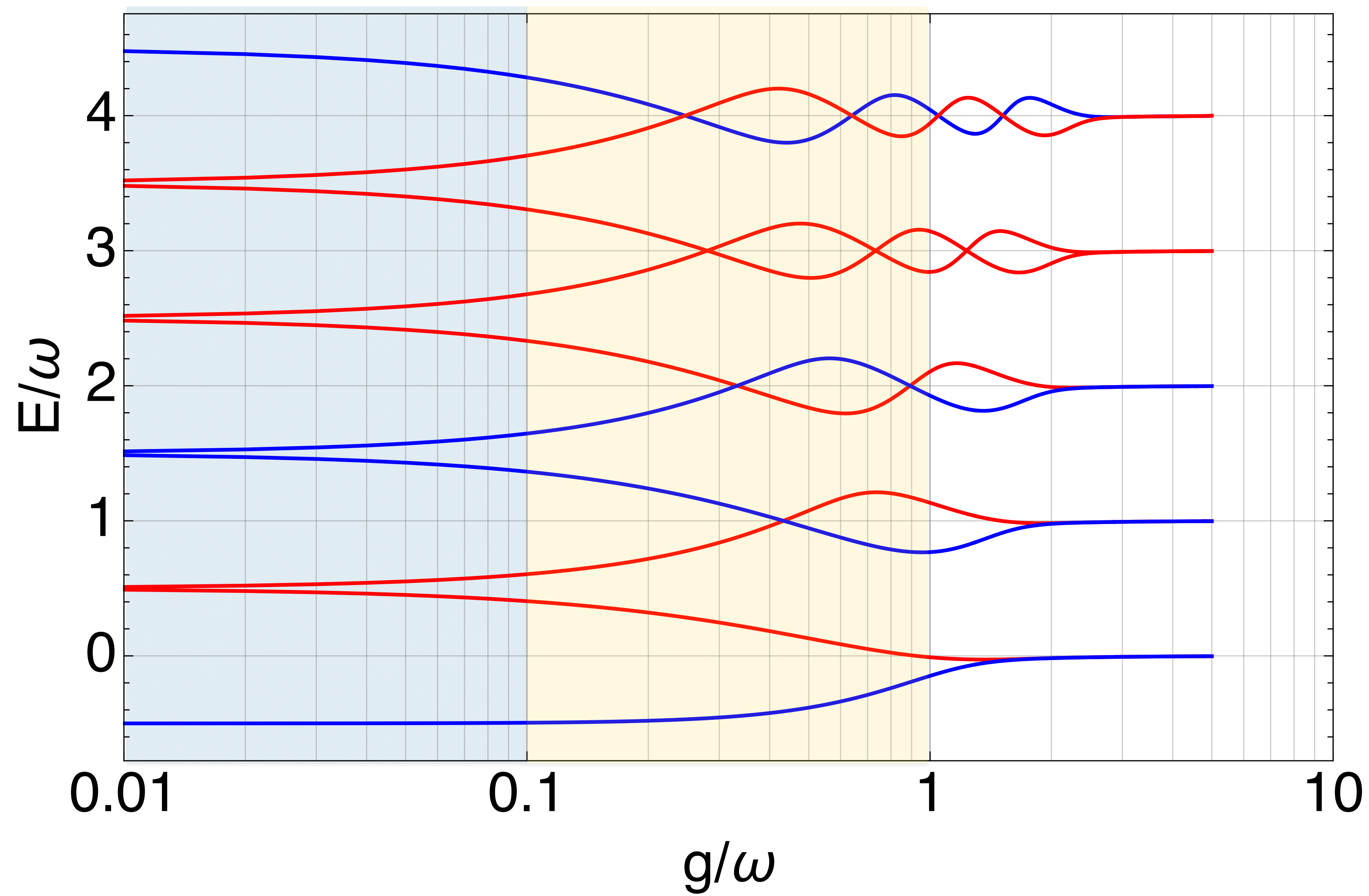
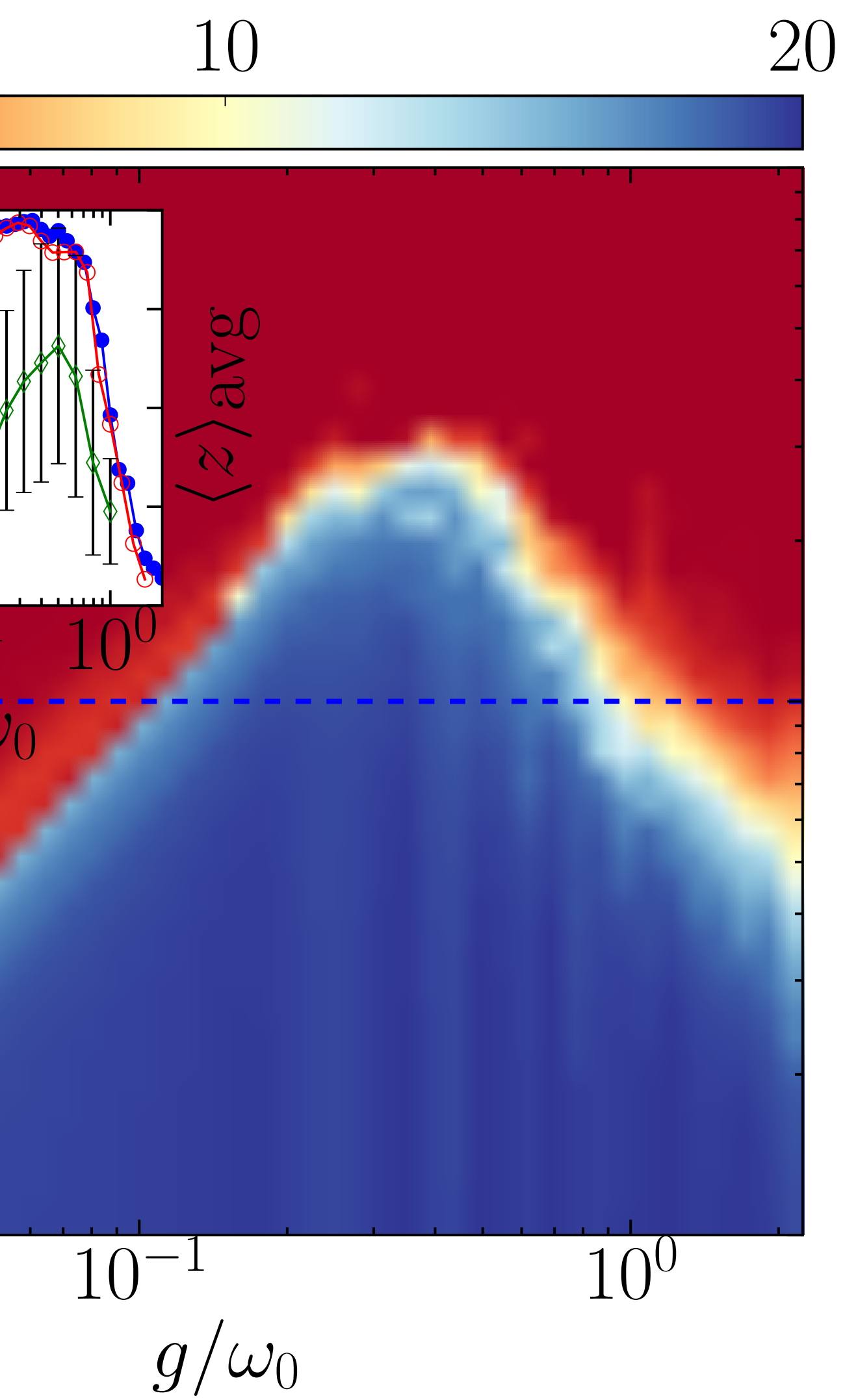
(Generalized Rotating Wave Approximation)

Feranchuk et al. (JPA, 1996); Irish (PRL, 2007)

Hwang & Choi (PRA, 2010; PRB, 2013)

$$|n, e/o\rangle = D_b(g/\omega) |n\rangle_b \times |e/o\rangle$$

$$E_{n,e/o} \approx \omega n + \mathcal{O}(e^{-g^2/2\omega^2}) \quad g \gg \omega$$



DUALITY

DUALITY

Hwang & Choi (PRA, 2010)

Hwang, Puebla, & Plenio (PRL, 2015)

$$H \rightarrow \tilde{H} = D^\dagger(\alpha) H D(\alpha), \quad \alpha \equiv \frac{1}{2} \sqrt{\frac{\Omega}{\omega} \left(\frac{4g^2}{\omega\Omega} - \frac{\omega\Omega}{4g^2} \right)}$$

$$\tilde{H} \approx \omega a^\dagger a + \frac{1}{2} \tilde{\Omega} \tilde{\sigma}^z - \tilde{g} (a^\dagger + a) \tilde{\sigma}^x$$

$$\frac{\tilde{\Omega}}{\Omega} \equiv \frac{4g^2}{\omega\Omega}$$

$$\frac{\tilde{g}}{g} \equiv \frac{\omega\Omega}{4g^2}$$

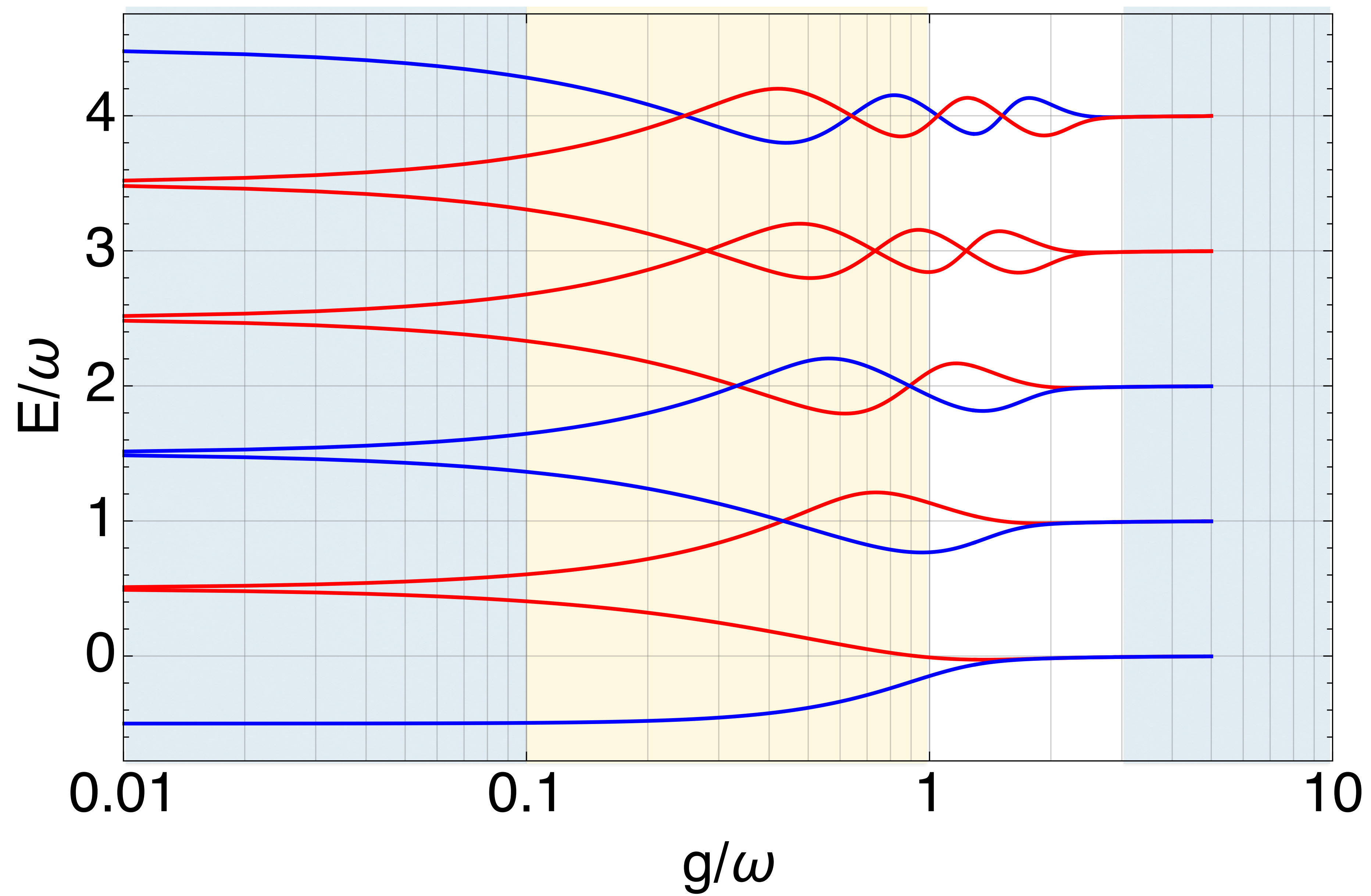
DUALITY!

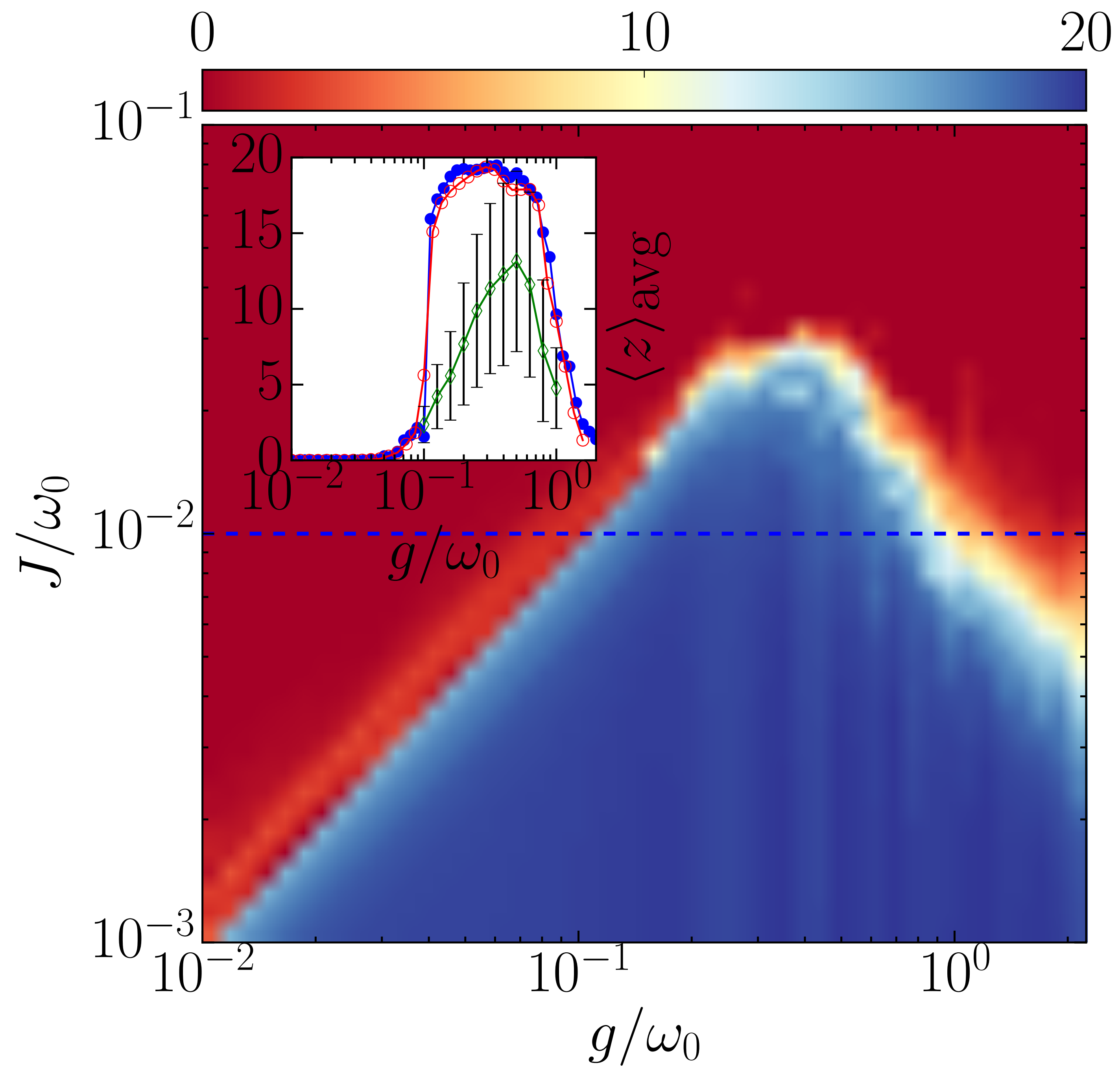
Hwang & Choi (PRA, 2010)

Hwang, Puebla, & Plenio (PRL, 2015)

$$\tilde{H} = \omega a^\dagger a + \frac{1}{2} \tilde{\Omega} \tilde{\sigma}^z - \tilde{g} (a^\dagger + a) \tilde{\sigma}^z$$

$$\frac{2\tilde{g}^2}{\omega\tilde{\Omega}} = \frac{\omega\Omega}{2g^2}$$





THANK YOU!