

Relativistic effects for lanthanide elements

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- Relativistic Hamiltonians
- Spin Separation - Unitary transformations of the Dirac Hamiltonian
- DKH and X2C Hamiltonian
- Hands on exercises: Relativistic calculations for Ce^{3+} ion
- Homework exercise: Relativistic calculations for Pr^{3+} ion

Relativistic effect will affect chemical properties of heavy elements

- atomic ionization potential and electron affinities
- electronic excitation energies (this affects dissociation asymptotes)
- change of bond length and strength when relativistic effects are included

Relativistic effect will affect chemical properties of heavy elements

- electronic states of transition metal, lanthanide, actinide complexes undergo spin orbit splittings
- special relativity effect increases the contraction of ionic radii in the lanthanide series by 10%

Configuration	Term	J	Level (cm^{-1})
$5p^64f$	2F	5/2	0
		7/2	2253
$5p^65d$	2D	3/2	49737
		5/2	52226
$5p^66s$	2S	1/2	86602
$5p^66p$	2P	1/2	122585
		3/2	127292
$5p^656$	2D	5/2	177198
		3/2	178913

Martin et. al (1978)

Relativistic effect will affect chemical properties of heavy elements

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- special relativity effect increases the contraction of ionic radii in the lanthanide series by 10%

Element	Radius Ln^{3+} [pm]
La	103
Ce	102
Pr	99
Nd	98.3
Pm	97
Sm	95.8
Eu	94.7
Gd	93.8
Tb	92.3
Dy	91.2
Ho	90.1
Er	89
Tm	88
Yb	86.8
Lu	86.1

Within the Born–Oppenheimer approximation electronic relativist and non-relativistic Hamiltonians have form

$$\hat{H} = \sum_i \hat{h}(i) + \frac{1}{2} \sum_{i \neq j} \hat{g}(i,j) + V_{NN},$$

- $\hat{h}(i)$ is one electron operator
- $\hat{g}(i,j)$ is electron-electron operator (Breit/Gaunt)
- V_{NN} nucleus-nucleus interaction

See for example mini-review T.Saue, ChemPhysChem 12 3077 (2011)

Dirac equation of an electron in external static field

$$\hat{h}_D = c\boldsymbol{\alpha}\hat{\mathbf{p}} + \hat{V} + E_0\hat{\beta}$$

\hat{V} is an external field (molecular) potential;

$E_0 = m_e c^2$ is a rest mass energy of the electron;

$\beta, \boldsymbol{\alpha} = (\alpha_x, \alpha_y, \alpha_z)$ quantities are 4x4 matrices

$$\beta = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \alpha_x = \begin{pmatrix} 0 & \sigma_x \\ \sigma_x & 0 \end{pmatrix}, \quad \alpha_y = \begin{pmatrix} 0 & \sigma_y \\ \sigma_y & 0 \end{pmatrix}, \quad \alpha_z = \begin{pmatrix} 0 & \sigma_z \\ \sigma_z & 0 \end{pmatrix}$$

expressed in terms of the Pauli matrices

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

The four-component wave function Ψ is formed of two two-component parts: Large Ψ^L , Small Ψ^S

$$\Psi = \begin{pmatrix} \Psi^L \\ \Psi^S \end{pmatrix}, \quad \Psi^L = \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}, \quad \Psi^S = \begin{pmatrix} \psi_3 \\ \psi_4 \end{pmatrix}.$$

\hat{h}_D couples large and small components

$$\begin{pmatrix} \hat{V} & c(\boldsymbol{\sigma} \cdot \mathbf{p}) \\ c(\boldsymbol{\sigma} \cdot \mathbf{p}) & \hat{V} - 2m_e c^2 \end{pmatrix} \begin{pmatrix} \Phi^L \\ \Phi^S \end{pmatrix} = E^+ \begin{pmatrix} \Phi^L \\ \Phi^S \end{pmatrix}$$

Please note that energy is shifted by subtracting rest mass of electron ($E^+ = E - m_e c^2$).

There are many approaches of carrying out unitary transformation \hat{U} which decouples solution of positive and negative energy

$$\hat{U}^\dagger \begin{pmatrix} \hat{h}_{11} & \hat{h}_{12} \\ \hat{h}_{21} & \hat{h}_{22} \end{pmatrix} \hat{U} = \begin{pmatrix} \hat{h}_+ & 0 \\ 0 & \hat{h}_- \end{pmatrix}$$

Review Books:

K. Dyall, K. Faegri "Introduction to Relativistic Quantum Chemistry"
M. Reiher, A. Wolf "Relativistic Quantum Chemistry"

Relativistic electronic-structure calculations employing a two-component no-pair formalism with external-field projection operators

Bernd A. Hess

*Fachbereich 9, Theoretische Chemie, Bergische Universität Gesamthochschule Wuppertal,
Gauss-Strasse 20, 5600 Wuppertal 1, West Germany*

(Received 16 January 1986)

A no-pair formalism employing external-field projection operators correct to second order in the potential is used to calculate the 1s energies of one-electron atoms and ground-state properties of the bromine and silver atoms in the framework of the multireference double-excitation configuration-

- Bernard Hess as the first designed practical algorithm of constructing projection operators corresponding to a particle in the external field of the nucleus (Douglas-Kroll).
- Douglas-Kroll-Hess Hamiltonian is bounded from below
- transformation to two-component Hamiltonian allows for partitioning the relativistic effects into two parts: the scalar relativistic effects and spin-orbit coupling.

Relativistic Hamiltonians implemented in MOLCAS

- arbitrary-order Douglas-Kroll-Hess (DKH) Hamiltonian
- Barysz–Sadlej–Snijders (BSS) Hamiltonian
- one-component exact decoupling (X2C) Hamiltonian

Basis sets to be used in relativistic calculations

- The ANO-RCC basis set (Douglas-Kroll Hamiltonian)
- ANO-R basis set - (X2C)

Spin-free formalism in MOLCAS

"A further approximation is made by neglecting all explicitly spin dependent operators for the variational determination of the wave function. The spin-orbit coupling terms of may be calculated by perturbation theory at a later stage."

B. Hess Phys. Rev. A 33 3742 (1986)

- The CASSCF/RASSCF wave functions are constructed for a specified total spin, which reflect permutation symmetry of the electrons.
- The CASSCF/RASSCF wave functions are not assumed to be eigenstates of \hat{S}_z .
- $(2S + 1)$ CASSCF/RASSCF wave functions for total spin S form representation of all possible spin states within the given multiplet.

Relativistic calculations of Ce³⁺ free ion

CASSCF/CASPT2/RASSI-SO
method

- Ce³⁺ electronic configuration
[Xe]4f¹5d¹6s²
- for calculation use abelian C_i
point group

Configuration	Term	J	Level (cm ⁻¹)	Ce ³⁺
5p ⁶ 4f	² F	5/2	0	
		7/2	2253	
5p ⁶ 5d	² D	3/2	49737	
		5/2	52226	
5p ⁶ 6s	² S	1/2	86602	
5p ⁶ 6p	² P	1/2	122585	
		3/2	127292	
5p ⁶ 56	² D	5/2	177198	
		3/2	178913	

Martin et. al (1978) NIST

Relativistic calculations of Ce³⁺ free ion

&GATEWAY

Title

Cerium

SYmmetry

xyz

Basis set

Ce.ano-rcc.Roos.25s22p15d11f4g2h.9s8p5d4f3g2h

Ce 0 0 0

end of basis

AUXShow

AMFI

ANGMOM=0.0 0.0 0.0

&SEWARD

- Ce.ano-rcc.Roos will activate DKH
- AMFI Atomic mean field integrals

Input file:

Cerium3.casscf_rassi.input

Relativistic calculations of Ce³⁺ free ion

&SCF &END

Title

Cerium3+ atom

Occupied

15 19

occnumbers

*1 2 3 4 5 6 7 8 9 10

2 2 2 2 2 2 2 2 2 2

2 2 2 2 2

2 2 2 2 2 2 2 2 2 2

2 2

0.1428571428571428 0.1428571428571428

0.1428571428571428 0.1428571428571428

0.1428571428571428 0.1428571428571428

0.1428571428571428

Iterations

40

End of Input

C_i	E	i	Orbital
A _g	+1	+1	s, d
A _u	+1	-1	p, f

A_g: $1s^2 2s^2 2p^6 3s^2 3p^6 3d^10 4s^2 4p^6 4d^{10}$

1+1+1+1+5+1+5=15

A_u: $2s^2 2p^6 3s^2 3p^6 3d^10 4s^2 4p^6 4d^10 5s^2 5p^6 5d^1 6s^2 6p^6 6d^1 7s^2 7p^6 7d^1 8s^2 8p^6 8d^1$

3+3+3+3+7=19

Relativistic calculations of Ce³⁺ free ion

```
&RASSCF &END
Title
Cerium 3+    CAS2 (4f5d6s)1    2Au
Symmetry
2
Spin
2
Nactel
1   0   0
Frozen
0   0
Inactive
15  12
Ras2
6   7
CiRoot
7  7
1 2 3 4 5 6 7
1 1 1 1 1 1 1
LumOrb
End of Input
>>COPY $Project.JobIph JOB001
```

Active space: one electron on 13 orbitals
 2A_u irrep
 $(2 \cdot 3 + 1) = 7$ degenerate 4f states L=3
command COPY saves CASSCF wave functions for further calculations

Relativistic calculations of Ce³⁺ free ion

```
&RASSCF &END
Title
Cerium 3+ ion,  CAS2:4f5d6s  States 2Ag
Symmetry
1
Spin
2
Nactel
1   0   0
Frozen
0   0
Inactive
15  12
Ras2
6   7
CiRoot
6  6
1 2 3 4 5 6
1 1 1 1 1 1
LumOrb
End of Input
>>COPY $Project.JobIph JOB002
```

Active space: one electron on 13 orbitals

2A_g irrep

$(2 \cdot 2 + 1) + (2 \cdot 0 + 1) = 6$ degenerate 5d
states L=2 plus one 6s

command COPY saves CASSCF wave
functions for further calculations

Relativistic calculations of Ce³⁺ free ion

```
&RASSI &END  
CIprint  
NROF JOBIPHS  
2 7 6  
1 2 3 4 5 6 7  
1 2 3 4 5 6  
Spin  
End of input
```

RASSI-SO calculations,
The Joblph files copied as JOB001,
JOB002 are used. The Nr of Joblphs
keyword is followed by: The number of
Joblph files (2), The number of states to
pick from each of JOB00n files (7,6), the
serial numbers of these states for the first
file (1,2,3,4,5,6,7), the serial numbers of
these states for the second file (1,2,3,4,5,6)

Results

Eigenvalues of complex Hamiltonian:

(Shifted by EMIN (a.u.) = -8852.0496747242)

S0 State	Relative EMIN(au)	Rel lowest level(eV)	D:o, cm**(-1)	J-value
1	-0.0062960408	0.0000000000	0.0000	2.5
2	-0.0062960408	0.0000000000	0.0000	2.5
3	-0.0062960408	0.0000000000	0.0000	2.5
4	-0.0062960408	0.0000000000	0.0000	2.5
5	-0.0062960408	0.0000000000	0.0000	2.5
6	-0.0062960408	0.0000000000	0.0000	2.5
7	0.0047220306	0.2998169948	2418.1871	3.5
8	0.0047220306	0.2998169948	2418.1871	3.5
9	0.0047220306	0.2998169948	2418.1871	3.5
10	0.0047220306	0.2998169948	2418.1871	3.5
11	0.0047220306	0.2998169948	2418.1871	3.5
12	0.0047220306	0.2998169948	2418.1871	3.5
13	0.0047220306	0.2998169948	2418.1871	3.5
14	0.0047220306	0.2998169948	2418.1871	3.5
15	0.1964870207	5.5180082094	44505.7377	1.5
16	0.1964870207	5.5180082094	44505.7377	1.5
17	0.1964870207	5.5180082094	44505.7377	1.5

Exercise 1: MS-CASPT2 and broken degeneracies of J energies

input file: Cerium3.caspt2_rassi.input

	S0	State	...	D:o, cm**(-1)	J-value
&CASPT2 &END					
Title	1	...		0.0000	2.5
Cerium 3+ ion 2Au ungerade states	2	...		0.0000	2.5
MultiState=7	3	...		17.7057	2.5
1 2 3 4 5 6 7	4	...		17.7057	2.5
End of input	5	...		21.8429	2.5
>>COPY \$Project.JobMix JOB001	6	...		21.8429	2.5
... long input ...	7	...		2415.7169	3.5
&RASSI &END	8	...		2415.7169	3.5
CIprint	9	...		2420.3363	3.5
NROF JOBIPHS	10	...		2420.3363	3.5
2 7 6	11	...		2441.4861	3.5
1 2 3 4 5 6 7	12	...		2441.4861	3.5
1 2 3 4 5 6	13	...		2449.1004	3.5
Spin	14	...		2449.1004	3.5
ejob	15	...		44918.2725	1.5
End of input	16	...		44918.2725	1.5
	17	...		44921.2176	1.5
	18	...		44921.2176	1.5
	19	...		47249.9050	2.5

One can see that that $(2J + 1)$ components of J level are not degenerated. This error is introduced by MS-CASPT2 method.

AVERAGED ENERGIES (2J+1)			
:: RASSI State	1	Total energy:	-8852.27226716 4f -8852.272183570
:: RASSI State	2	Total energy:	-8852.27225926 4f -8852.272183570
:: RASSI State	3	Total energy:	-8852.27225182 4f -8852.272183570
:: RASSI State	4	Total energy:	-8852.27224297 4f -8852.272183570
:: RASSI State	5	Total energy:	-8852.27220762 4f -8852.272183570
:: RASSI State	6	Total energy:	-8852.27205902 4f -8852.272183570
:: RASSI State	7	Total energy:	-8852.27199714 4f -8852.272183570
:: RASSI State	8	Total energy:	-8852.06751354 5d -8852.067494002
:: RASSI State	9	Total energy:	-8852.06749892 5d -8852.067494002
:: RASSI State	10	Total energy:	-8852.06748936 5d -8852.067494002
:: RASSI State	11	Total energy:	-8852.06748638 5d -8852.067494002
:: RASSI State	12	Total energy:	-8852.06748181 5d -8852.067494002
:: RASSI State	13	Total energy:	-8851.90695416 6s -8851.906954160

```
&RASSI &END
CIprint
NROF JOBIPHS
2 7 6
1 2 3 4 5 6 7
1 2 3 4 5 6
Spin
HDIAG
-8852.272183570
-8852.272183570
-8852.272183570
-8852.272183570
-8852.272183570
-8852.272183570
-8852.272183570
-8852.067494002
-8852.067494002
-8852.067494002
-8852.067494002
-8851.906954160
End of input
```

HDIAG The values give an energy for each wave function, to replace the diagonal elements of the Hamiltonian matrix.
Please check input file:
`Cerium3.hdiag_rassi.input`

