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### Crystal field theory Point groups

- Symmetry lowered by "crystal field":  $G \subset O(3)$
- Discrete set of symmetry operations remain
- Conjugacy classes:  $R, S \in \mathcal{C}$  if  $\exists T \in G : R = T^{-1}ST$
- Representations:

$$P_R\psi_j = \sum_{i=1}^{l_{\Gamma}} [\Gamma(R)]_{ij}\psi_i$$

• Characters:

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$$\chi_{\Gamma}(R) = \sum_{i=1}^{l_{\Gamma}} [\Gamma(R)]_{ii} = \operatorname{Tr}[\Gamma(R)]$$

Symmetry

### Crystal field theory Point groups

• Example: character table of  $O_h = O \otimes \{E, i\}$ 

$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	$O_h$	E	$8C_3$	$6C'_2$	$6C_4$	$3C_2$	i	$8S_6$	$6\sigma_d$	$6S_4$	$3\sigma_h$
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$A_{1g}$	1	1	1	1	1	1	1	1	1	1
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$A_{2g}$	1	1	-1	-1	1	1	1	-1	-1	1
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$E_{g}$	2	-1	0	0	2	2	-1	0	0	2
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$T_{1g}$	3	0	-1	1	-1	3	0	-1	1	-1
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$T_{2g}$	3	0	1	-1	-1	3	0	1	-1	-1
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$A_{1u}$	1	1	1	1	1	-1	-1	-1	-1	-1
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$A_{2u}$	1	1	-1	-1	1	-1	-1	1	1	-1
$T_{1u}$ 3 0 -1 1 -1 -3 0 1 -1 1	$E_{u}$	2	-1	0	0	2	-2	1	0	0	-2
	$T_{1u}$	3	0	-1	1	-1	-3	0	1	-1	1
$T_{2u}$ 3 0 1 -1 -1 -3 0 -1 1 1	$T_{2u}$	3	0	1	-1	-1	-3	0	-1	1	1

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$O_h$	E	$8C_3$	$6C'_2$	$6C_4$	$3C_2$	i	$8S_6$	$6\sigma_d$	$6S_4$	$3\sigma_h$
$A_{1g}$	1	1	1	1	1	1	1	1	1	1
$A_{2g}$	1	1	-1	-1	1	1	1	-1	-1	1
$E_g$	2	-1	0	0	2	2	-1	0	0	2
$T_{1g}$	3	0	-1	1	-1	3	0	-1	1	-1
$T_{2g}$	3	0	1	-1	-1	3	0	1	-1	-1
$A_{1u}$	1	1	1	1	1	-1	-1	-1	-1	-1
$A_{2u}$	1	1	-1	-1	1	-1	-1	1	1	-1
$E_u$	2	-1	0	0	2	-2	1	0	0	-2
$T_{1u}$	3	0	-1	1	-1	-3	0	1	-1	1
$T_{2u}$	3	0	1	-1	-1	-3	0	-1	1	1

# Crystal field theory

Point groups

- Symmetry lowered by "crystal field":  $O(3) \supset G$
- IRs of O(3) (i.e. atomic states) reducible in G:

$$\mathcal{D}^{(\mathcal{I})} = \bigoplus_{j=1}^{r} a_j \Gamma^{(j)} = a_1 \Gamma^{(1)} \oplus a_2 \Gamma^{(2)} \oplus \ldots \oplus a_r \Gamma^{(r)}$$

 $(\mathcal{I} \text{ angular momentum of interest}, r \text{ number of IR for group } G)$ 

• Simply calculable thanks to character table:

$$a_j = \frac{1}{h} \sum_{i=1}^r N_i \underbrace{\chi(\mathcal{C}_i)}_{O(3)} \underbrace{\left[\chi^{(j)}(\mathcal{C}_i)\right]^*}_G$$

 $(N_i \text{ number of symm. operations in class } C_i; h = \sum_{i=1}^r N_i = \text{order of } G)$ 

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### Crystal field theory Point groups

٠	Example:	additional	classes	and IRs	of $O_h^*$ :
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$O_h^*$	E	R	$8C_3$	$8RC_3$	$6C_2' + 6RC_2'$	$6C_4$	$6RC_4$	$3C_2 + 3RC_2$	i	• • •
$E_{1/2g}$	2	-2	1	-1	0	$\sqrt{2}$	$-\sqrt{2}$	0	2	• • •
$E_{5/2g}$	2	-2	1	-1	0	$-\sqrt{2}$	$\sqrt{2}$	0	2	
$F_{3/2g}$	4	-4	-1	1	0	0	0	0	4	•••
$E_{1/2u}$	2	-2	1	-1	0	$\sqrt{2}$	$-\sqrt{2}$	0	-2	• • •
$E_{5/2u}$	2	-2	1	-1	0	$-\sqrt{2}$	$\sqrt{2}$	0	-2	• • •
$F_{3/2u}$	4	-4	-1	1	0	0	0	0	-4	•••



### Crystal field theory Symmetry-adapted basis states

• Different choices rationalized from "coupling schemes"

	$\hat{H} =$	$\hat{H}_0' + \hat{H}_1' + \hat{H}_2 +$	$\hat{H}_3$
	$\operatorname{term}$	$\operatorname{multiplet}$	crystal field
	splitting $(RS)$	splitting $(SO)$	$\operatorname{splitting}$
3d	1-10 eV	$0.1 \ \mathrm{eV}$	1  eV
4d	1-10  eV	0.1-1  eV	$1 \mathrm{~eV}$
5d	1-10  eV	$1  \mathrm{eV}$	$1 \mathrm{~eV}$
4f	1-10  eV	$0.1-1 \ \mathrm{eV}$	$0.01 \ \mathrm{eV}$
5f	1-10  eV	$1  \mathrm{eV}$	$0.01\text{-}0.1~\mathrm{eV}$
5p	1-10  eV	$1  \mathrm{eV}$	$1 \mathrm{eV}$
6p	1-10  eV	$1  \mathrm{eV}$	$1 \mathrm{~eV}$

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### Crystal field theory

Symmetry-adapted basis states: weak field

- Natural for  $4f^N$  (lanthanides) and  $5f^N$  (heavy actinides) configurations
- 2J + 1 degenerate atomic multiplet  ${}^{2S+1}L_J$  split by CF:

$$\mathcal{D}^{(J)} = \bigoplus_{j=1}^{r} a_j \Gamma^{(j)}$$

• Reduction coefficients:

$$|\alpha SLJ; \ a\Gamma\gamma\rangle = \sum_{M=-J}^{J} \langle JM| Ja\Gamma\gamma\rangle |\alpha SLJM\rangle$$

• Example:  $Pr^{3+}$ ;  $4f^2$ ;  $^{3}H_4$ :

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$${}^{3}H_{4} = A_{1g} \oplus E_{g} \oplus T_{1g} \oplus T_{2g} \qquad (O_{h})$$

Orystal field theory Symmetry-adapted basis states: intermediate field • Natural for  $3d^N$  (transition metals) configurations • (2L+1)(2S+1) degenerate atomic term  ${}^{2S+1}L$  split by CF:  $D^{(L)} = \bigoplus_{i=1}^r a_i \Gamma^{(i)}$ • Example:  $\operatorname{Cr}^{3+}$ ;  $3d^3$ ;  ${}^4F$ :  ${}^4F = {}^4A_{2g} \oplus {}^4T_{1g} \oplus {}^4T_{2g}$  ( $O_h$ )

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### Crystal field theory Symmetry-adapted basis states: intermediate field

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• Spin degeneracy further reduced through spin-orbit coupling:

$$\mathcal{D}^{(S)} = \bigoplus_{j=1}^{\prime} a_j \Gamma^{(j)}$$

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 ${}^{4}A_{2g}$ 

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• CF multiplets obtained from direct product (spin-orbit interaction):

$$\Gamma^{(i)} \to \Gamma^{(i)} \otimes \mathcal{D}^{(S)} = \bigoplus_{j=1}^{r} a_j \left( \Gamma^{(i)} \otimes \Gamma^{(j)} \right) = \bigoplus_{j=1}^{r} \bigoplus_{k=1}^{r} a_j a_k \Gamma^{(k)}$$

• Clebsch-Gordan and reduction coefficients:  $|\alpha(Sa_s\Gamma^{(j)}, La_L\Gamma^{(i)})a\Gamma^{(k)}\gamma^{(k)}\rangle = \sum_{\substack{M_SM_L\gamma^{(i)}\gamma^{(j)}}} \langle SM_S | a_s\Gamma^{(j)}\gamma^{(j)} \rangle \langle LM_L | a_L\Gamma^{(i)}\gamma^{(i)} \rangle \langle \Gamma^{(i)}\gamma^{(i)}\Gamma^{(j)}\gamma^{(j)} | a\Gamma^{(k)}\gamma^{(k)} \rangle | \alpha SM_SLM_L \rangle$ 

Symmetry

# Crystal field theory

Symmetry-adapted basis states: intermediate field

• Spin degeneracy further reduced through spin-orbit coupling:

$$\mathcal{D}^{(S)} = \bigoplus_{j=1}^r a_j \Gamma^{(j)}$$

• CF multiplets obtained from direct product (spin-orbit interaction):

$$\Gamma^{(i)} \to \Gamma^{(i)} \otimes \mathcal{D}^{(S)} = \bigoplus_{j=1}^{r} a_j \left( \Gamma^{(i)} \otimes \Gamma^{(j)} \right) = \bigoplus_{j=1}^{r} \bigoplus_{k=1}^{r} a_j a_k \Gamma^{(k)}$$

• Example:  $\operatorname{Cr}^{3+}$ ;  $3d^3$ ;  ${}^4F$ ;  ${}^4T_{1g}$ :

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$$\mathcal{D}^{3/2} = F_{3/2g}$$
  
$$T_{1g} \otimes F_{3/2g} = E_{1/2g} \oplus E_{5/2g} \oplus 2F_{3/2g}$$

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	6	5	4	3	2	1	0
1		$(\overset{+}{3},\overset{+}{2})$	$(\overset{+}{3},\overset{+}{1})$	$(\overset{+}{3},\overset{+}{0}),(\overset{+}{2},\overset{+}{1})$	$(\overset{+}{3},\overset{+}{-1}),(\overset{+}{2},\overset{+}{0})$	$(\overset{+}{3},\overset{+}{-2}),(\overset{+}{2},\overset{+}{-1}),(\overset{+}{1},\overset{+}{0})$	(3,-3), (2,-2), (1,-1)
		$^{3}H$	$^{3}H$	${}^{3}H$ ${}^{3}F$	$^{3}H$ $^{3}F$	${}^{3}H$ ${}^{3}F$ ${}^{3}P$	${}^{3}H$ ${}^{3}F$ ${}^{3}P$
0	$({}^+3,{}^-3)$	$(\overset{+}{3},\overset{-}{2})$ $(\overset{-}{3},\overset{+}{2})$	$(\overset{+}{3},\overset{-}{1})\\(\overset{-}{3},\overset{+}{1})\\(\overset{+}{2},\overset{-}{2})$	$(\overset{+}{3},\overset{-}{0}), (\overset{-}{3},\overset{+}{0})$ $(\overset{+}{2},\overset{-}{1}), (\overset{-}{2},\overset{+}{1})$	$ \begin{array}{c} \stackrel{+}{(3,-1)}, \stackrel{-}{(3,-1)} \\ \stackrel{+}{(2,0)}, \stackrel{-}{(2,0)} \\ \stackrel{+}{(1,1)} \end{array} $	$ \stackrel{+}{\overset{-}(3,-2)}, \stackrel{-}{\overset{-}(3,-2)}, \stackrel{+}{\overset{-}(2,-1)}, \stackrel{+}{\overset{-}(2,-1)} \\ \stackrel{-}{\overset{-}(2,-1)}, \stackrel{+}{\overset{+}(1,0)}, \stackrel{-}{\overset{-}(1,0)} $	$ \begin{array}{c} (\overset{+}{3},\overset{-}{-3}),(\overset{-}{3},\overset{-}{-3}),(\overset{+}{2},\overset{-}{-2})\\ (\overset{+}{2},\overset{+}{-2}),(\overset{+}{1},\overset{-}{-1}),(\overset{+}{1},\overset{+}{-1})\\ (\overset{+}{0},\overset{-}{0}) \end{array} $
	$^{1}I$	${}^{1}I$ ${}^{3}H$	${}^{1}I  {}^{1}G  {}^{3}H$	${}^{1}I  {}^{1}G                   $	${}^{1}I  {}^{1}G  {}^{1}D \ {}^{3}H  {}^{3}F$	${}^{1}I {}^{1}G {}^{1}D \\ {}^{3}H {}^{3}F {}^{3}P$	${}^{1}I  {}^{1}G  {}^{1}D  {}^{1}S \ {}^{3}H  {}^{3}F  {}^{3}P$

# Solution: Configurations, terms and multiplets for $Pr^{3+}$

•  $4f^2$  configuration

- Number of Slater determinants:  $\frac{14!}{2!12!} = 91$
- Pauli's exclusion principle must be respected: use  $M_L, M_S$  table
- In summary:
  - Triplets:  ${}^{3}H$ ,  ${}^{3}F$ ,  ${}^{3}P$
  - Singlets:  ${}^{1}I$ ,  ${}^{1}G$ ,  ${}^{1}D$ ,  ${}^{1}S$

• 4f5d configuration

- Number of Slater determinants:  $14 \times 10 = 140$
- Pauli's exclusion principle is automatically respected: straightforward coupling

  - $S = |s_1 s_2|, \dots, s_1 + s_2 = 0, 1$  or  $1/2 \otimes 1/2 = 0 \oplus 1$   $L = |\ell_1 \ell_2|, \dots, \ell_1 + \ell_2 = 1, 2, 3, 4, 5$  or  $3 \otimes 2 = 1 \oplus 2 \oplus 3 \oplus 4 \oplus 5$
- All combinations:
  - Triplets:  ${}^{3}H$ ,  ${}^{3}G$ ,  ${}^{3}F$ ,  ${}^{3}D$ ,  ${}^{3}P$
  - Singlets:  ${}^{1}H$ ,  ${}^{1}G$ ,  ${}^{1}F$ ,  ${}^{1}D$ ,  ${}^{1}P$



• Singlets do not split (S = 0, J = L)

• ...

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### Solution: Configurations, terms and multiplets for Pr<sup>3+</sup> Energy level scheme

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- Compare with experiment / empirical assignment
- Compare with *ab initio* wave function calculations (see also later lectures):



## Exercise: Crystal field levels $Pr^{3+}$ in YAG

Assume that the  $Pr^{3+}$  ion is incorporated on an  $Y^{3+}$  site of a YAG crystal.

- Predict how the atomic  ${}^{2S+1}L$  terms split in the crystal field (with  $D_2$  symmetry). How many occurrences of every spin-symmetry combination  ${}^{2S+1}\Gamma$  do you expect for both considered electron configurations?
- Oberive how every <sup>2S+1</sup>Γ is further split by spin-orbit coupling. What are the corresponding symmetry labels?

$D_2^*$	E	R	$C_2 + RC_2 \ (z)$	$C_2 + RC_2 (y)$	$C_2 + RC_2 \ (x)$	
A	1	1	1	1	1	
$B_1$	1	1	1	-1	-1	
$B_2$	1	1	-1	1	-1	
$B_3$	1	1	-1	-1	1	
$E_{1/2}$	2	-2	0	0	0	
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### Solution: Crystal field levels $Pr^{3+}$ in YAG

- Descent from spherical symmetry (L part):
  - S term:  $\mathcal{D}^{(0)} = A$
  - P term:  $\mathcal{D}^{(1)} = B_1 \oplus B_2 \oplus B_3$
  - *D* term:  $\mathcal{D}^{(2)} = 2A \oplus B_1 \oplus B_2 \oplus B_3$
  - F term:  $\mathcal{D}^{(3)} = A \oplus 2B_1 \oplus 2B_2 \oplus 2B_3$
  - G term:  $\mathcal{D}^{(4)} = 3A \oplus 2B_1 \oplus 2B_2 \oplus 2B_3$
  - *H* term:  $\mathcal{D}^{(5)} = 2A \oplus 3B_1 \oplus 3B_2 \oplus 3B_3$
  - *I* term:  $\mathcal{D}^{(6)} = 4A \oplus 3B_1 \oplus 3B_2 \oplus 3B_3$

# Solution: Crystal field levels $Pr^{3+}$ in YAG

- Descent from spherical symmetry (L part)
- $4f^2$  configuration:  $3^3A$ ,  $6^3B_1$ ,  $6^3B_2$ ,  $6^3B_3$ ,  $10^1A$ ,  $6^1B_1$ ,  $6^1B_2$ ,  $6^1B_3$

		$^{3}A$	${}^{3}B_{1}$	${}^{3}B_{2}$	${}^{3}B_{3}$	$^{1}A$	${}^{1}B_{1}$	${}^{1}B_{2}$	$^{1}B_{3}$	
	$^{3}H$	2	3	3	3					
	${}^{3}F$	1	2	2	2					
	$^{3}P$	0	1	1	1					
	$^{1}I$					4	3	3	3	
	${}^{1}G$					3	2	2	2	
	$^{1}D$					2	1	1	1	
	$^{1}S$					1	0	0	0	
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# Solution: Crystal field levels $Pr^{3+}$ in YAG

- Descent from spherical symmetry (L part)
- $4f^2$  configuration:  $3^3A$ ,  $6^3B_1$ ,  $6^3B_2$ ,  $6^3B_3$ ,  $10^1A$ ,  $6^1B_1$ ,  $6^1B_2$ ,  $6^1B_3$
- 4f5d configuration:  $8^3A$ ,  $9^3B_1$ ,  $9^3B_2$ ,  $9^3B_3$ ,  $8^1A$ ,  $9^1B_1$ ,  $9^1B_2$ ,  $9^1B_3$

	$^{2S+1}A$	${}^{2S+1}B_1$	${}^{2S+1}B_2$	${}^{2S+1}B_3$
$^{2S+1}H$	2	3	3	3
$^{2S+1}G$	3	2	2	2
$^{2S+1}F$	1	2	2	2
${}^{2S+1}D$	2	1	1	1
$^{2S+1}P$	0	1	1	1



## Solution: Crystal field levels $Pr^{3+}$ in YAG

- Descent from spherical symmetry (L part)
- $4f^2$  configuration:  $3^3A$ ,  $6^3B_1$ ,  $6^3B_2$ ,  $6^3B_3$ ,  $10^1A$ ,  $6^1B_1$ ,  $6^1B_2$ ,  $6^1B_3$
- 4f5d configuration:  $8^{3}A$ ,  $9^{3}B_{1}$ ,  $9^{3}B_{2}$ ,  $9^{3}B_{3}$ ,  $8^{1}A$ ,  $9^{1}B_{1}$ ,  $9^{1}B_{2}$ ,  $9^{1}B_{3}$
- Descent from spherical symmetry (S part)

	A	$B_1$	$B_2$	$B_3$
$^{3}A$		1	1	1
${}^{3}B_{1}$	1		1	1
${}^{3}B_{2}$	1	1		1
${}^{3}B_{3}$	1	1	1	
$^{1}A$	1			
${}^{1}B_{1}$		1		
${}^{1}B_{2}$			1	
${}^{1}B_{3}$				1



- Descent from spherical symmetry (L part)
- $4f^2$  configuration:  $3^3A$ ,  $6^3B_1$ ,  $6^3B_2$ ,  $6^3B_3$ ,  $10^1A$ ,  $6^1B_1$ ,  $6^1B_2$ ,  $6^1B_3$
- 4f5d configuration:  $8^{3}A$ ,  $9^{3}B_{1}$ ,  $9^{3}B_{2}$ ,  $9^{3}B_{3}$ ,  $8^{1}A$ ,  $9^{1}B_{1}$ ,  $9^{1}B_{2}$ ,  $9^{1}B_{3}$
- Descent from spherical symmetry (S part)
- $4f^2$  configuration: 28A,  $21B_1$ ,  $21B_2$ ,  $21B_3$  (91 SO states in total)
- 4f5d configuration: 35A,  $35B_1$ ,  $35B_2$ ,  $35B_3$  (140 SO states in total)

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