







# Radiative transitions

Light-matter interaction

- QED Hamiltonian:  $\hat{H}_{\rm tot} = \hat{H}_{\rm mat} + \hat{H}_{\rm rad} + \hat{H}_{\rm int}$
- Electron in EM field:  $\vec{p_i} \rightarrow \vec{p_i} + e\vec{A}$ 
  - $\vec{A}(\vec{r},t) = \sum_{\vec{k},\sigma} A_0(k) \vec{\epsilon}_{\vec{k}\sigma} \left( a_{\vec{k}\sigma} e^{i\vec{k}\cdot\vec{r}} + a_{\vec{k}\sigma}^* e^{-i\vec{k}\cdot\vec{r}} \right)$ •  $\vec{\nabla}_{\vec{l}} \cdot \vec{A} = 0$

• 
$$\mathrm{e}^{\mathrm{i}\vec{k}\cdot\vec{r}} \approx 1 + \mathrm{i}\vec{k}\cdot\bar{r}$$

• After a bit of algebra:

$$\hat{H}_{\text{int}} = \sum_{i=1}^{N} \left( \underbrace{\underline{e\vec{r_i} \cdot \vec{E_0}(k)}_{\text{E1}}}_{\text{E1}} + \underbrace{\frac{e}{2m_e} \left(\hat{\vec{\ell_i}} + g_e \hat{\vec{s}_i}\right) \cdot \vec{B_0}(k)}_{\text{M1}} + \underbrace{\frac{1}{2} e\vec{r_i} \otimes \vec{r_i} : \vec{k} \otimes \vec{E_0}(k)}_{\text{E2}} \right)$$

Electronic spectra

• Relative strengths E1 : M1 : E2  $\approx 1 : 10^{-5} : 10^{-6}$ 

### Radiative transitions Experimental observables

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• Absorption spectrum:

$$S_{\rm abs}(\omega) = C\omega \sum_{f} \left| \langle f | \, \hat{\vec{D}_e} \, | i \rangle \right|^2 \delta(E_f - E_i - \hbar\omega)$$

• Emission spectrum:

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$$S_{\rm em}(\omega) = C\omega^3 \sum_{f} \left| \langle f | \hat{\vec{D}}_e | i \rangle \right|^2 \delta(E_f - E_i - \hbar\omega)$$

with  $\hat{\vec{D}}_e = -e \sum_{i=1}^N \vec{r_i}$  for E1 transitions.

Electronic spectra

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## Electron-vibrational transitions

• Born-Oppenheimer wave functions:

$$|\Psi_1^{(\text{tot})}\rangle = |\Psi_1\rangle |\zeta_{1n_1}\rangle \quad \text{and} \quad |\Psi_2^{(\text{tot})}\rangle = |\Psi_2\rangle |\zeta_{2n_2}\rangle$$

• Vibrational wave function:

$$\zeta_{\alpha n_{\alpha}}(Q_{\alpha}) = \prod_{i=1}^{3\nu-6} \zeta_{\alpha i n_{\alpha i}}(Q_i)$$

• Transition dipole operator:

$$\hat{\vec{D}} = \underbrace{-e\sum_{i=1}^{N}\vec{r_i}}_{\vec{D}_{\rm e}} + \underbrace{e\sum_{j=1}^{\nu}Z_j\vec{R_j}}_{\vec{D}_{\rm N}}$$

• Matrix element:

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$$\langle \Psi_2^{(\text{tot})} | \, \hat{\vec{D}} \, | \Psi_1^{(\text{tot})} \rangle = \langle \zeta_{2n_2} | \, \langle \Psi_2 | \, \hat{\vec{D}}_e \, | \Psi_1 \rangle \, | \zeta_{1n_1} \rangle + \underbrace{\langle \Psi_2 | \Psi_1 \rangle}_{=0} \langle \zeta_{2n_2} | \, \hat{\vec{D}}_N \, | \zeta_{1n_1} \rangle.$$

Electronic spectra

# Electron-vibrational transitions

• FC approximation:

$$\hat{\vec{D}}_{e}(Q) = \hat{\vec{D}}_{e}(Q_{1}^{(0)}) + \sum_{i=1}^{3\nu-6} \left(\frac{\partial \hat{\vec{D}}_{e}}{\partial Q_{i}}\right)_{Q_{1}^{(0)}} \left(Q_{i} - Q_{1,i}^{(0)}\right) \\ + \frac{1}{2} \sum_{i=1}^{3\nu-6} \sum_{j=1}^{3\nu-6} \left(\frac{\partial^{2} \hat{\vec{D}}_{e}}{\partial Q_{i} \partial Q_{j}}\right)_{Q_{1}^{(0)}} \left(Q_{i} - Q_{1,i}^{(0)}\right) \left(Q_{j} - Q_{1,j}^{(0)}\right) + \mathcal{O}(Q^{3})$$

• Matrix element:

$$\langle \Psi_2^{(\text{tot})} | \, \hat{\vec{D}}_{e} \, | \Psi_1^{(\text{tot})} \rangle = \underbrace{\langle \Psi_2 | \, \hat{\vec{D}}_{e}(0) \, | \Psi_1 \rangle}_{\text{purely electronic}} \underbrace{\langle \zeta_{2n_2}(Q) | \, \zeta_{1n_1}(Q) \rangle}_{\text{FC integrals}}$$

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## Electron-vibrational transitions

• FC approximation:

$$\hat{\vec{D}}_{e}(Q) = \hat{\vec{D}}_{e}(Q_{1}^{(0)}) + \sum_{i=1}^{3\nu-6} \left( \frac{\partial \hat{\vec{D}}_{e}}{\partial Q_{i}} \right)_{Q_{1}^{(0)}} \left( Q_{i} - Q_{1,i}^{(0)} \right)$$

$$+ \frac{1}{2} \sum_{i=1}^{3\nu-6} \sum_{j=1}^{3\nu-6} \left( \frac{\partial^{2} \hat{\vec{D}}_{e}}{\partial Q_{i} \partial Q_{j}} \right)_{Q_{1}^{(0)}} \left( Q_{i} - Q_{1,i}^{(0)} \right) \left( Q_{j} - Q_{1,j}^{(0)} \right) + \mathcal{O}(Q^{3})$$
• Matrix element:
$$\langle \Psi_{2}^{(\text{tot})} | \ \hat{\vec{D}}_{e} | \Psi_{1}^{(\text{tot})} \rangle = \underbrace{\langle \Psi_{2} | \ \hat{\vec{D}}_{e}(0) | \Psi_{1} \rangle}_{\text{purely electronic}} \underbrace{\langle \zeta_{2n_{2}}(Q) | \zeta_{1n_{1}}(Q) \rangle}_{\text{FC integrals}}$$

Electronic spectra

Franck-Condon integrals Which modes are important?

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- If  $\Delta Q_i = Q_{2,i}^{(0)} Q_{1,i}^{(0)} = 0$ : only zero-phonon line (ZPL) contributes • If  $\Delta Q_i = Q_{2,i}^{(0)} - Q_{1,i}^{(0)} \neq 0$ : non-trivial contribution
- Consider Taylor expansion of PESs:

$$U_{1}(Q) = U_{1}^{(0)} + \frac{1}{2} \sum_{i=1}^{3\nu-6} \left( \frac{\partial^{2}U_{1}}{\partial Q_{i}^{2}} \right)_{Q_{1}^{(0)}} (Q_{i} - Q_{1,i}^{(0)})^{2} + \dots$$
  

$$U_{2}(Q) = U_{2}^{(0)} + \sum_{i=1}^{3\nu-6} \left( \frac{\partial U_{2}}{\partial Q_{i}} \right)_{Q_{2}^{(0)}} (Q_{i} - Q_{2,i}^{(0)}) + \frac{1}{2} \sum_{i=1}^{3\nu-6} \left( \frac{\partial^{2}U_{2}}{\partial Q_{i}^{2}} \right)_{Q_{2}^{(0)}} (Q_{i} - Q_{2,i}^{(0)})^{2} + \dots$$

• Hellmann-Feynman:

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$$rac{\partial U_2}{\partial Q_i} = \langle \Psi_2 | rac{\partial \dot{H}_{el}}{\partial Q_i} | \Psi_2 
angle,$$

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• Contributing  $Q_i$ 's can be found from  $\Gamma(\Psi_2) \otimes \Gamma(Q_i) \otimes \Gamma(\Psi_2) \ni \Gamma_1$ 

Electronic spectra

# Franck-Condon integrals

#### Calculation

- Numerical integration
- Very simplistic, yet useful: single configurational coordinate model
  - One mode Q
  - Harmonic potentials
  - Equal curvatures

$$U_1(Q) = U_1(0) + \frac{1}{2}KQ^2$$
$$U_2(Q) = U_2(0) + A_2Q + \frac{1}{2}KQ^2$$

• FC integral:

J.J

$$\begin{aligned} \langle \zeta_{2n_2}(Q) | \zeta_{1n_1}(Q) \rangle &= \sqrt{\frac{n_1!}{n_2!}} (-S)^{n_2 - n_1} e^{-S/2} L_{n_1}^{n_2 - n_1}(S) \\ \langle \zeta_{20}(Q) | \zeta_{1n}(Q) \rangle &= \frac{e^{-S} S^n}{n!} \end{aligned}$$
  
with  $S = \frac{1}{2} \frac{M \omega_Q}{\hbar} \Delta Q^2$  and  $\omega_Q = \sqrt{\frac{K}{M}}.$ 

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# Spectral shape

• Sum over all electron-vibrational contributions:

$$S_{1\to2}(\omega) = C_{1\to2}\omega^{\beta} \sum_{i,j=1}^{3\nu-6} \sum_{n_{1i},n_{2j}=0}^{\infty} \left| \left\langle \zeta_{2n_{2j}}(Q_j) | \zeta_{1n_{1i}}(Q_i) \right\rangle \right|^2 \frac{\delta[E_0 \pm (n_{2j}\epsilon_{2j} - n_{1i}\epsilon_{1i}) - \hbar\omega]}{\mathrm{e}^{\frac{n_{1i}\epsilon_{1i}}{k_BT}} - 1}.$$

• For single configurational coordinate model:

$$C_{1\to 2}\omega^{\beta}\sum_{n=-\infty}^{\infty}\exp\left[\frac{n\hbar\omega_Q}{2k_BT} - S\coth\frac{\hbar\omega_Q}{2k_BT}\right] I_n\left[S\operatorname{csch}\left(\frac{\hbar\omega_Q}{2k_BT}\right)\right]\delta\left(E_0\pm n\hbar\omega_Q - \hbar\omega\right)$$

• Low-temperature limit:

$$S_{1\to 2}(\omega) = C_{1\to 2}\omega^{\beta} \sum_{n=0}^{\infty} \frac{\mathrm{e}^{-S}S^n}{n!} \delta(E_0 \pm n\hbar\omega - \hbar\omega),$$

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Electronic spectra

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# Oscillator strengths

• Oscillator strength:

$$f_{1\to 2} = C\omega_0 \left| \langle \Psi_2 | \, \hat{\vec{D}}_e \, | \Psi_1 \rangle \right|^2$$

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• Two flavours:

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- Dipole:  $f_{1\to2}^{D} = C\omega_0 e^2 \left| \langle \Psi_2 | \sum_{i=1}^N \vec{r_i} | \Psi_1 \rangle \right|^2$  Velocity:  $f_{1\to2}^{V} = C \frac{1}{\omega_0} \left( \frac{e\hbar}{m_e} \right)^2 \left| \langle \Psi_2 | \sum_{i=1}^N \vec{\nabla_i} | \Psi_1 \rangle \right|^2$   $f_{1\to2}^{D} = f_{1\to2}^{V}$  for exact eigenfunctions of  $\hat{H}$

## Oscillator strengths Molcas: RASSI-SO output













# Exercise: Potential energy curves en spectra for $YAG:Pr^{3+}$

- Introduce an "ad hoc" breathing mode,  $Q_{\rm Pr-O}$  for the  $\rm PrO_8^{13-}$  cluster in YAG.
- Perform single-point SA-CASSCF/MS-CASPT2/RASSI-SO calculations to get the states of the  $4f^2$  and 4f5d configurations for a range of geometries along the breathing mode.
- Construct potential energy curves for all states and plot the configuration coordinate diagram.
- Make a table with equilibrium values for  $Q_{Pr-O}$  and vibrational frequencies for a small number of relevant electronic states.
- From the RASSI-SO output, extract oscillator strengths and simulate the absorption spectrum of  $Pr^{3+}$  in YAG.

Electronic spectra

• Simulate the emission spectrum from the lowest 4f5d state.

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# Literature

